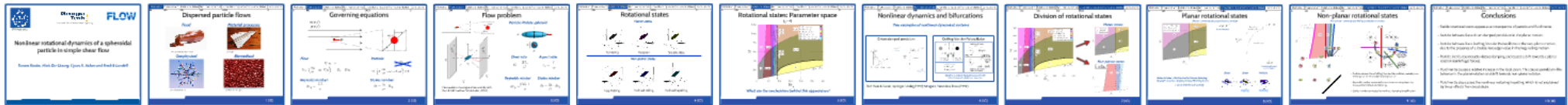


# Thank you!





KTH Mechanics



G.W. Woodruff School of Mechanical Engineering



# Nonlinear rotational dynamics of a spheroidal particle in simple shear flow

*Tomas Rosén, Minh Do-Quang, Cyrus K. Aidun and Fredrik Lundell*

# Dispersed particle flows

## Food



Ref: GorillaAttack/Shutterstock

## Material processes



Ref: Motif

## Geophysical



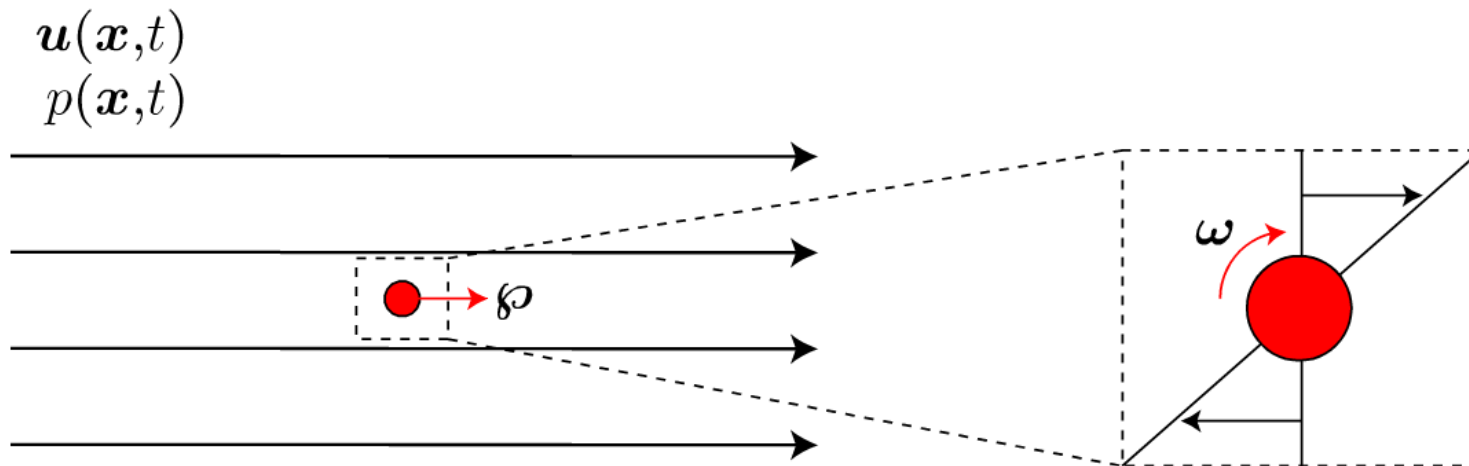
Ref: Snowcrystals.com

## Biomedical



Ref: Wellcome Images

# Governing equations



**Flow:**

$$\nabla \cdot \mathbf{u} = 0$$

$$Re_p \cdot \left( \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = -\nabla p + \nabla^2 \mathbf{u}$$

**Reynolds number:**

$$Re_p = \frac{UL}{\nu}$$

**Particle:**

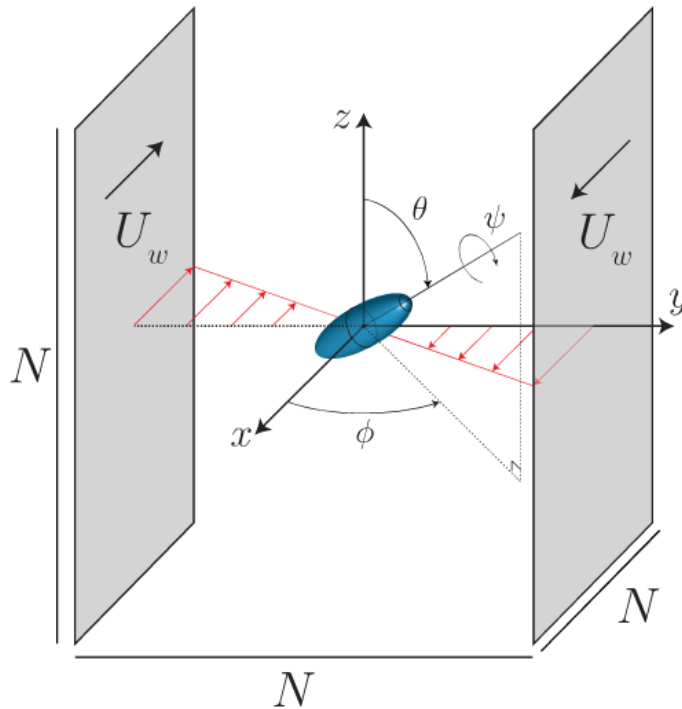
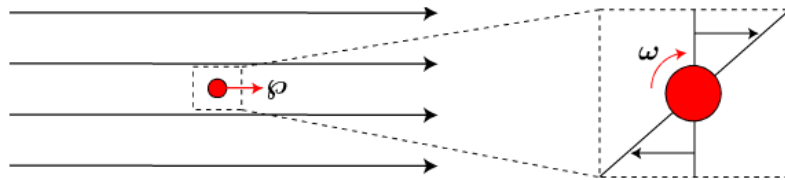
~~$$St \cdot \frac{\partial \boldsymbol{\rho}}{\partial t} = \mathbf{F}$$~~

$$St \cdot \left( \mathbf{I} \cdot \frac{\partial \boldsymbol{\omega}}{\partial t} + \boldsymbol{\omega} \times (\mathbf{I} \cdot \boldsymbol{\omega}) \right) = \mathbf{T}$$

**Stokes number:**

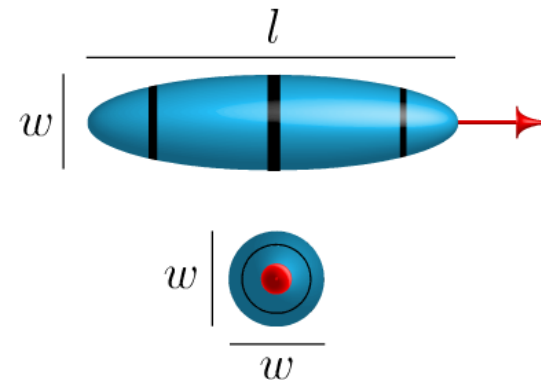
$$St = \frac{\rho_p}{\rho_f} \cdot Re_p$$

# Flow problem



Flow problem investigated numerically with the LB-EBF method (Wu & Aidun, 2010)

*Particle: Prolate spheroid*



*Shear rate*

$$G = \frac{2U_w}{N}$$

*Aspect ratio*

$$r_p = \frac{l}{w}$$

*Reynolds number*

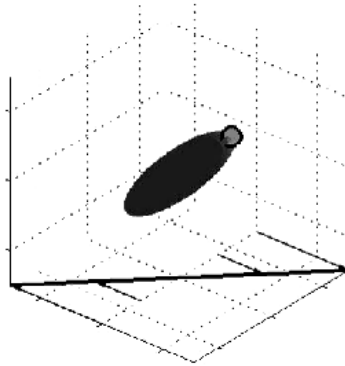
$$Re_p = \frac{(Gl) \cdot l}{\nu}$$

*Stokes number*

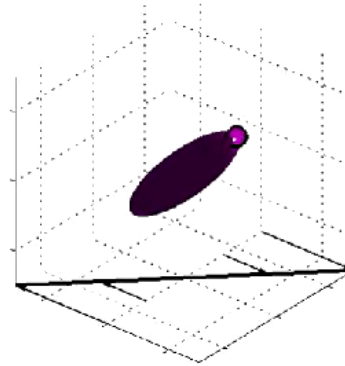
$$St = \frac{\rho_p}{\rho_f} \cdot Re_p$$

# Rotational states

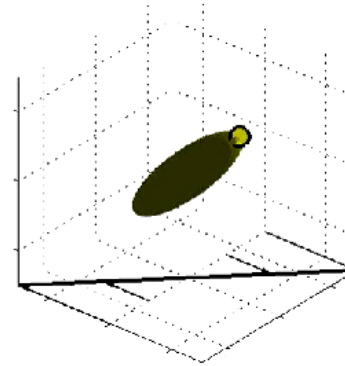
## Planar states



Tumbling

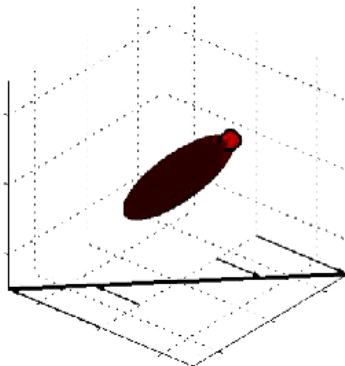


Rotation

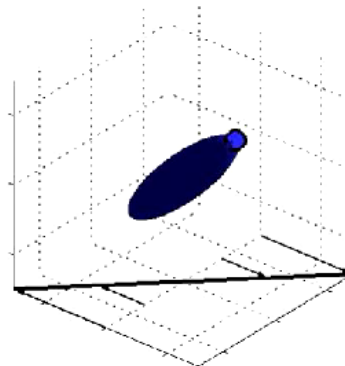


Steady state

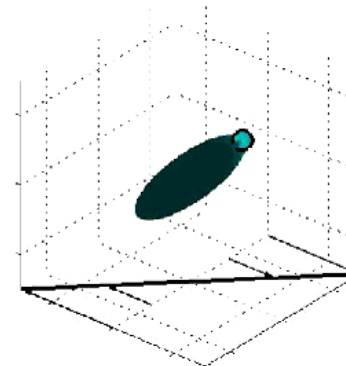
## Non-planar states



Log-Rolling



Inclined rolling

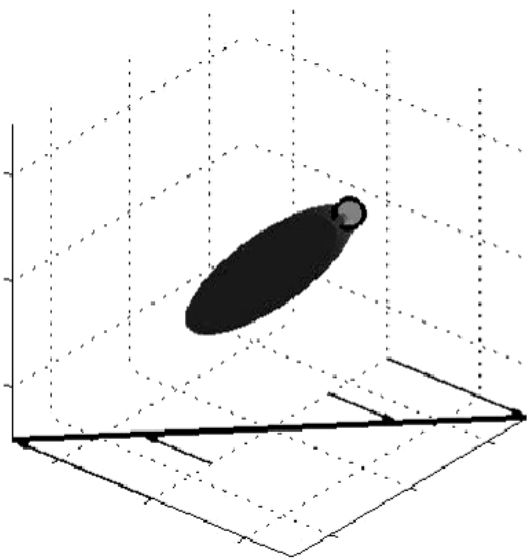


Inclined kayaking

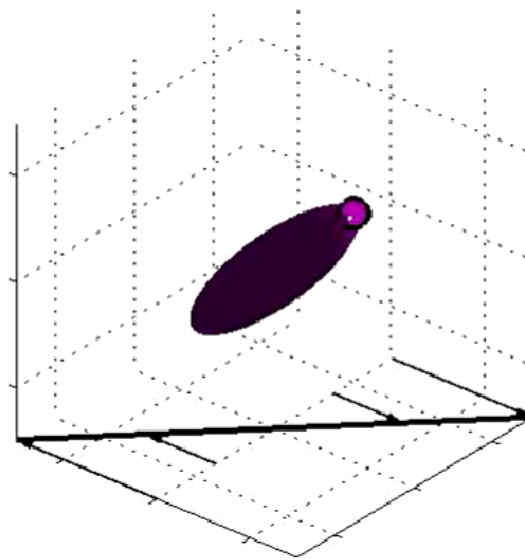


# Rotational states

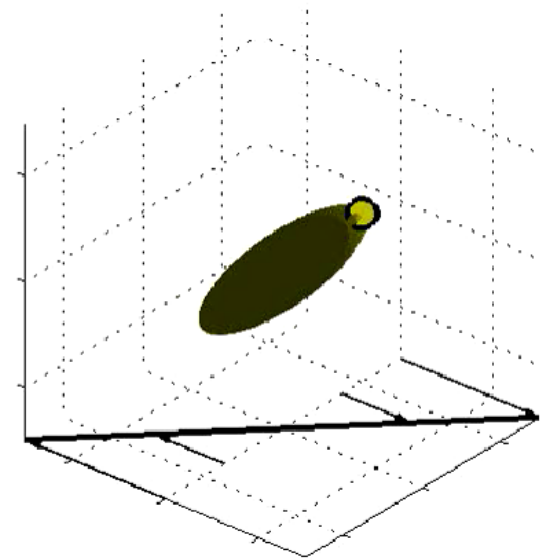
## *Planar states*



Tumbling

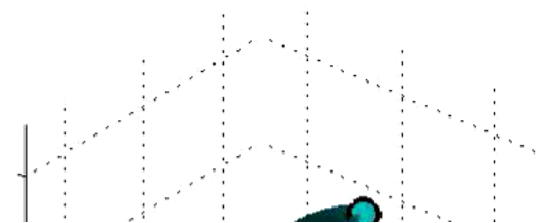
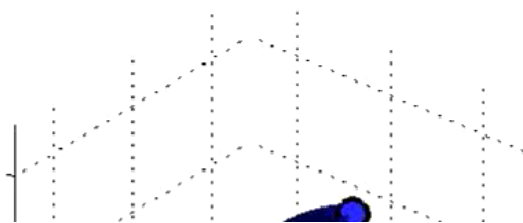
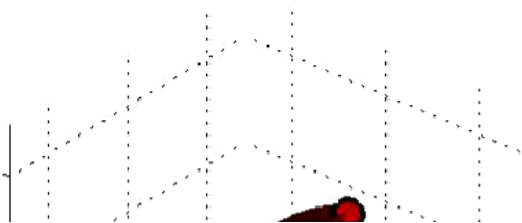


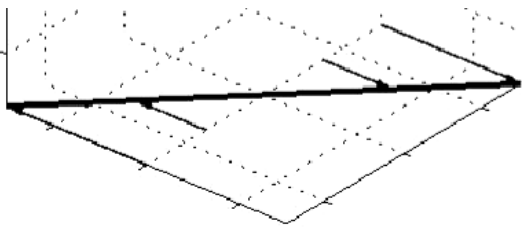
Rotation



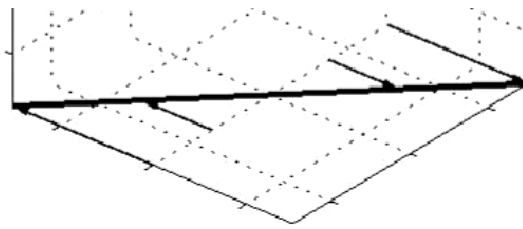
Steady state

## *Non-planar states*

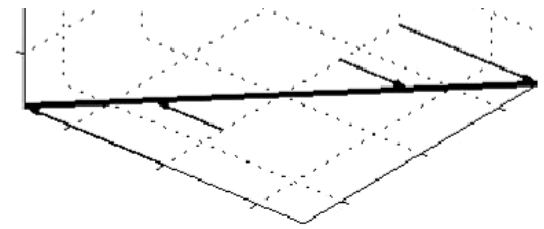




Tumbling

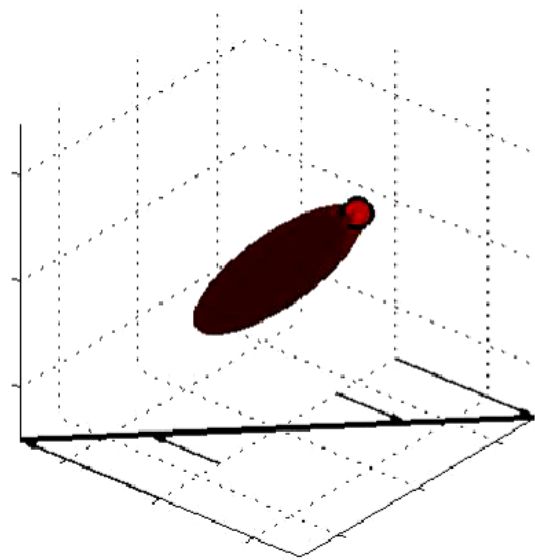


Rotation

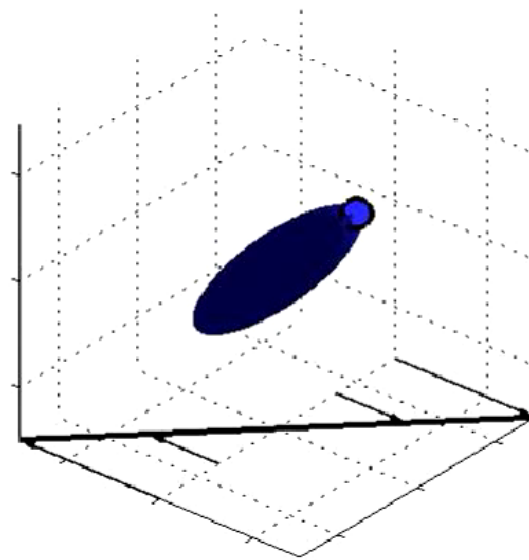


Steady state

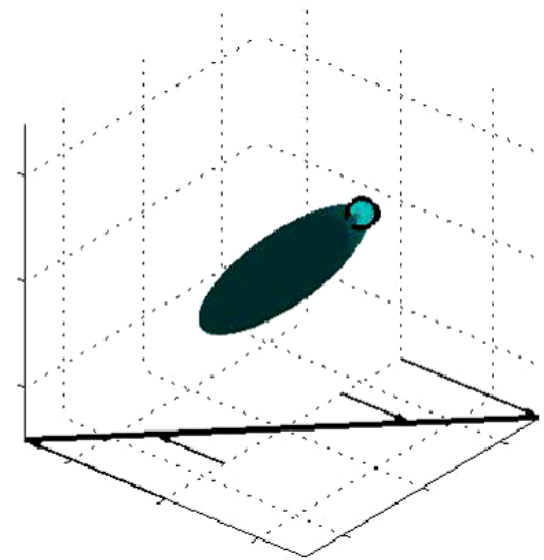
*Non-planar states*



Log-Rolling



Inclined rolling

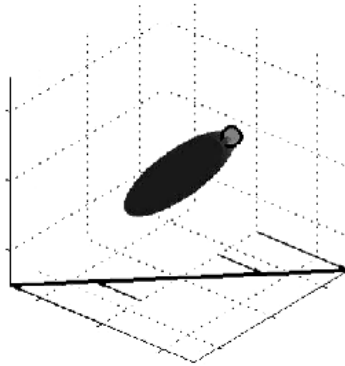


Inclined kayaking

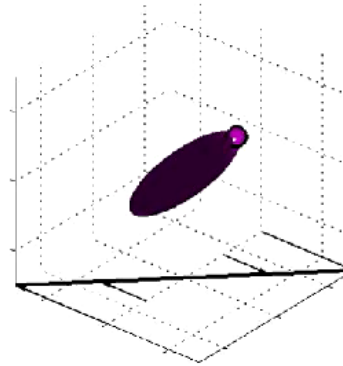


# Rotational states

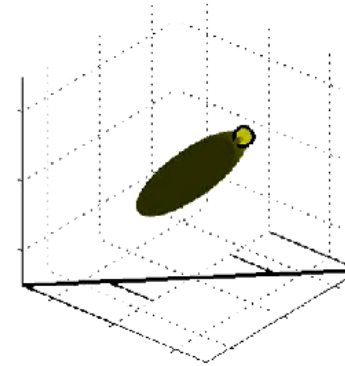
## Planar states



Tumbling

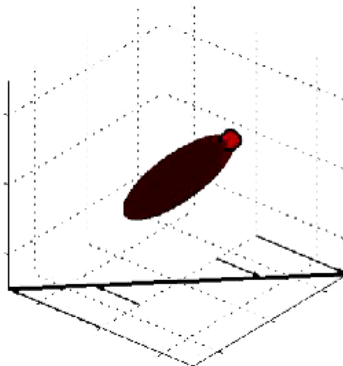


Rotation

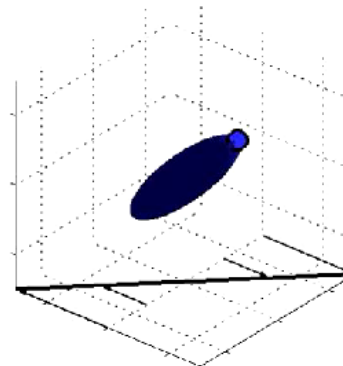


Steady state

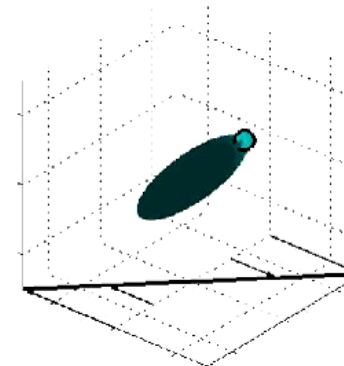
## Non-planar states



Log-Rolling



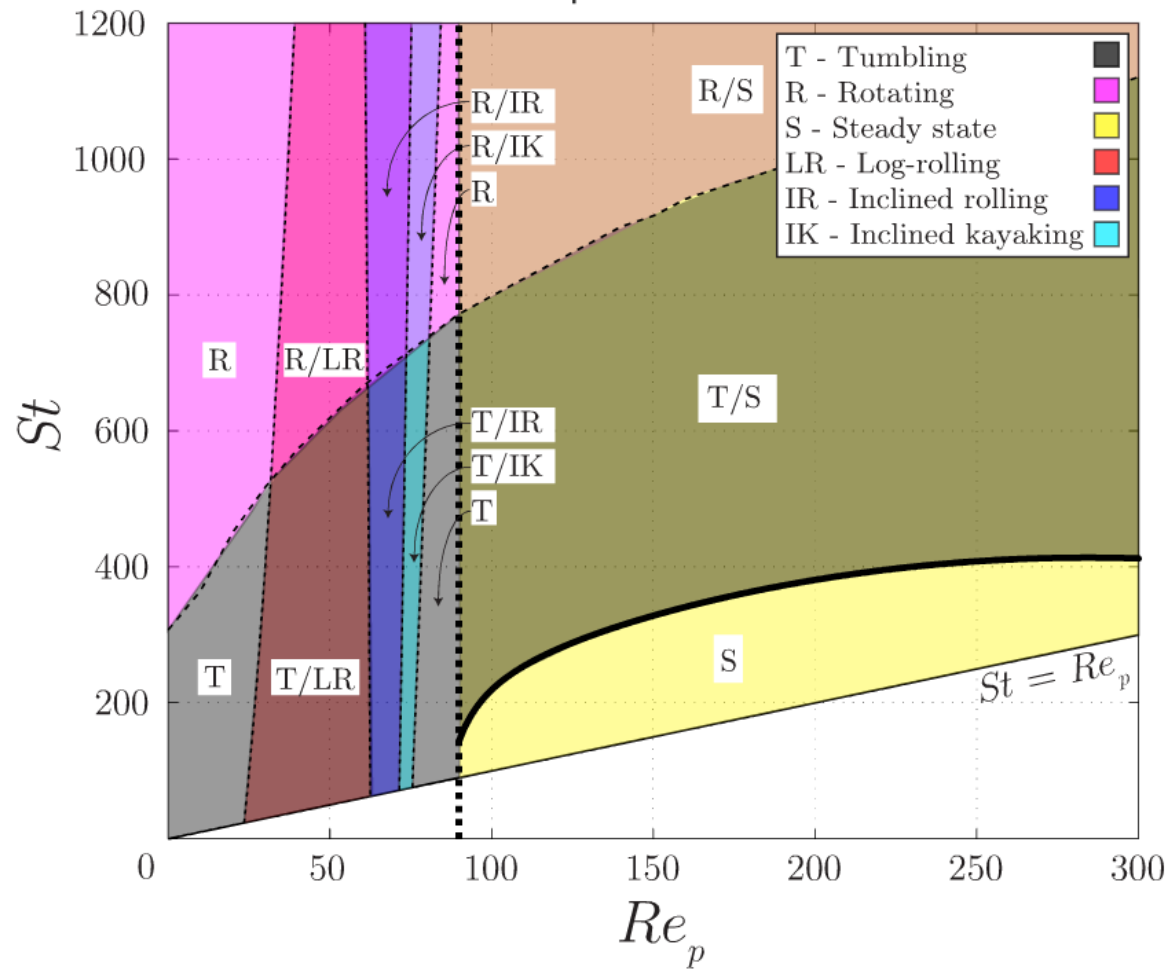
Inclined rolling



Inclined kayaking

# Rotational states: Parameter space

( $r_p = 4$ )

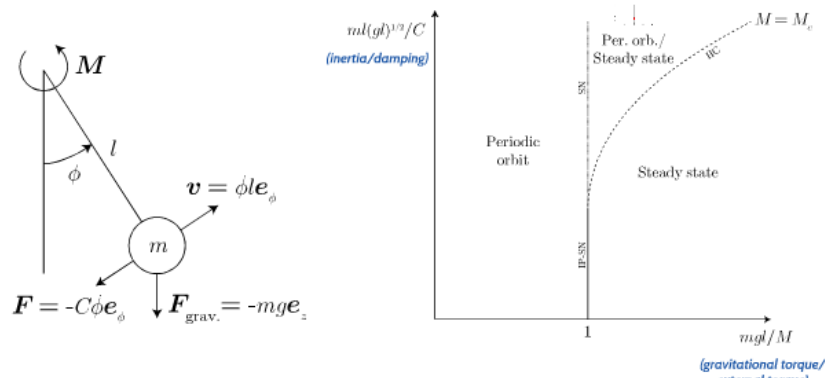


*What are the mechanisms behind this appearance?*

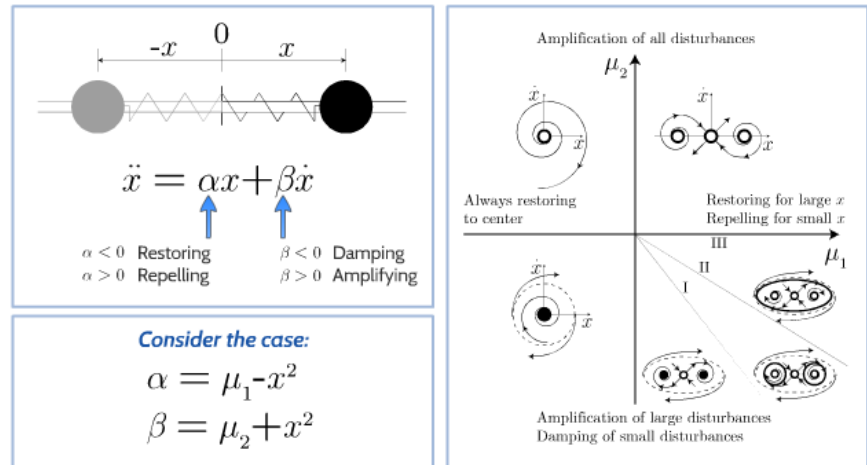
# Nonlinear dynamics and bifurcations

## Two examples of nonlinear dynamical systems

### Driven damped pendulum



### Duffing-Van der Pol oscillator



Consider the case:

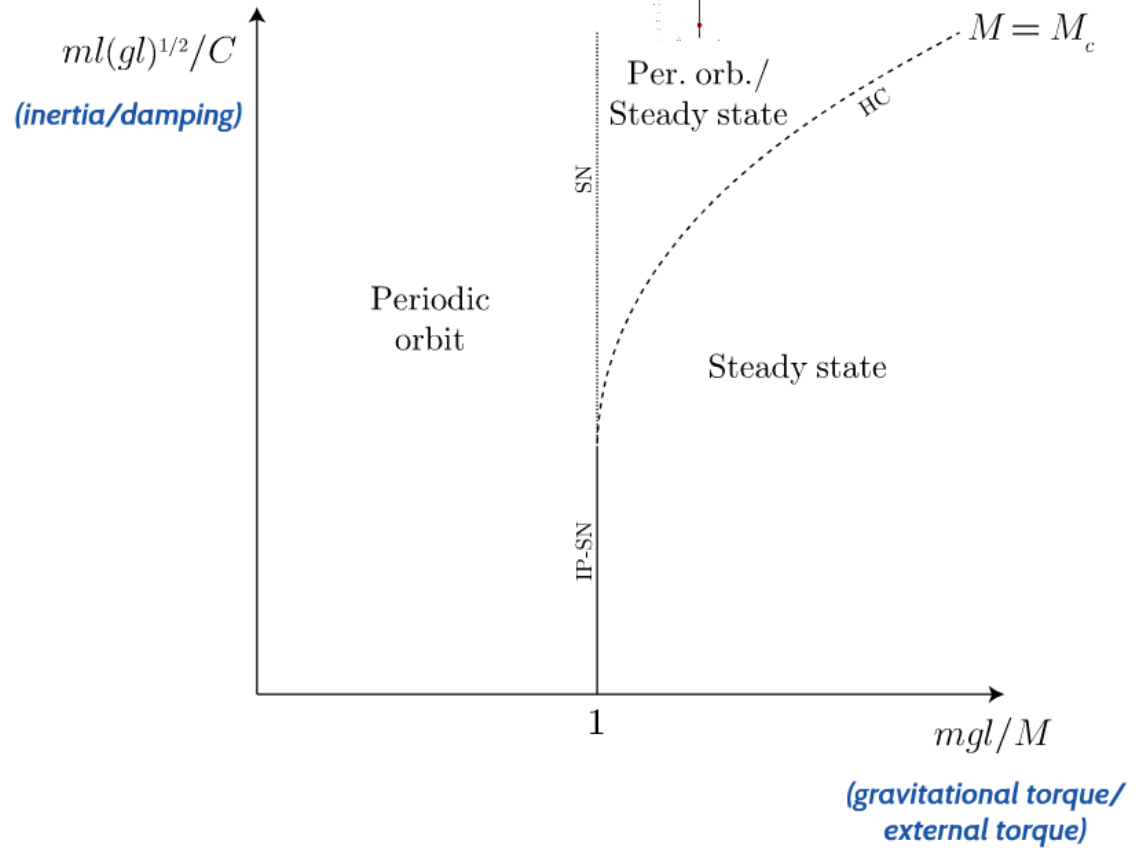
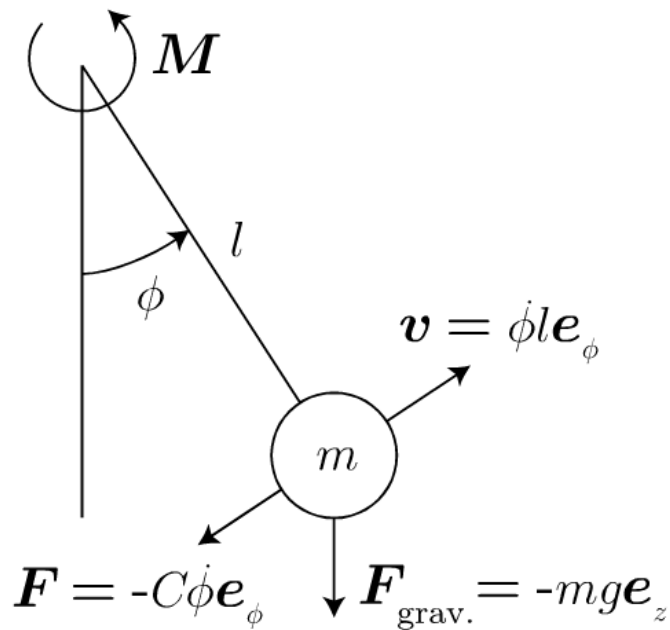
$$\alpha = \mu_1 - x^2$$

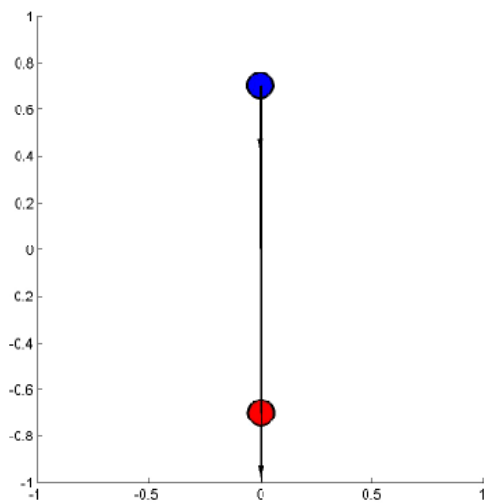
$$\beta = \mu_2 + x^2$$

Typical for a nonlinear dynamical system with odd symmetry around a double-zero eigenvalue

Ref: Hale & Kocak, Springer-Verlag (1991); Strogatz, Westview Press (1994).

# Driven damped pendulum

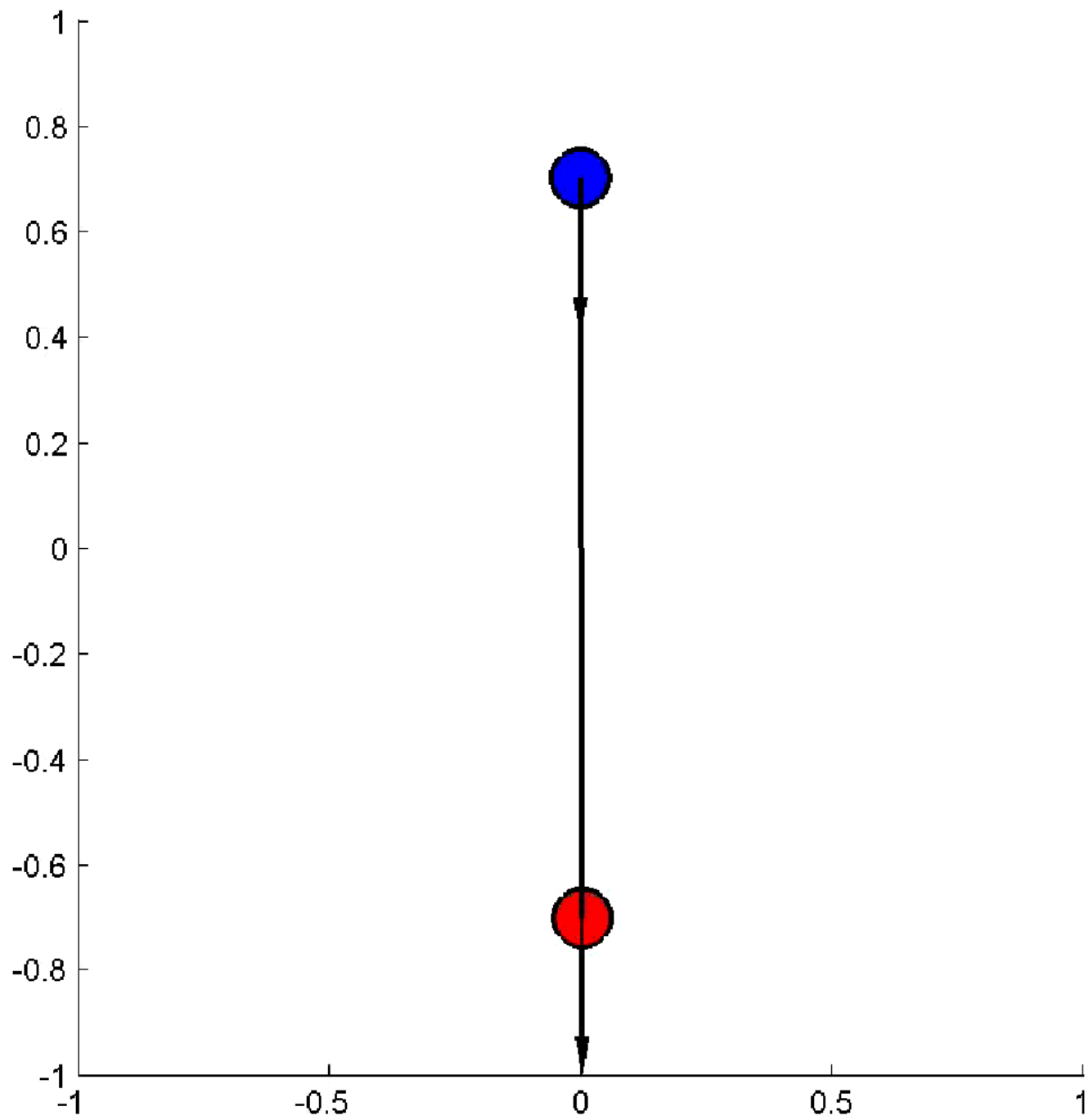




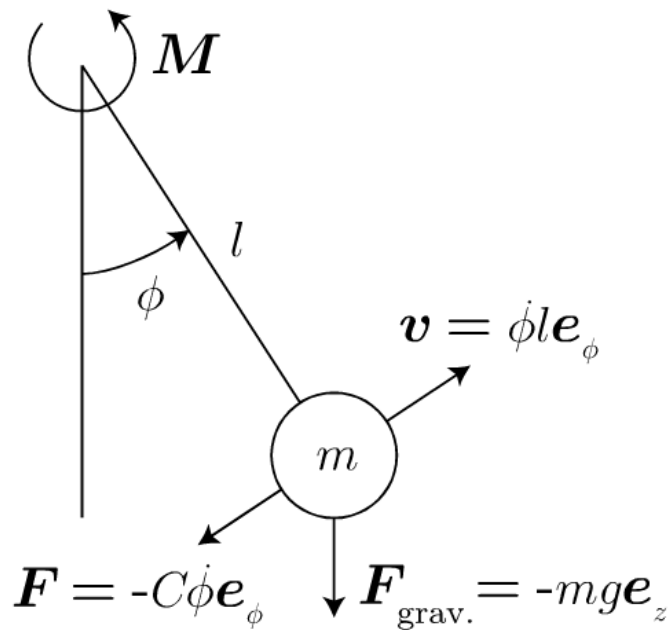
Per. orb./  
Steady state

SN

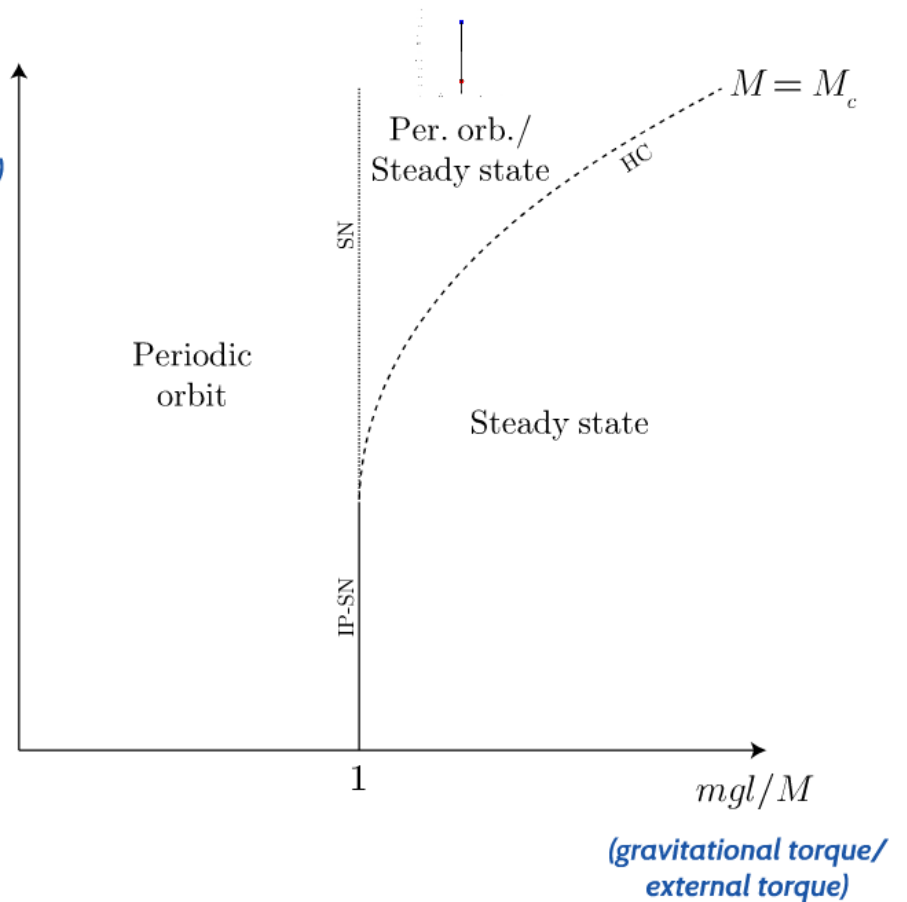
HC



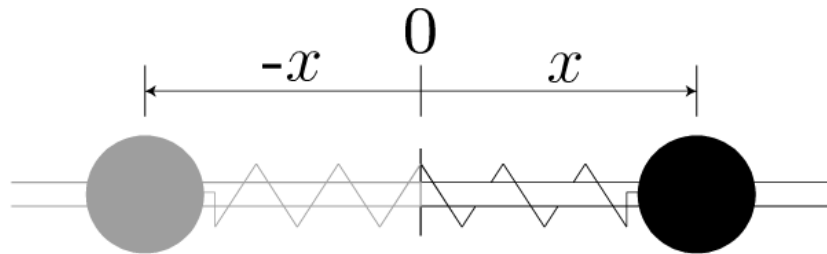
# Driven damped pendulum



$ml(gl)^{1/2}/C$   
(inertia/damping)



# Duffing-Van der Pol oscillator



$$\ddot{x} = \alpha x + \beta \dot{x}$$

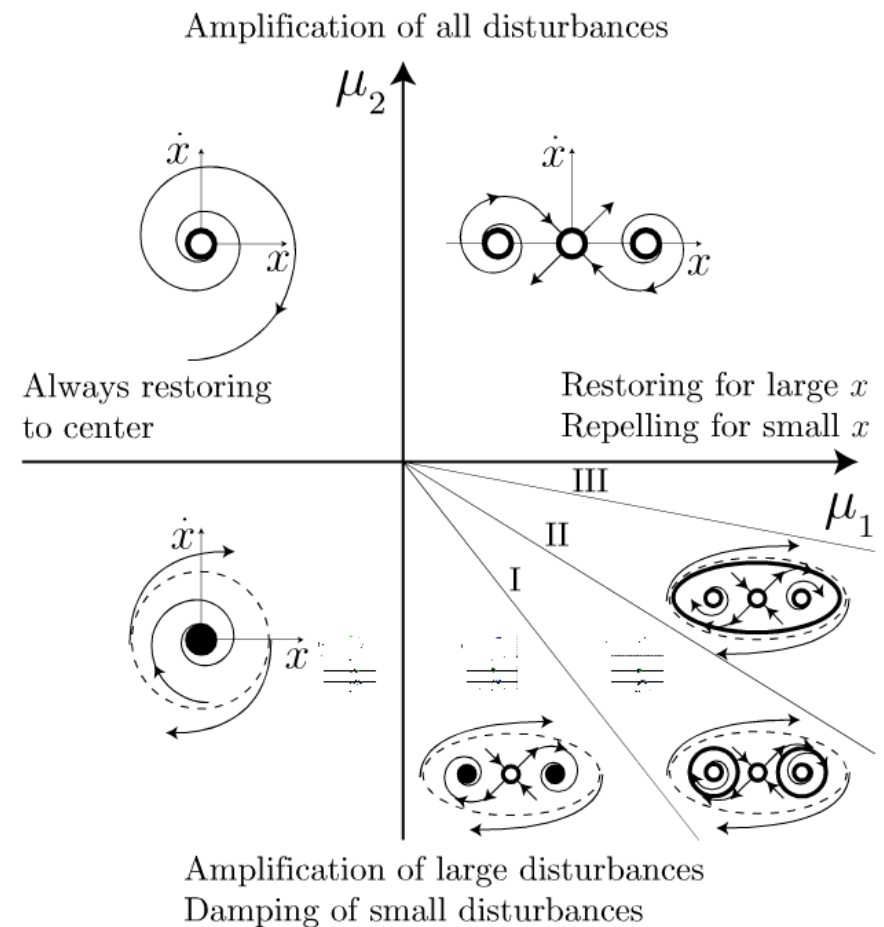
$\alpha < 0$  Restoring  
 $\alpha > 0$  Repelling

$\beta < 0$  Damping  
 $\beta > 0$  Amplifying

**Consider the case:**

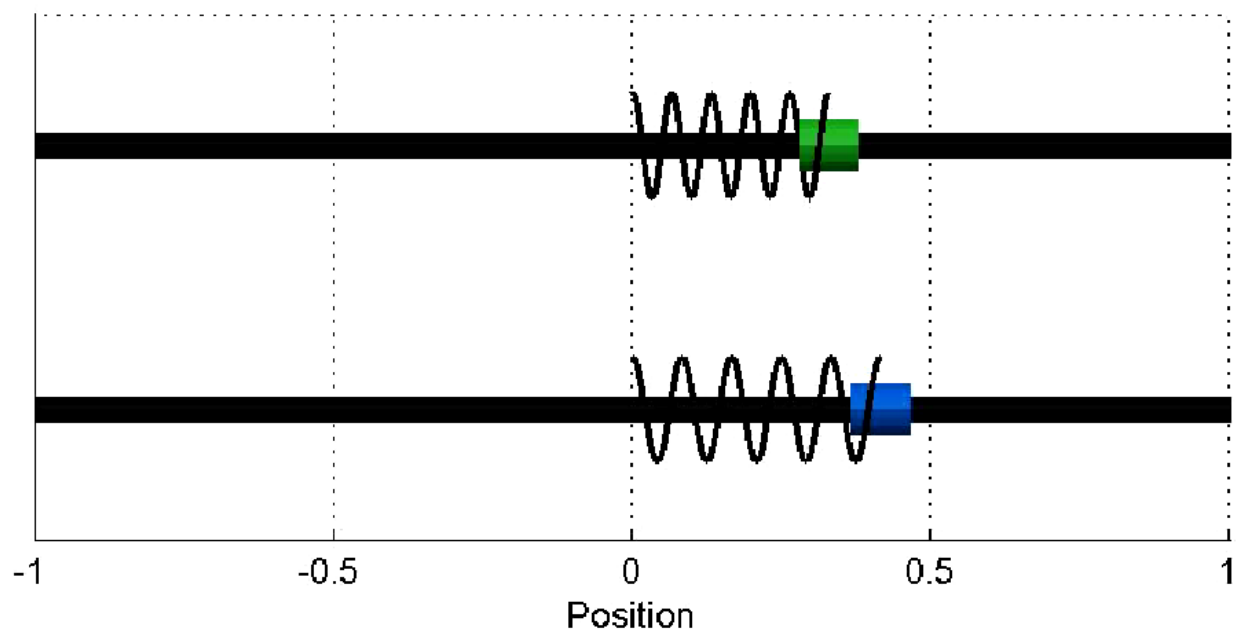
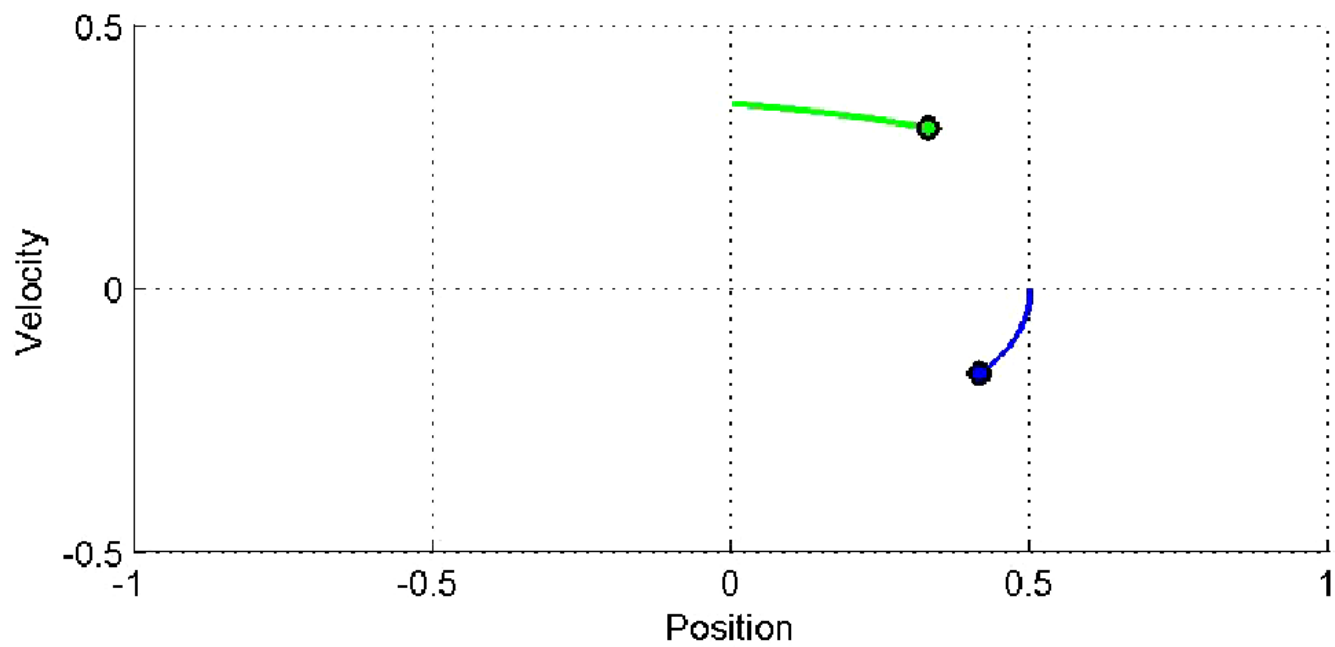
$$\alpha = \mu_1 - x^2$$

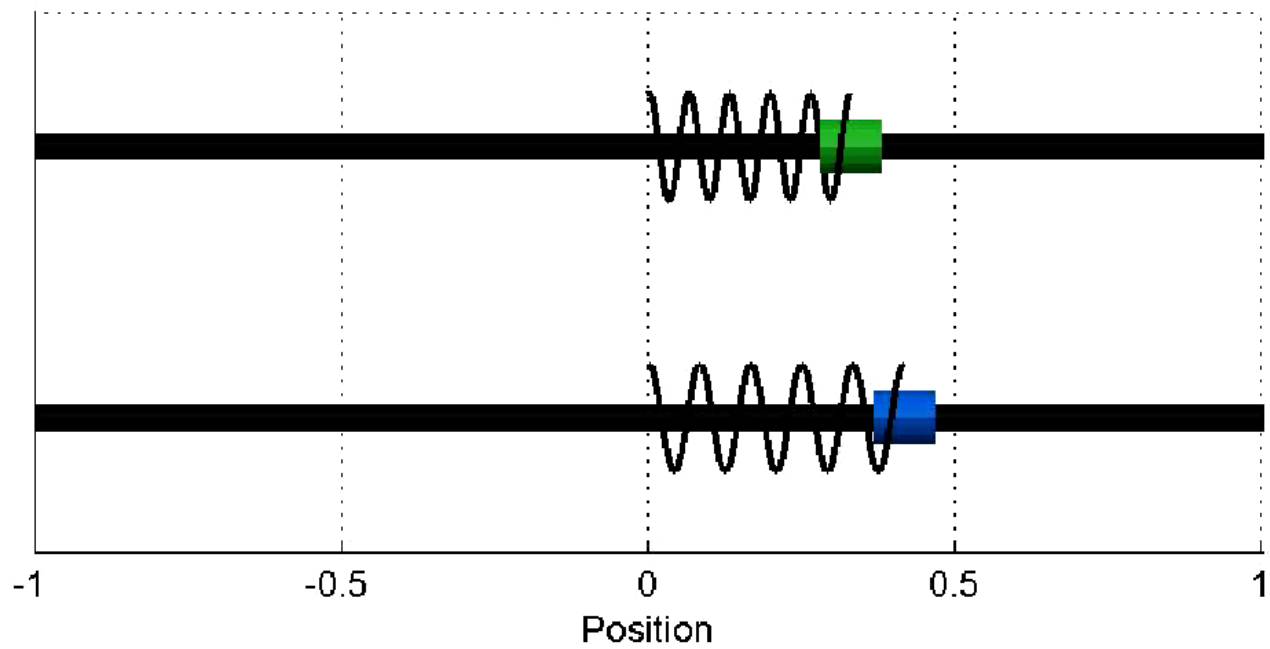
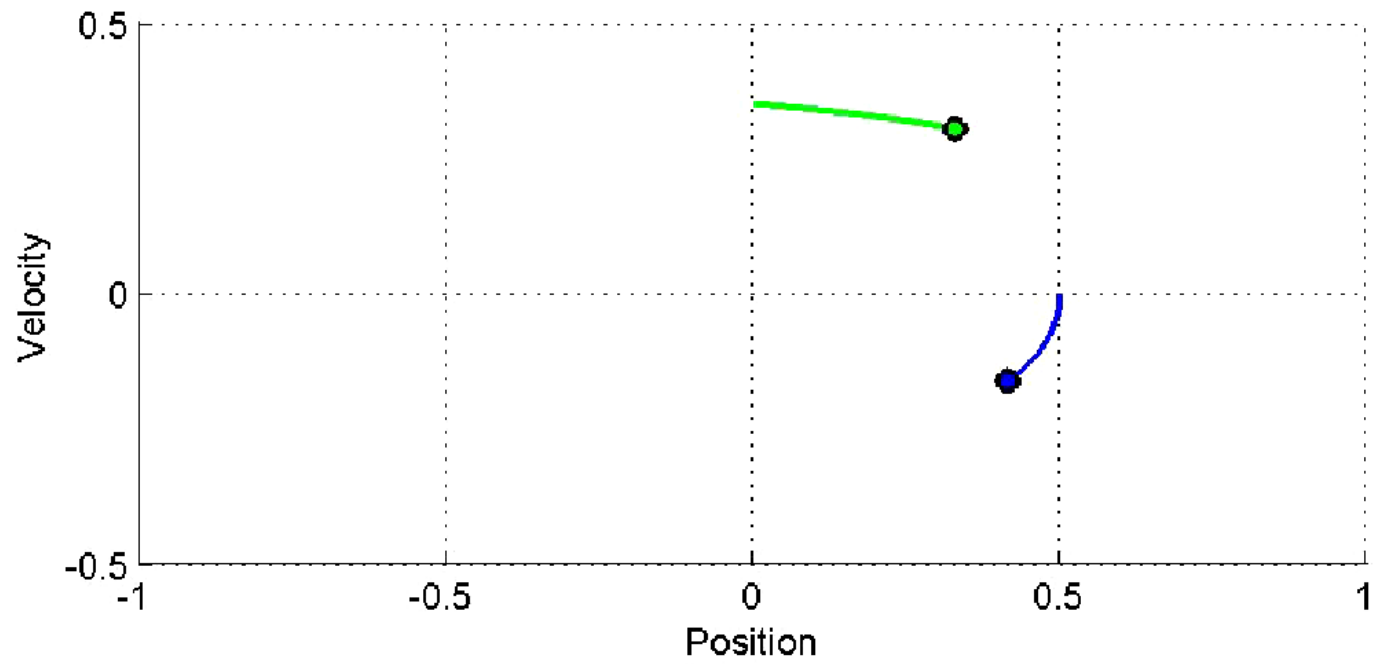
$$\beta = \mu_2 + x^2$$

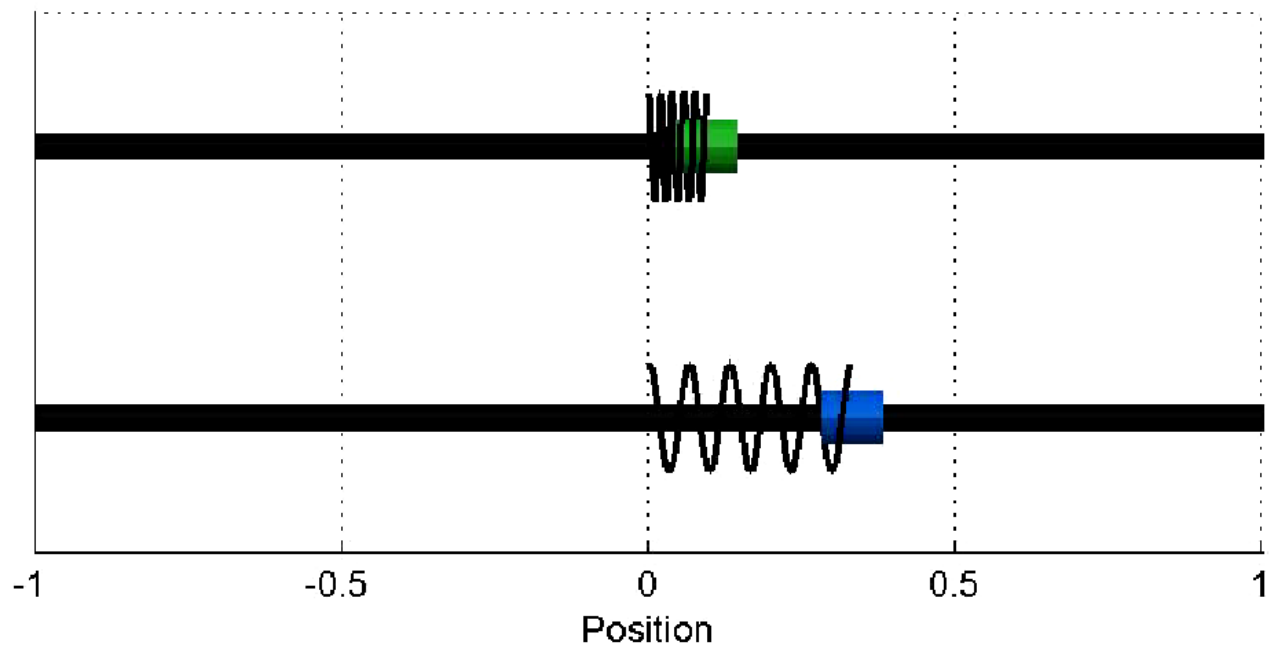
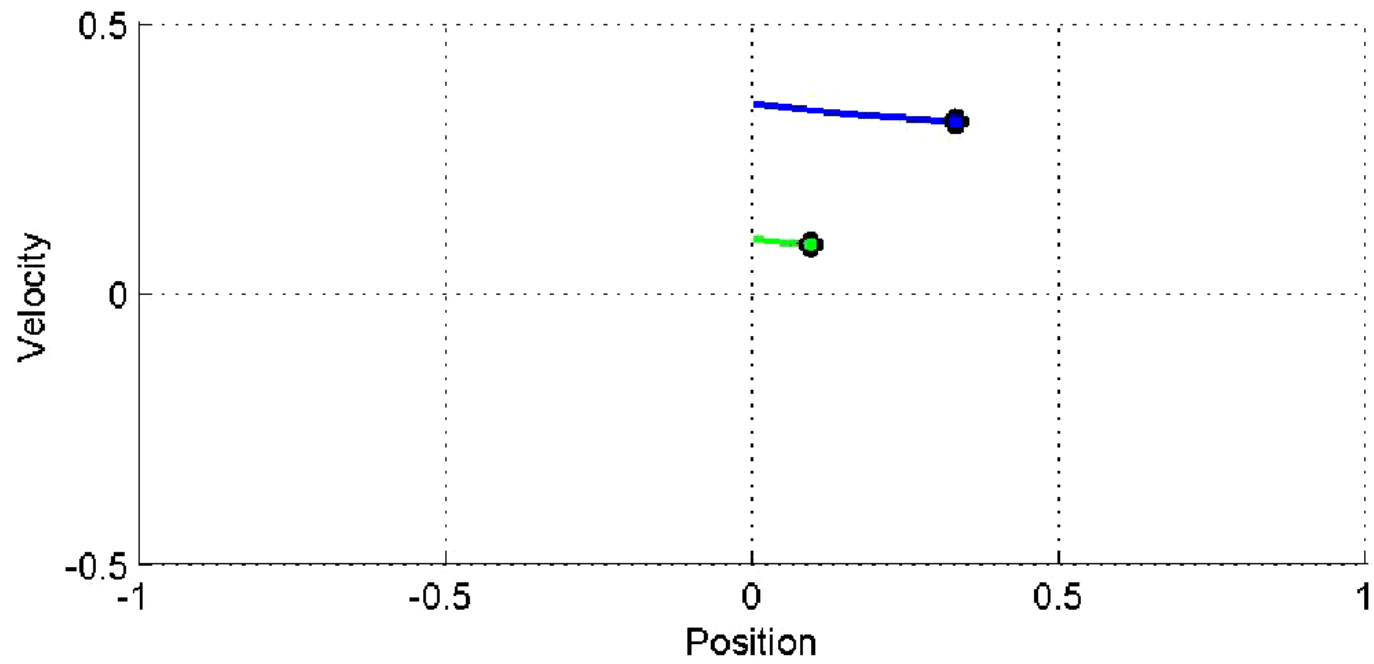


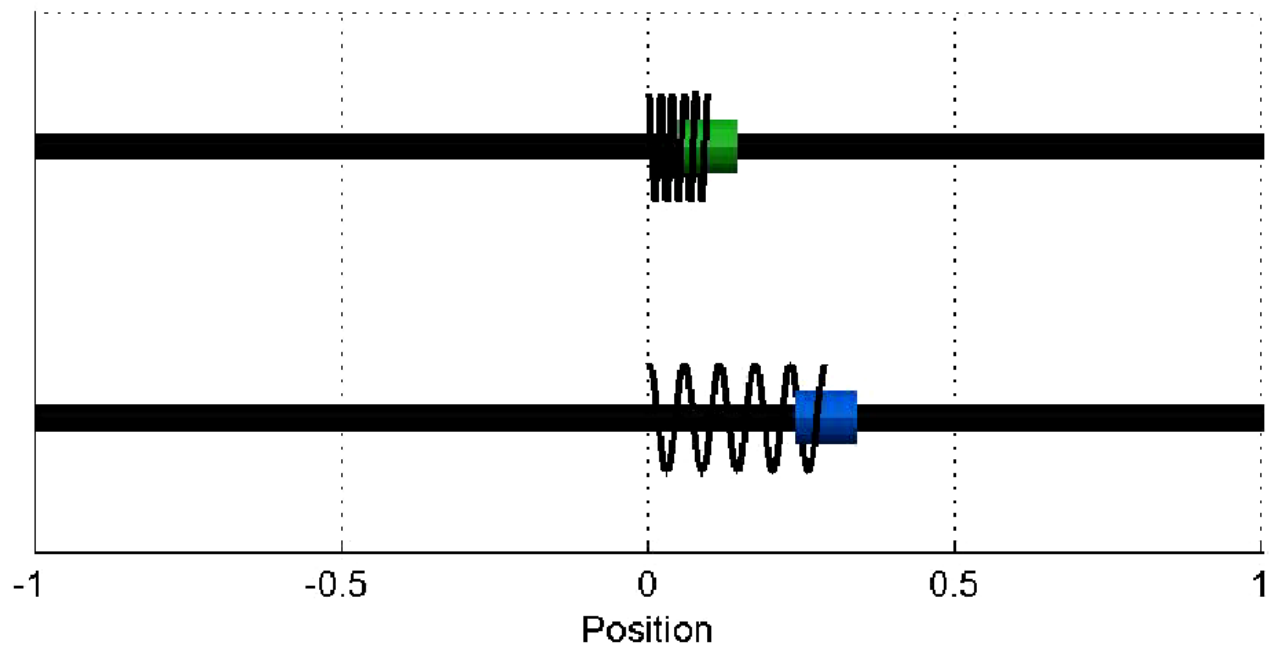
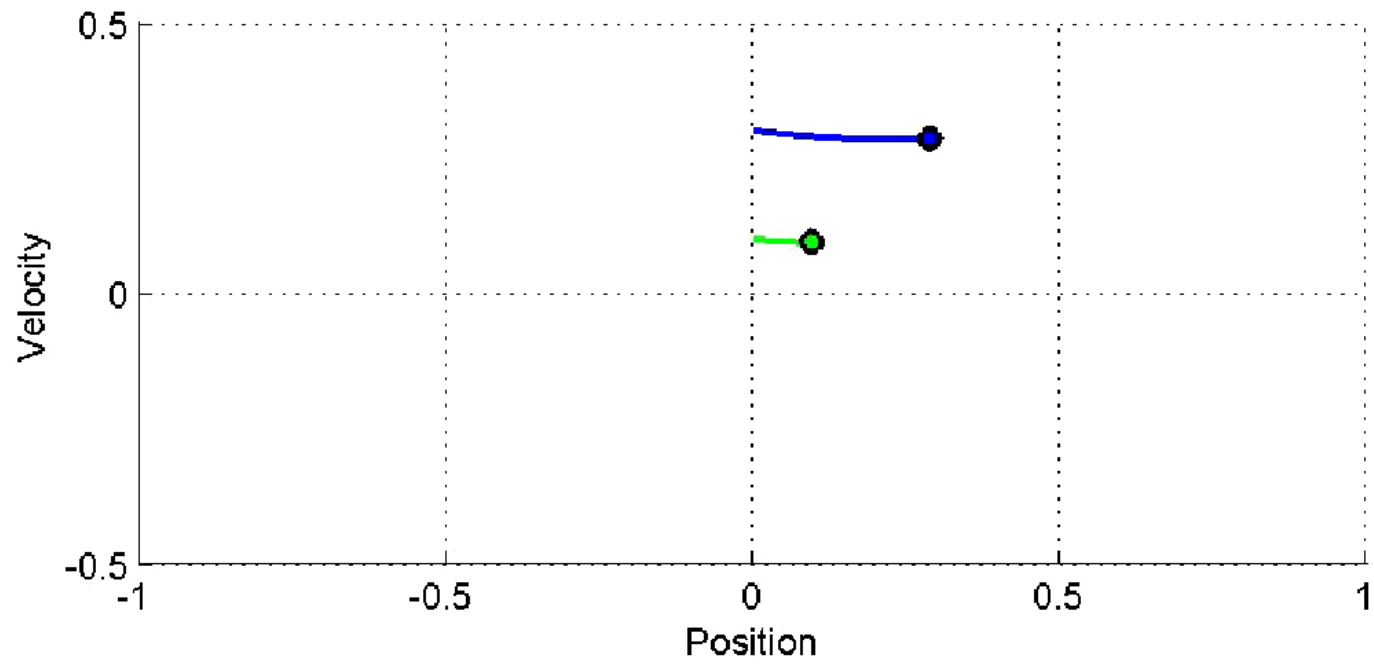
**Typical for a nonlinear dynamical system with odd symmetry around a double-zero eigenvalue**



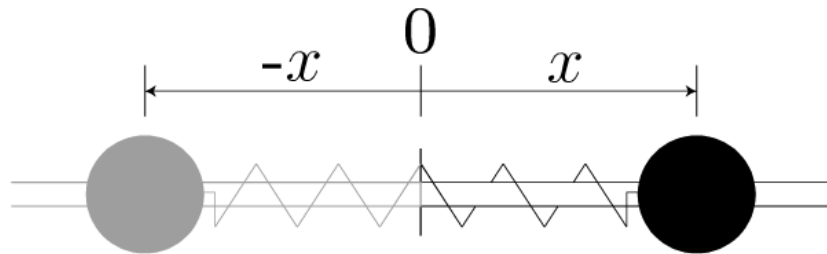








# Duffing-Van der Pol oscillator



$$\ddot{x} = \alpha x + \beta \dot{x}$$

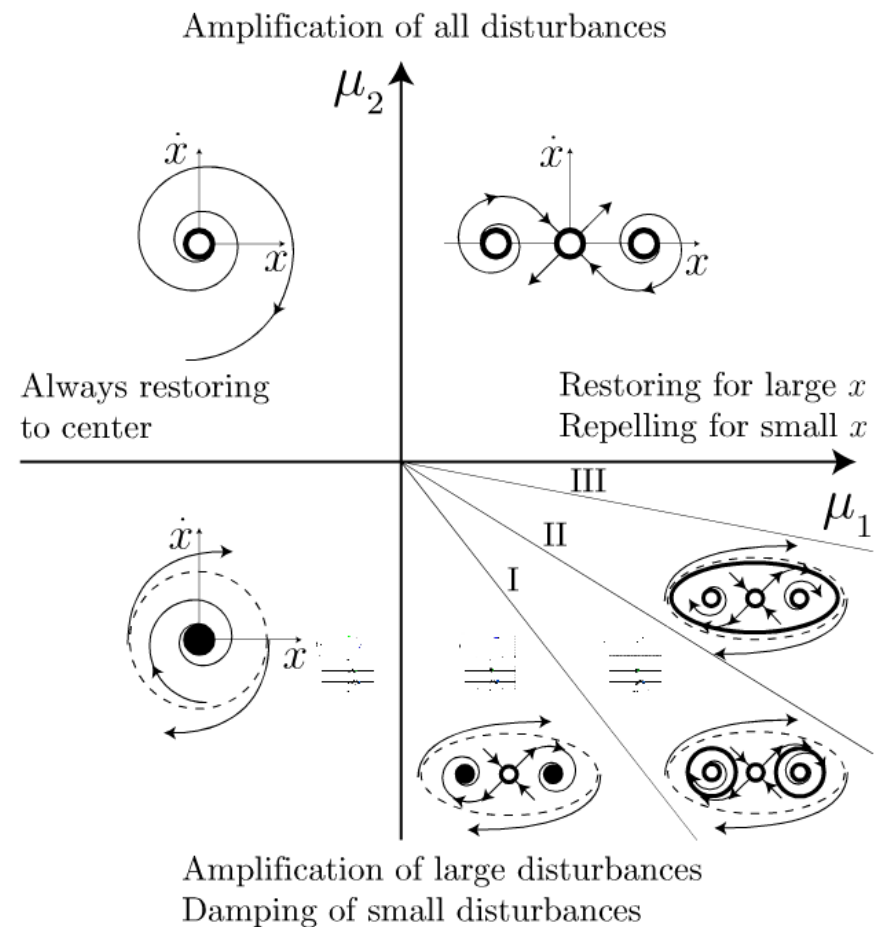
$\alpha < 0$  Restoring  
 $\alpha > 0$  Repelling

$\beta < 0$  Damping  
 $\beta > 0$  Amplifying

**Consider the case:**

$$\alpha = \mu_1 - x^2$$

$$\beta = \mu_2 + x^2$$

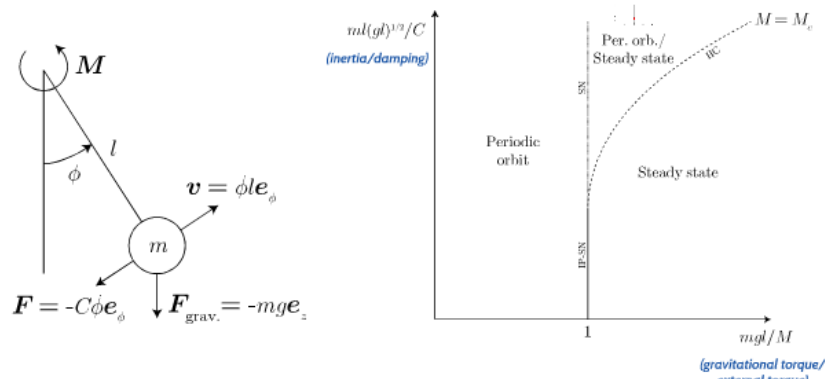


**Typical for a nonlinear dynamical system with odd symmetry around a double-zero eigenvalue**

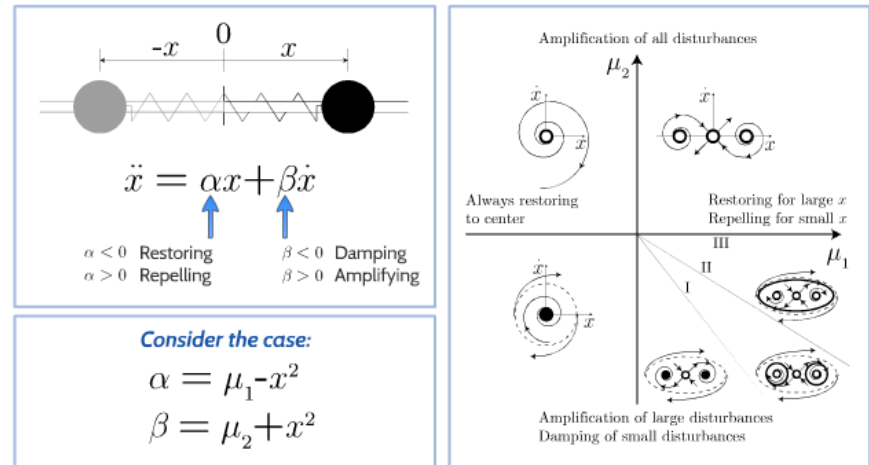
# Nonlinear dynamics and bifurcations

## Two examples of nonlinear dynamical systems

### Driven damped pendulum



### Duffing-Van der Pol oscillator



Consider the case:

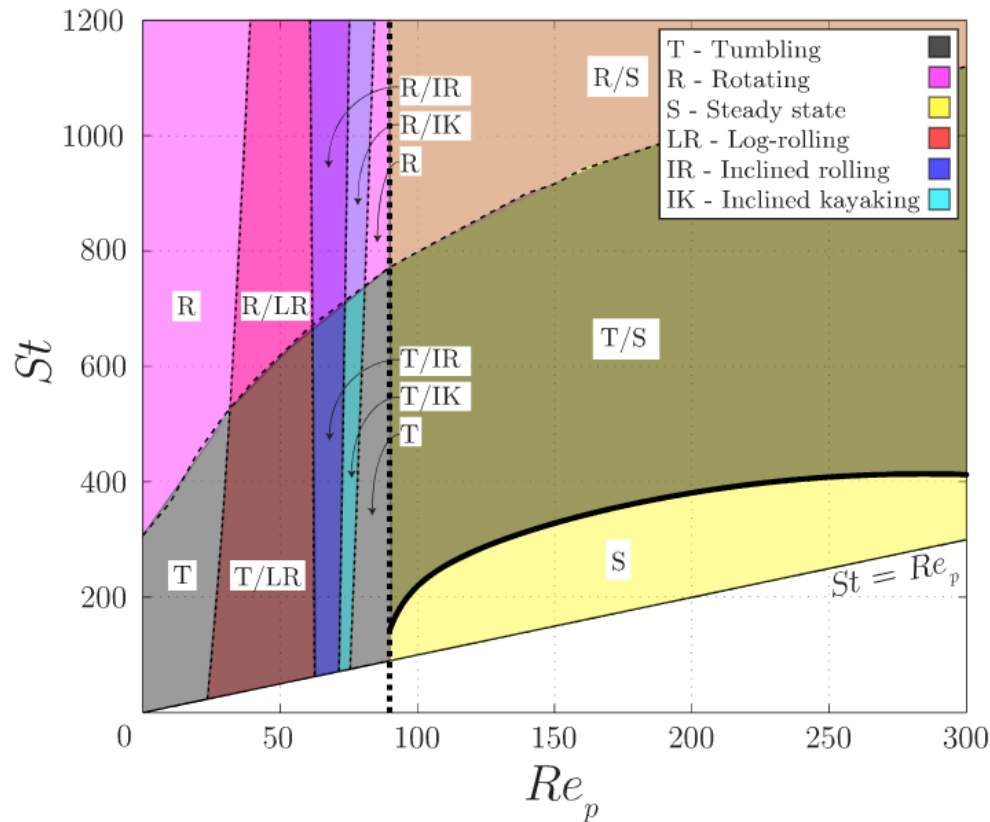
$$\alpha = \mu_1 - x^2$$

$$\beta = \mu_2 + x^2$$

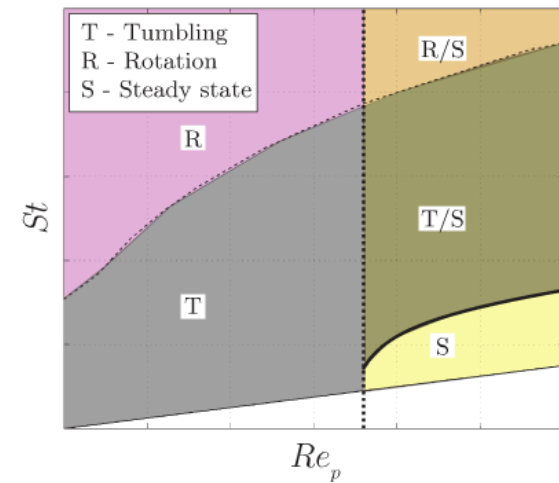
Typical for a nonlinear dynamical system with odd symmetry around a double-zero eigenvalue

Ref: Hale & Kocak, Springer-Verlag (1991); Strogatz, Westview Press (1994).

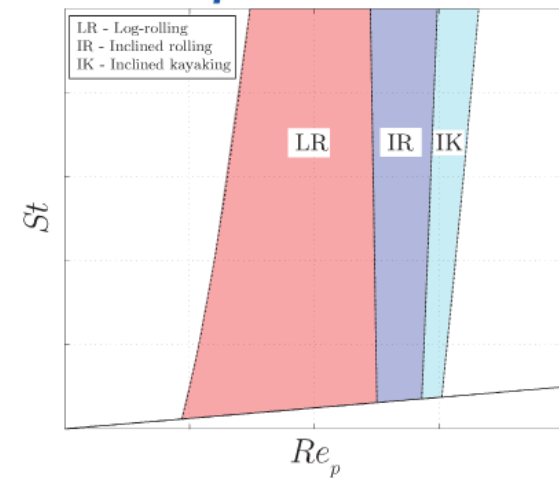
# Division of rotational states



## Planar states

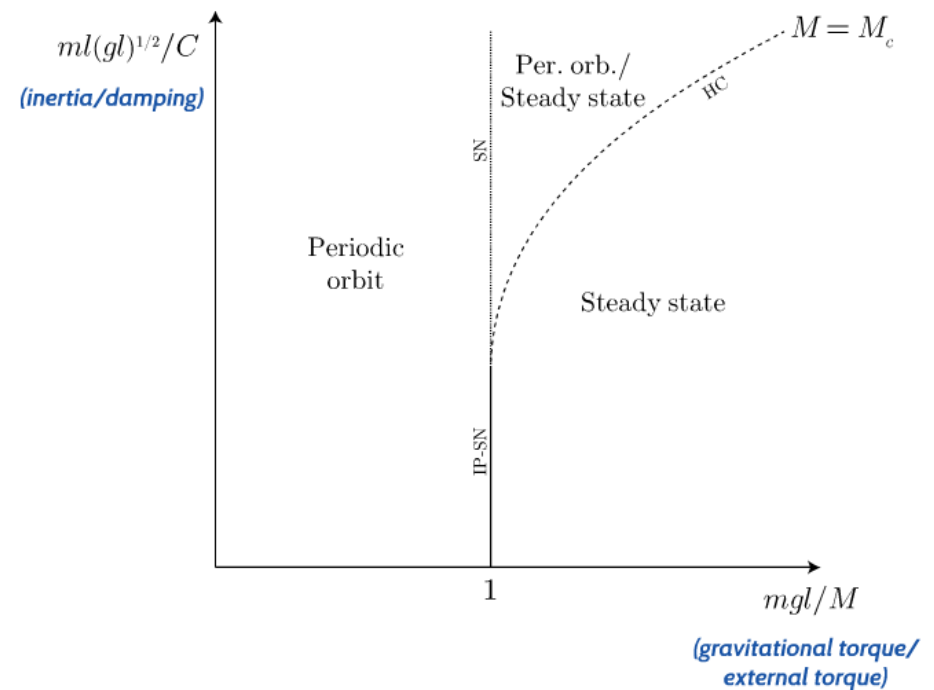
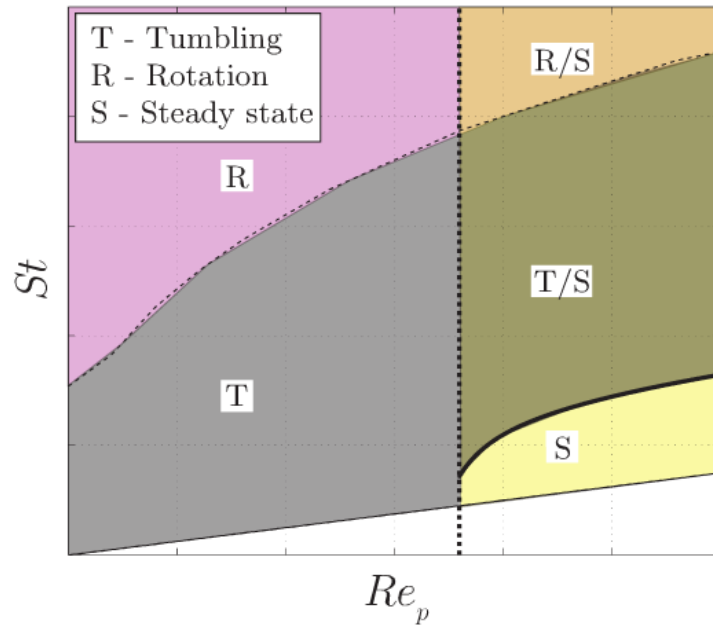


## Non-planar states

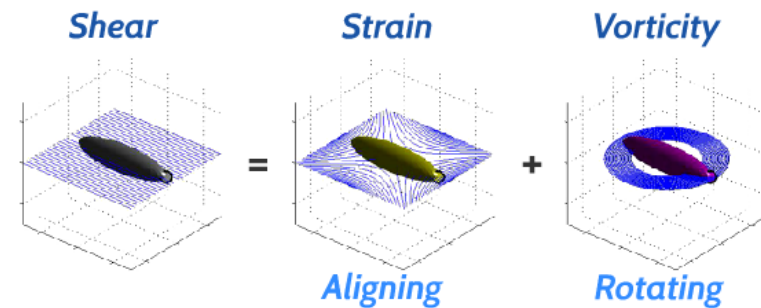


# Planar rotational states

(Planar = symmetry axis perpendicular to vorticity)



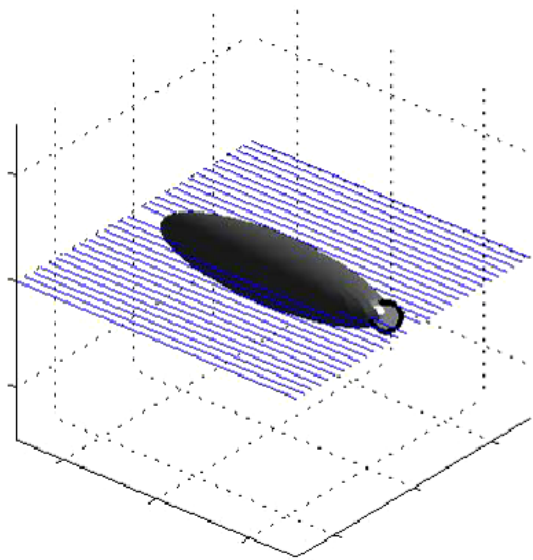
Stokes number = Particle inertia/viscous damping  
 Reynolds number = Aligning torque/Rotating torque  
 = Local strain/Local vorticity





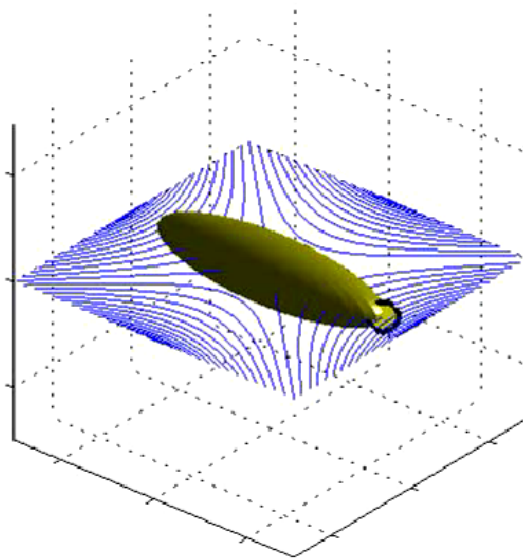
(gravitational torque  
external torque)

**Shear**



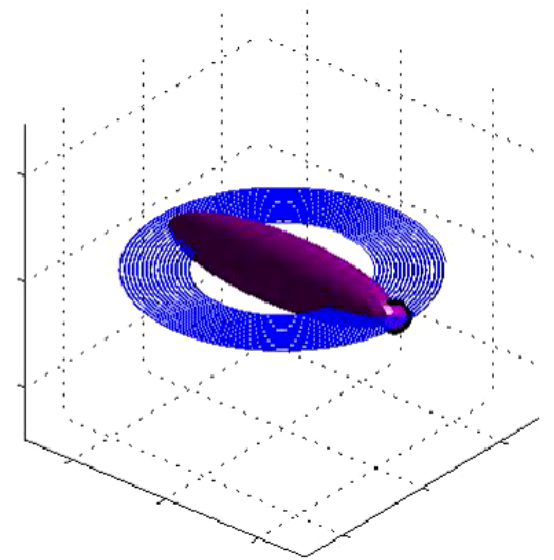
=

**Strain**



+

**Vorticity**

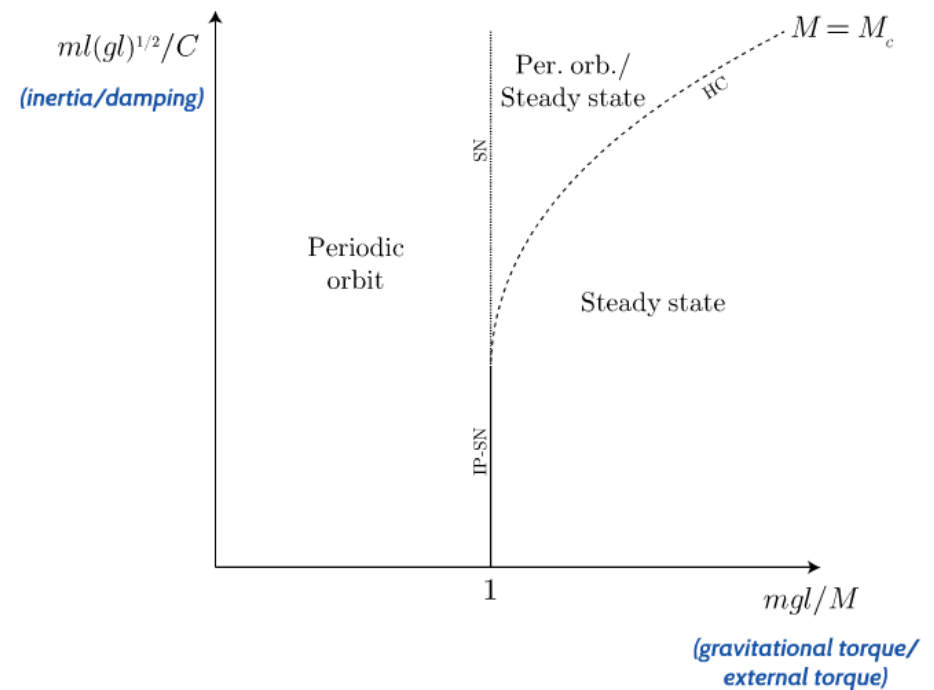
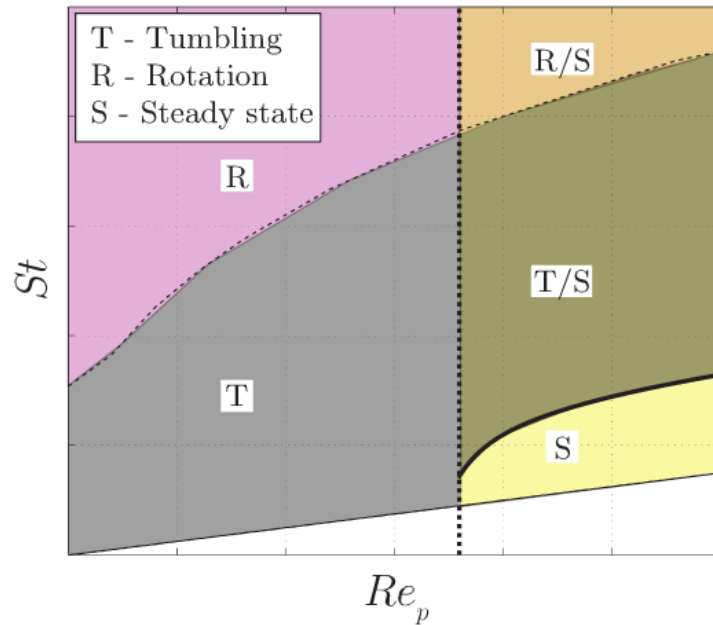


**Aligning**

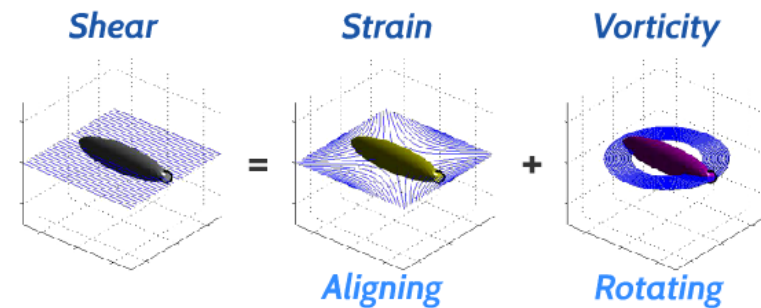
**Rotating**

# Planar rotational states

(Planar = symmetry axis perpendicular to vorticity)

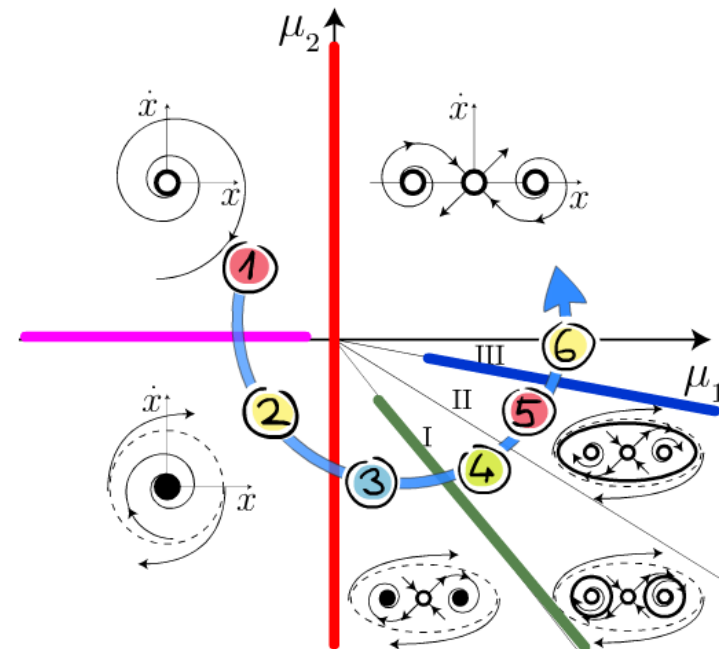
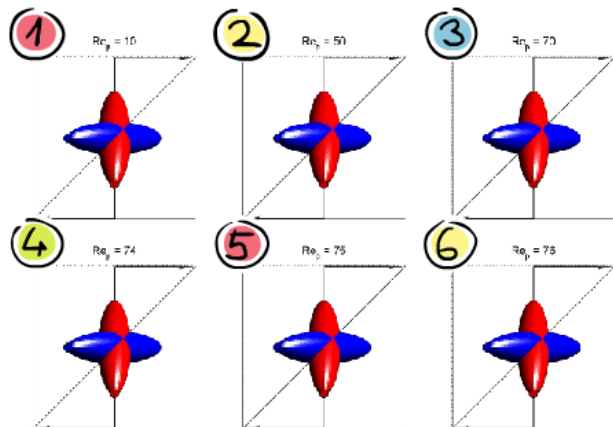
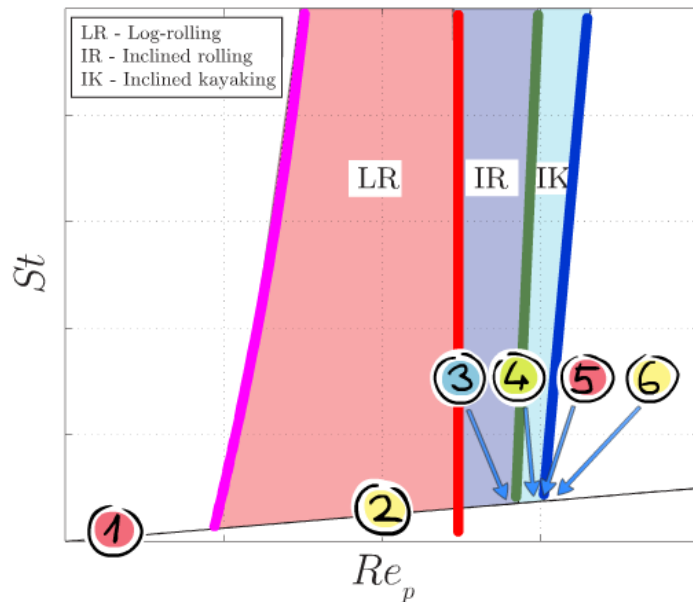


Stokes number = Particle inertia/viscous damping  
 Reynolds number = Aligning torque/Rotating torque  
 = Local strain/Local vorticity



# Non-planar rotational states

(Non-planar = symmetry axis NOT perpendicular to vorticity)

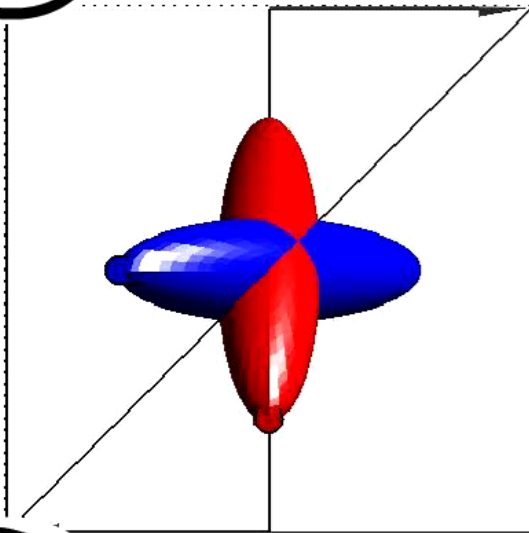


- Particle behaves like a Duffing-Van der Pol oscillator, probably due to the presence of a double zero eigenvalue.
- Reynolds number connected to nonlinear restoring/repelling forces and amplification/damping.
- Stokes number connected to nonlinear damping/amplification.

$P$

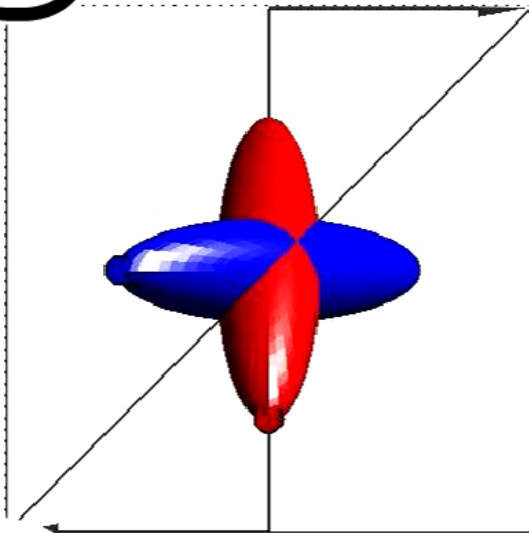
1

$Re_p = 10$



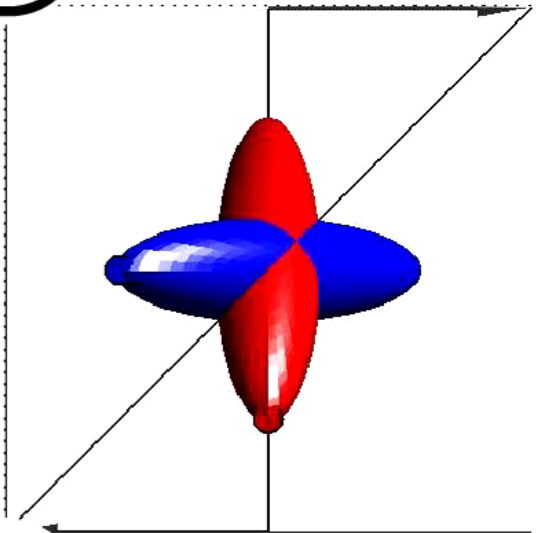
2

$Re_p = 50$



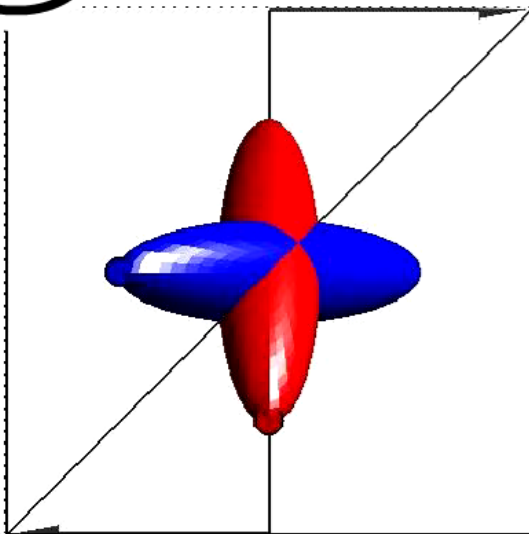
3

$Re_p = 70$



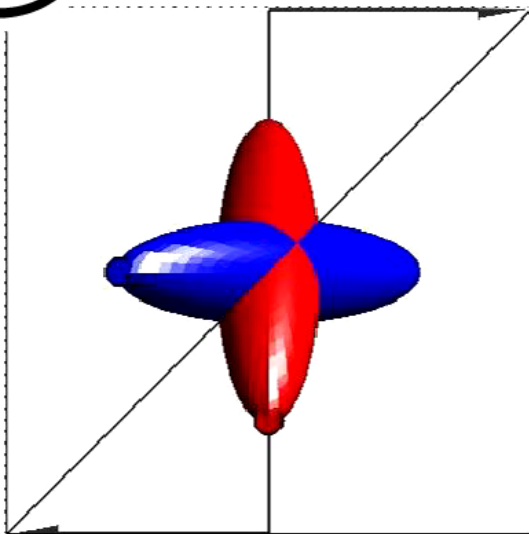
4

$Re_p = 74$



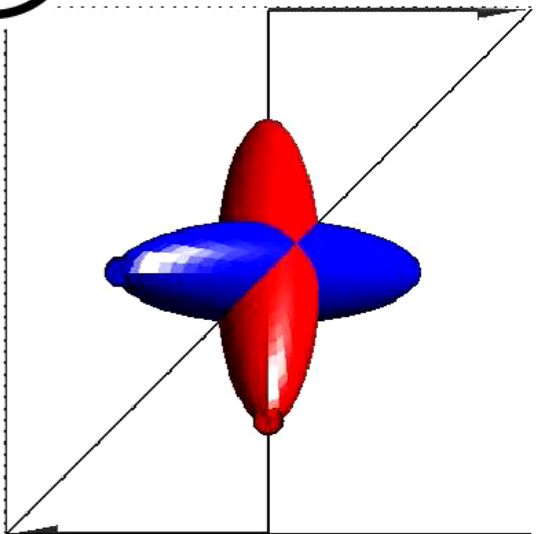
5

$Re_p = 75$



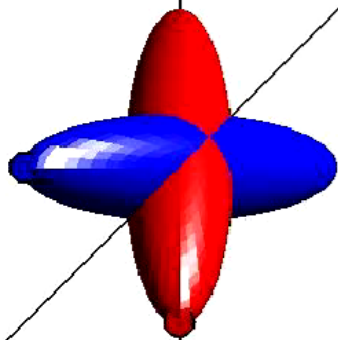
6

$Re_p = 76$



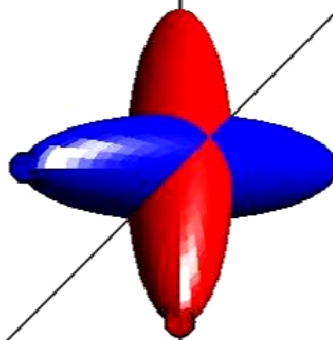
1

$Re_p = 10$



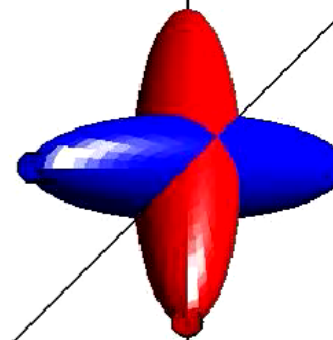
2

$Re_p = 50$



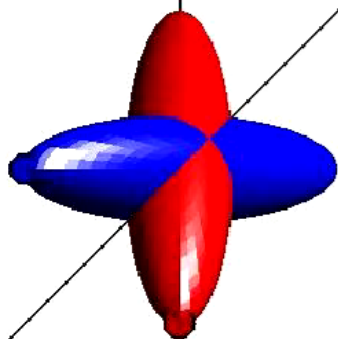
3

$Re_p = 70$



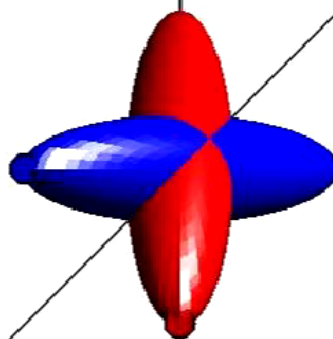
4

$Re_p = 74$



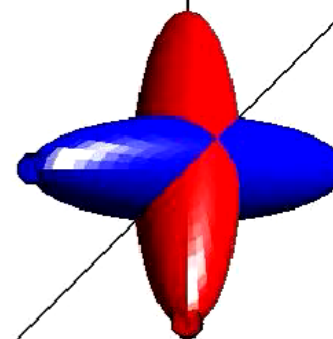
5

$Re_p = 75$



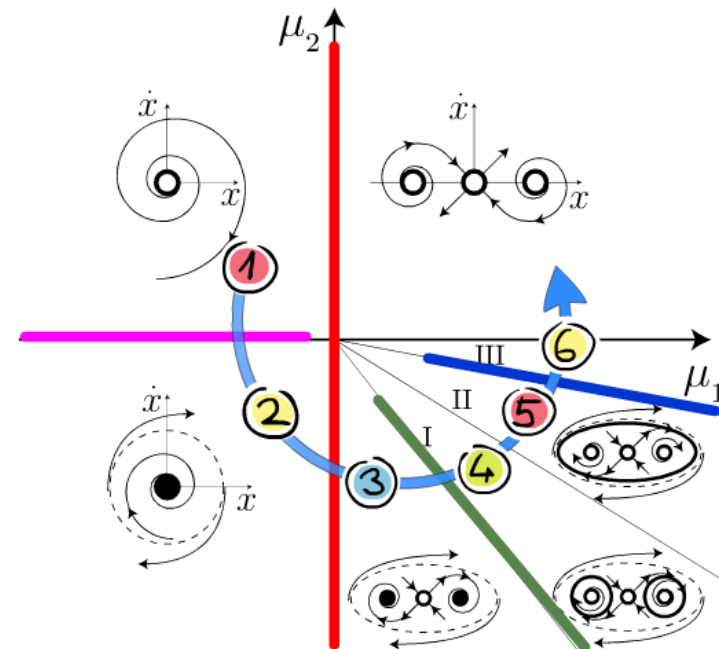
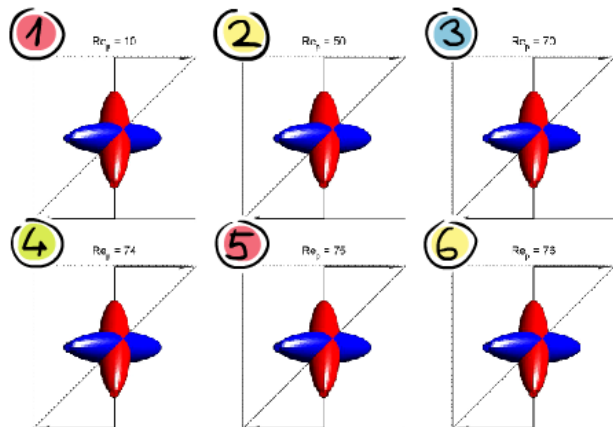
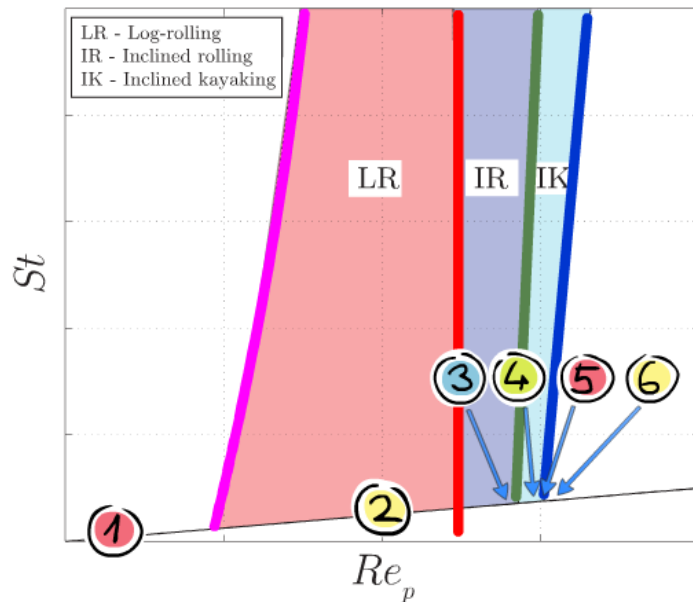
6

$Re_p = 76$



# Non-planar rotational states

(Non-planar = symmetry axis NOT perpendicular to vorticity)



- Particle behaves like a Duffing-Van der Pol oscillator, probably due to the presence of a double zero eigenvalue.
- Reynolds number connected to nonlinear restoring/repelling forces and amplification/damping.
- Stokes number connected to nonlinear damping/amplification.



# Conclusions

- Stable rotational states appear as a consequence of particle and fluid inertia.
- Particle behaves like a driven damped pendulum in the planar motion.
- Particle behaves like a Duffing-Van der Pol oscillator in the non-planar motion, due to the presence of a double zero eigenvalue in the Log-rolling motion.
- Particle inertia counteracts viscous damping and causes a drift towards a planar rotation (centrifugal forces).
- Fluid inertia causes a relative increase in the local strain. This causes pendulum-like behavior in the planar rotation and drift towards non-planar rotation.
- Fluid inertia also causes the nonlinear restoring/repelling, which is not explained by linear effects from local strain.

# Thank you!

