

Joint 5th SIG43–FP1005 Workshop on

Fiber Suspension Flow Modelling



Rotation Statistics of Fibers in Wall-Shear Turbulence

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Motivation of the Study: Why are we looking at these statistics?



Statistical characterization of fiber angular velocity is crucial to:

1. Quantify rotational dispersion and rotation rates

Parameterization of rotation rates as a function of the ratio:

fiber length Lagrangian integral flow scale

2. Validate rheological models

How does the extra non-Newtonian stress, produced by the addition of fibers to the Newtonian carrier fluid, depend on fiber orientation distribution?



Aim of the study: examine the effect of local shear and turbulence anisotropy on fiber rotation at varying fiber inertia/length



Methodology - Carrier Fluid



• Examples: flow of air at 1.8 m/s in a 4 cm high channel flow of water at 3.8 m/s in a 0.5 cm high channel



Methodology - Fibers



Fibers are modelled as **prolate ellipsoidal particles**.

Lagrangian particle tracking.

Simplifying assumptions: dilute flow, <u>one-way</u> <u>coupling</u>, Stokes flow ($Re_P < 1$), pointwise particles (particle size is smaller than the smallest flow scale).

Periodicity in x ed y, elastic rebound at the wall and **conservation of angular momentum**.

200,000 fibers tracked, *random* initial position and orientation, linear and angular velocities equal to those of the fluid at fiber's location.



Methodology – Fiber Kinematics



Kinematics: described by (1) position of the fiber center of mass and (2) fiber orientation.



• Euler parameters: e_0 , e_1 , e_2 , e_3

$$e_0 = \cos\left[\frac{1}{2}(\psi+\varphi)\right]\cos\left(\frac{\theta}{2}\right)$$
 ,...

• Rotation matrix: $\mathbf{x}' = R_{Eul}\mathbf{x}''$

$$R_{eul} = \begin{bmatrix} e_0^2 + e_1^2 - e_2^2 - e_3^2 & 2(e_1e_2 + e_0e_3) & 2(e_1e_3 - e_0e_2) \\ 2(e_1e_2 - e_0e_3) & e_0^2 - e_1^2 + e_2^2 - e_3^2 & 2(e_2e_3 + e_0e_1) \\ 2(e_1e_3 + e_0e_2) & 2(e_2e_3 - e_0e_1) & e_0^2 - e_1^2 - e_2^2 + e_3^2 \end{bmatrix}$$

- X_G, Y_G, Z_G
- 3 frames of reference (to define orientation)
- Euler angles: φ , ψ , θ (singularity problems)





Methodology – Fiber Dynamics



Rotational dynamics: Euler equations with Jeffery moments.

• Euler Equations: (2nd cardinal law)

$$I_{x'x'}\dot{\omega}_{x'} + \omega_{y'}\omega_{z'}(I_{z'z'} - I_{y'y'}) = M_{x'}^{est}$$
$$I_{y'y'}\dot{\omega}_{y'} + \omega_{x'}\omega_{z'}(I_{z'z'} - I_{x'x'}) = M_{y'}^{est}$$
$$I_{z'z'}\dot{\omega}_{z'} + \omega_{x'}\omega_{y'}(I_{y'y'} - I_{x'x'}) = M_{z'}^{est}$$

(in the particle frame)





Methodology – Fiber Dynamics



Translational Dynamics: hydrodynamic resistance (Brenner, 1963).

- First cardinal law: 0
 - Brenner's law: 0 (form drag and skin drag)
- $m_P \frac{d\mathbf{v}}{dt} = \sum_i \mathbf{F}_i = \mathbf{F}_{drag}$ (inertia and drag only!!)

 $\mathbf{F}'_{drag} = \mu \pi a \mathbf{\bar{k}'} (\mathbf{u'} - \mathbf{v'})$ (in the fiber frame) Resistance

Tensor

In the inertial (Eulerian) frame: 0

$$\mathbf{\underline{u}}' = R_{eul} \mathbf{\underline{u}}$$

$$\mathbf{\overline{K}}_{(\varphi,\theta,\psi)} = R_{eul}^{T} \mathbf{\overline{K}}' R_{eul}$$

$$\Rightarrow \mathbf{F}_{drag} = \mu \pi a \mathbf{\overline{K}}_{(\varphi,\theta,\psi)} (\mathbf{u} - \mathbf{v})$$

$$\left\{ \begin{array}{c} m_{P} \frac{d\mathbf{v}}{dt} = \mu \pi a \mathbf{\overline{K}}_{(\varphi,\theta,\psi)} (\mathbf{u} - \mathbf{v}) \\ \frac{d\mathbf{x}_{(G)}}{dt} = \mathbf{v} \end{array} \right\} \mathbf{v}(t) \mathbf{x}_{G}(t) \quad \text{(via numerical integration)}$$

$$Once \text{ fiber orientation is known, fiber translational motion can be computed!}$$

IULIUII Call







The physics of turbulent fiber dispersion is determined by a small set of parameters

- Aspect ratio: $\lambda = \frac{b}{a}$ (chosen values: $\lambda = 1.001, 3, 10, 50$)
- Stokes number: $St = \tau^+ = \frac{\tau_P}{\tau_F}$ (chosen values: $\underline{\tau^+=1, 5, 30, 100}$) • $\tau^+>1$: large inertia ("stones") • $\tau^+<1$ • $\tau^+<1$ • $\tau^+\sim1$

Input parameters:⁺ S, τ , λ

• Specific density: $S = \frac{\rho_P}{\rho_F}$



"Cartoon" of fiber's elongation









Mean (space & time-averaged) angular velocity in the near-wall region







Mean (space & time-averaged) angular velocity in the near-wall region







Root Mean Square of fiber angular velocities in the near-wall region



Strong effect of fiber inertia near the wall, length-related effects weaken!





Root Mean Square of fiber angular velocities in the near-wall region



Strong effect of fiber inertia near the wall, length-related effects weaken!





Root Mean Square of fiber angular velocities in the near-wall region







Root Mean Square of fiber angular velocities in the near-wall region







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Root Mean Square of fiber angular velocities in the near-wall region









Can fiber rotation be described within the theory of diffusion as a Ornstein-Uhlenbeck (OU) process?

The OU process is completely characterized by:

- $\boldsymbol{\cdot} \textbf{ statistically-stationary Gaussian distribution}$
- exponentially decaying autocovariance

$$R(\tau) = \alpha^2 \exp(-\tau/T_*)$$

where:

- $lpha^2$ variance of the Gaussian distribution
- T_* integral timescale of the process

We aim at assessing the applicability of such description in wall-shear turbulence







We compute Lagrangian autocorrelations for fibers with different inertia and elongation conditioning the statistics to 3 specific regions of the flow

$$R_{\Omega_i,\Omega_i}(\tau) = \frac{\langle \Omega_i'(\mathbf{x}_p(t_0), t_0) \Omega_i'(\mathbf{x}_p(t_0+\tau), t_0+\tau) \rangle}{\langle \Omega_i'(\mathbf{x}_p(t_0), t_0)^2 \rangle^{1/2} \langle \Omega_i'(\mathbf{x}_p(t_0+\tau), t_0+\tau)^2 \rangle^{1/2}}$$









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Conditioned Lagrangian autocorrelation curves in wall-shear turbulence

Log layer



In the log layer, autocorrelation curves display a neat exponential decay. This trend becomes less clear as the wall region is approached.





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Conditioned Lagrangian autocorrelation curves in wall-shear turbulence



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Conditioned Lagrangian autocorrelation curves in wall-shear turbulence



In the log layer, autocorrelation curves display a neat exponential decay. This trend becomes less clear as the wall region is approached.







$$\Gamma_{\Omega}(\tau) = \int_0^{\tau} R_{\Omega_i,\Omega_i}(t) \mathrm{d}t$$







$$\Gamma_{\Omega}(\tau) = \int_0^{\tau} R_{\Omega_i,\Omega_i}(t) dt \implies R_{\Omega_i,\Omega_i}(\tau) = e^{-\frac{\tau}{T_L^{ii}}}$$





$$\Gamma_{\Omega}(\tau) = \int_{0}^{\tau} R_{\Omega_{i},\Omega_{i}}(t) dt \longrightarrow R_{\Omega_{i},\Omega_{i}}(\tau) = e^{-\frac{\tau}{T_{L}^{ii}}} \longrightarrow \Gamma_{\Omega_{i}}(\tau) = T_{L}^{ii} \left(1 - e^{-\frac{\tau}{T_{L}^{ii}}}\right)$$





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Conclusions ...





Fiber angular velocity is a useful measure of fibers-turbulence interaction in wall-bounded flows: its statistical characterization provides useful indications for modeling turbulent fiber dispersion

Angular velocity statistics depend both on fiber elongation (quantitatively) and fiber inertia (also qualitatively!)

Fiber rotation exhibits autocorrelation curves for the angular velocities that decay exponentially in the log layer (but not in the wall layer...)

