

Internal and surface waves in turbulence



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[c] University of Rochester - Department of Mechanical Engineering

1. Stably-stratified turbulence: internal waves
2. Viscosity-stratified turbulence: interfacial waves

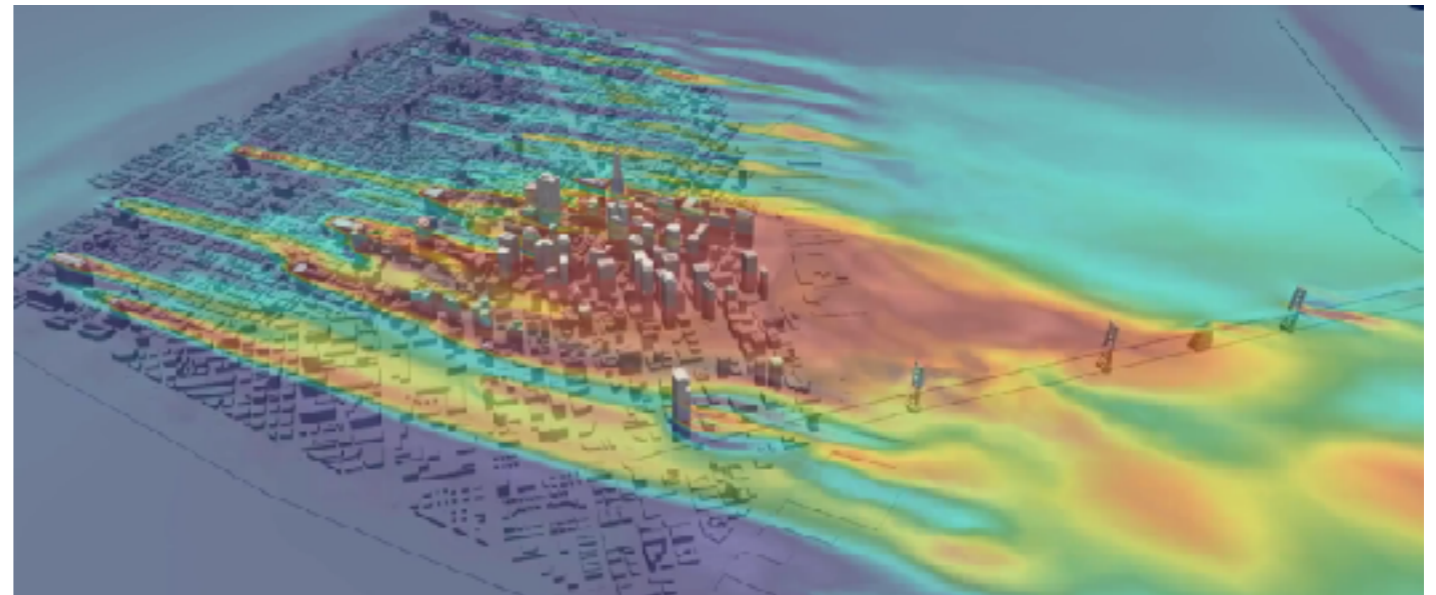
(Thermally) Stably-stratified wall-bounded turbulence



Algae/phytoplankton surfacing in lakes and oceans (which are usually stably stratified).

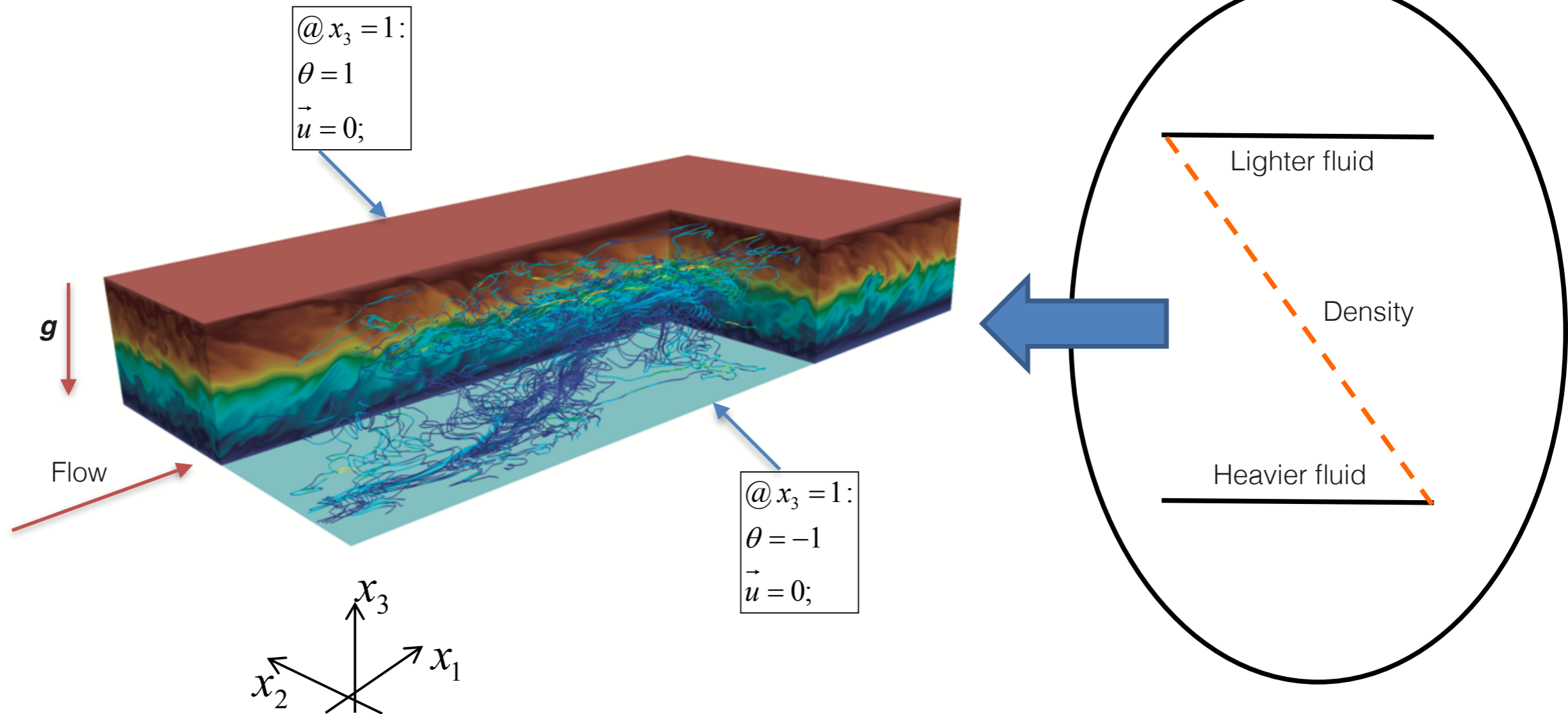
Algae blooms
(lake Erie, US)

Dispersion of pollutants in winter
(when the atmospheric boundary layer
is often stably-stratified)



Aim of the work: study the interaction between *turbulence* and *stable stratification* in a wall-bounded flow, at “large” Re;

We take the simplest possible configuration:



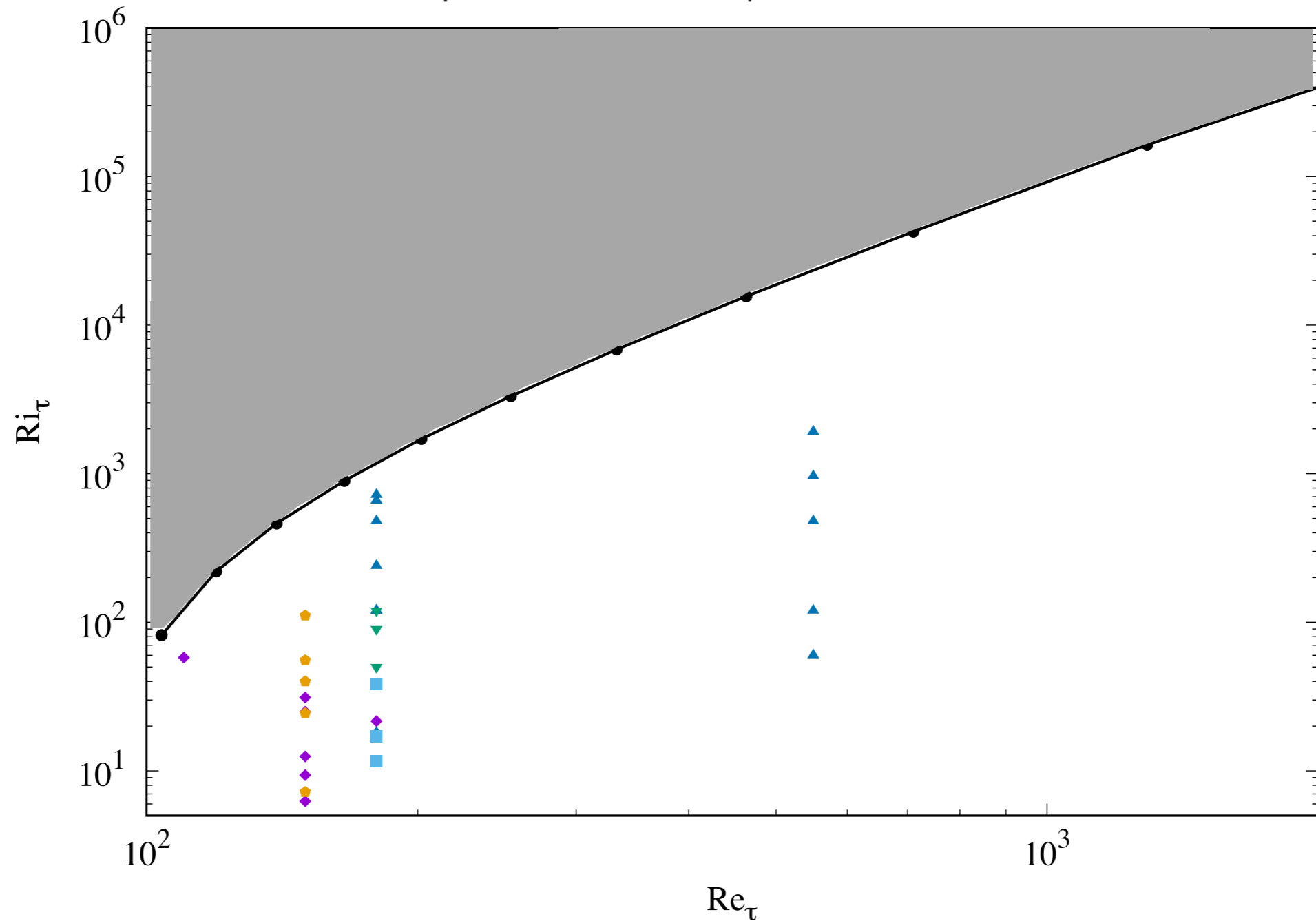
$$Re_\tau = \frac{u_\tau h}{\nu_0}$$

turbulence importance

$$Ri_\tau = \frac{g \Delta \rho h}{\rho_0 u_\tau^2}$$

stratification importance

$$Pr = \frac{\nu_0}{\kappa_0}$$



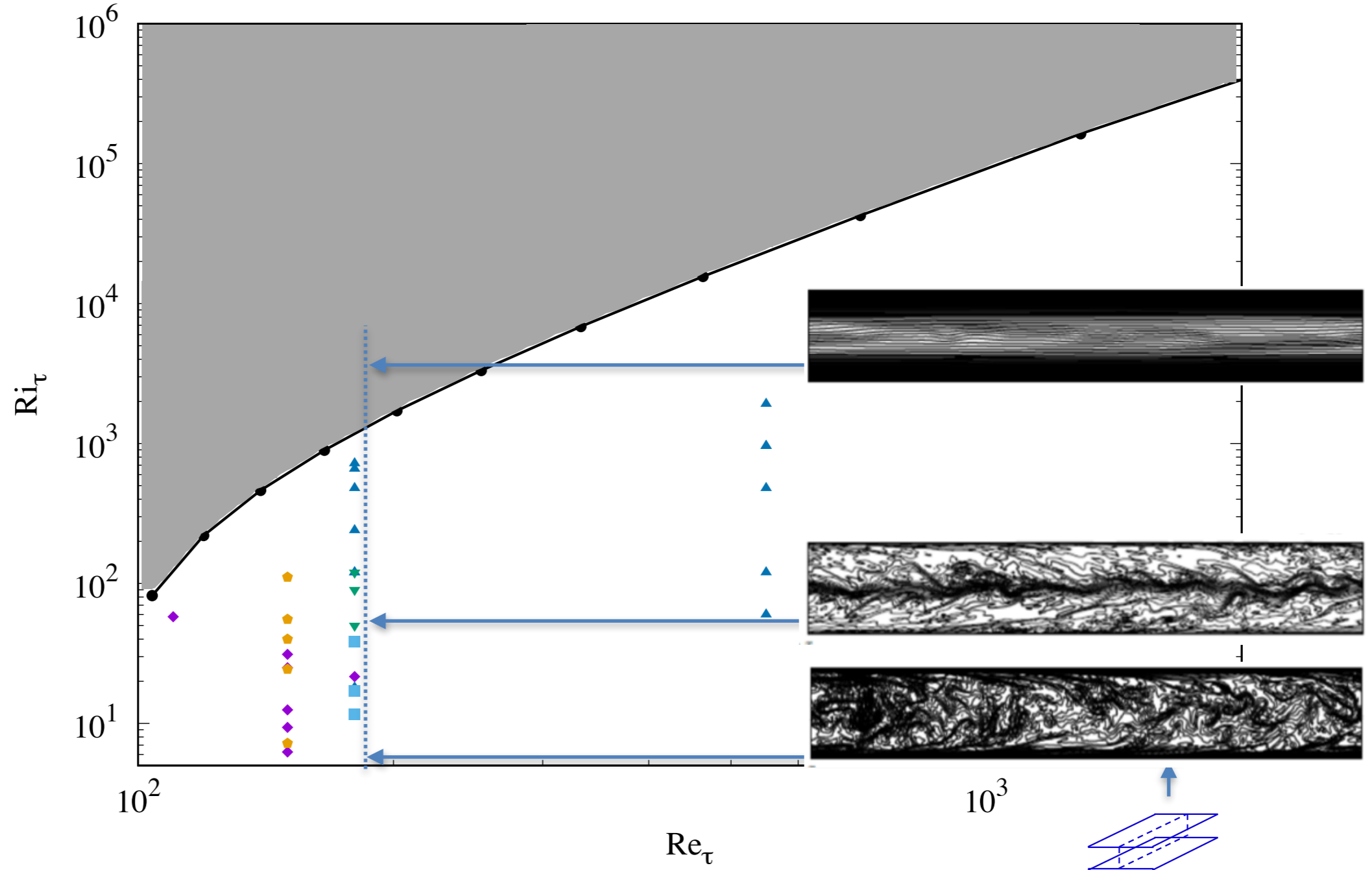
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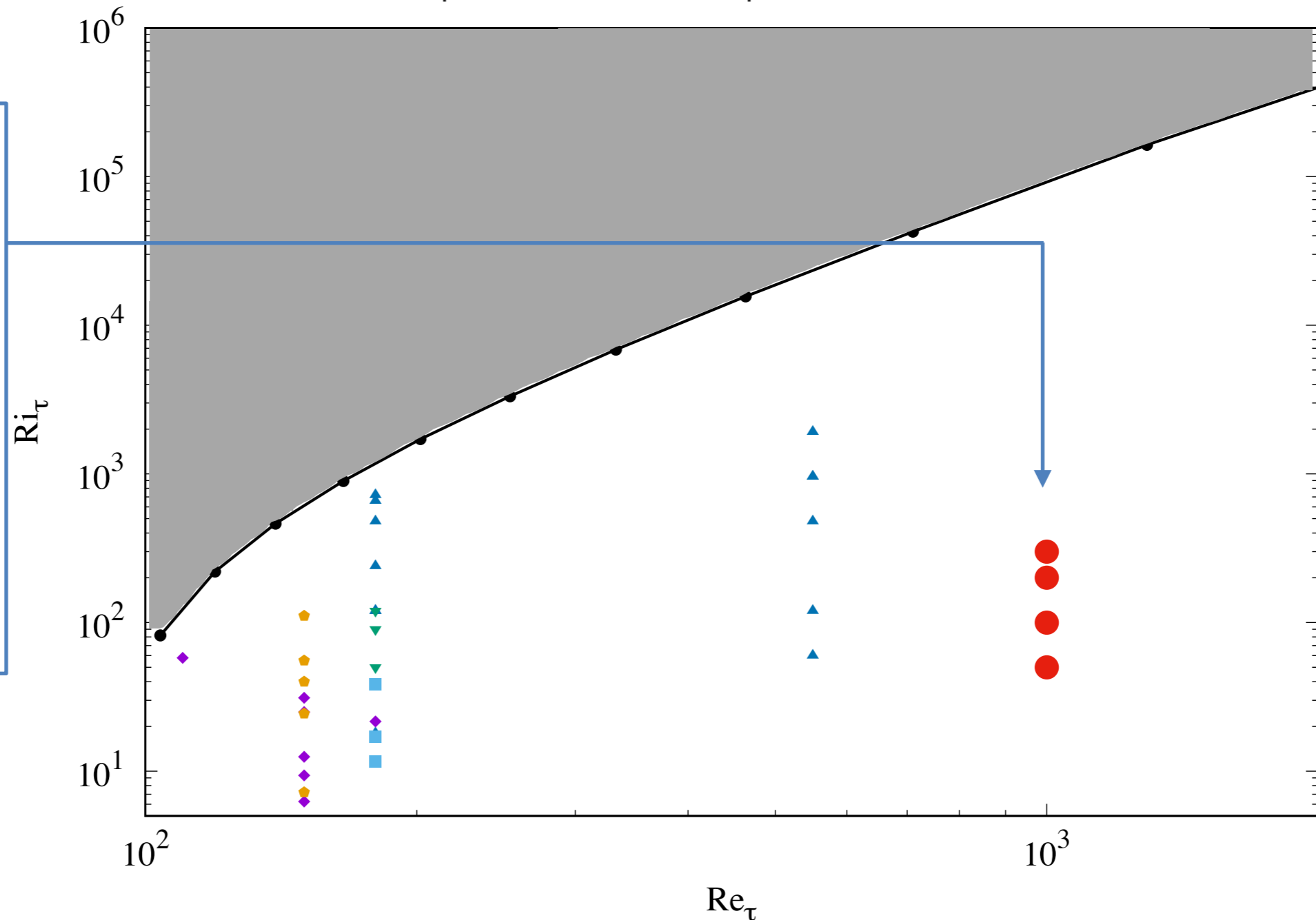
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stratification importance

$$Pr = \frac{\nu_0}{\kappa_0}$$



- No study in closed Poiseuille channel at $Re_t > 550$
- We want to move to $Re_t = 1000$, moderate stratification
- DNS data crucial for turbulence modeling (Pr_t , Ri_g , mixing efficiency... Previous talk!)

$$\frac{\partial u_i}{\partial x_i} = 0$$

$$\frac{Du_i}{Dt} = -\frac{\partial p'}{\partial x_i} + \frac{1}{Re_\tau} \frac{\partial^2 u_i}{\partial x_j^2} - Ri_\tau \theta \delta_{3,i} + \delta_{1,i}$$

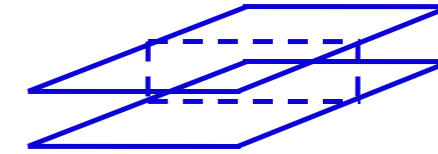
$$\frac{D\theta}{Dt} = \frac{1}{Re_\tau Pr} \frac{\partial^2 \theta}{\partial x_j^2}$$

- Direct Numerical Simulation, channel turbulence with air
- Imposed pressure gradient
- Pseudo-Spectral Fourier-Chebyshev method

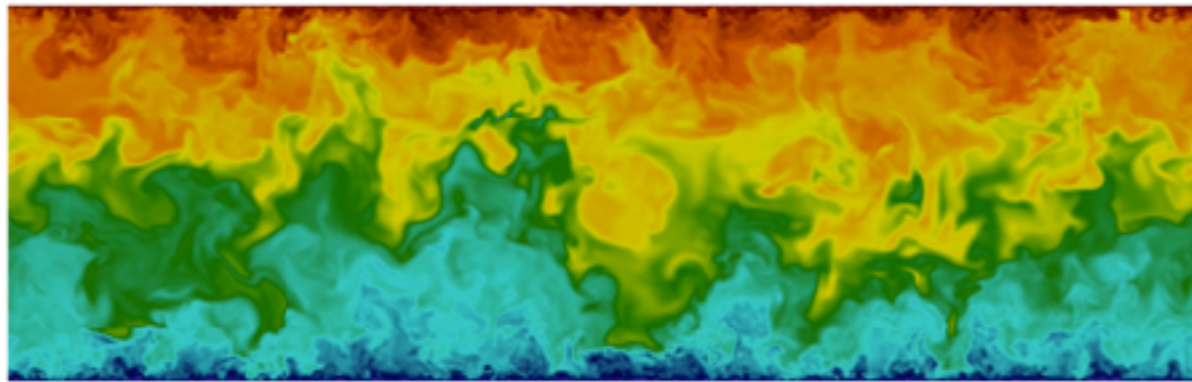
Simulation	Pr	Re_τ	Ri_τ	Grid
S0	0.71	1000	0	$1024 \times 1024 \times 1025$
S50	0.71	1000	50	$1024 \times 1024 \times 1025$
S100	0.71	1000	100	$1024 \times 1024 \times 1025$
S200	0.71	1000	200	$1024 \times 1024 \times 1025$
S300	0.71	1000	300	$1024 \times 1024 \times 1025$

$$Re_\tau = \frac{u_\tau h}{\nu_0} = 1000, \quad Pr = \frac{\nu_0}{\kappa_0} = 0.71, \quad Ri_\tau = \frac{g\beta_0 \Delta\theta h}{u_\tau^2}.$$

Visualization of temperature on a cross section



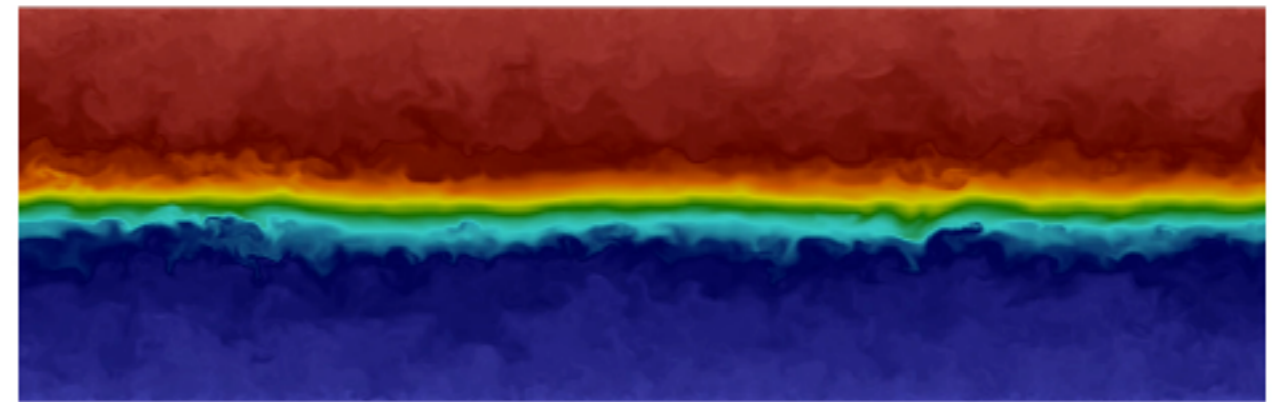
$$Ri_{\tau} = 0$$



Neutrally buoyant

- Higher mixing
- Mixing is driven by turbulence structures

$$Ri_{\tau} = 200$$

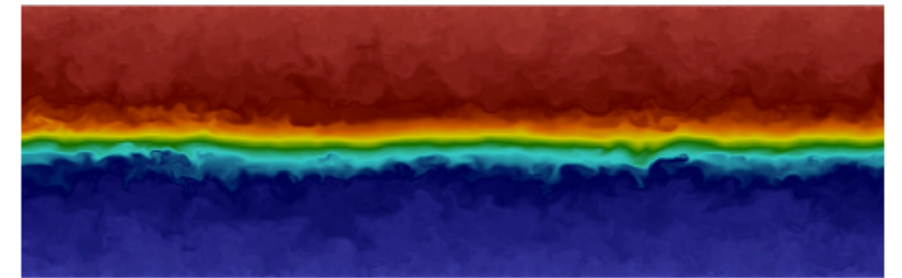
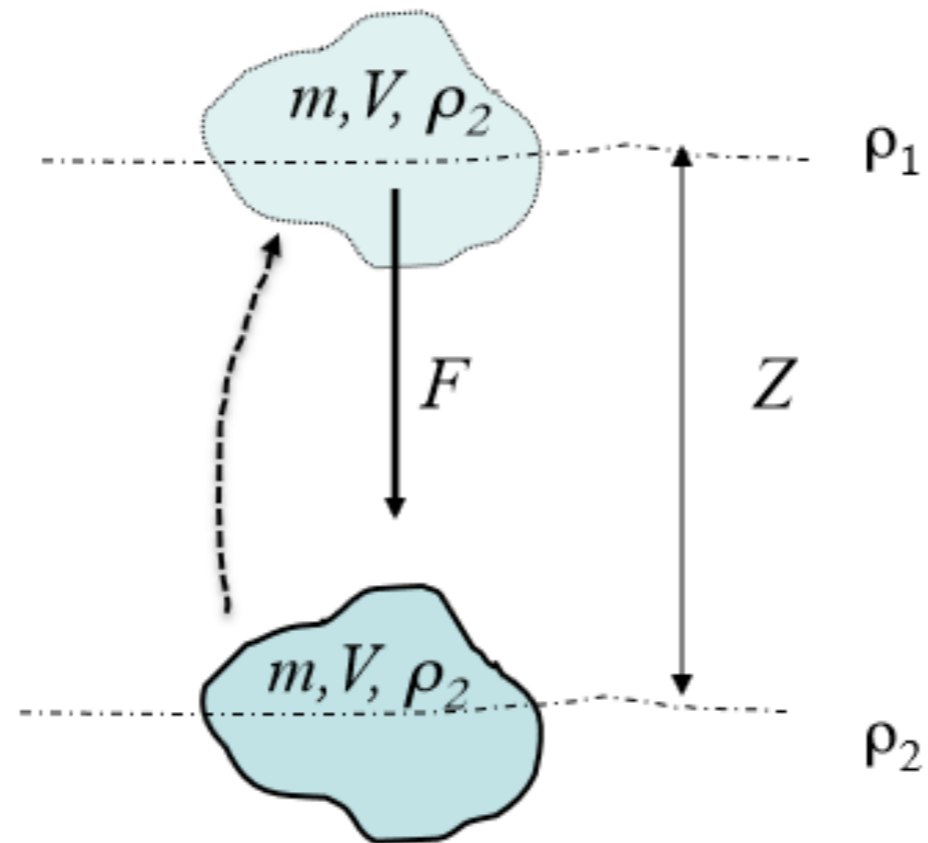
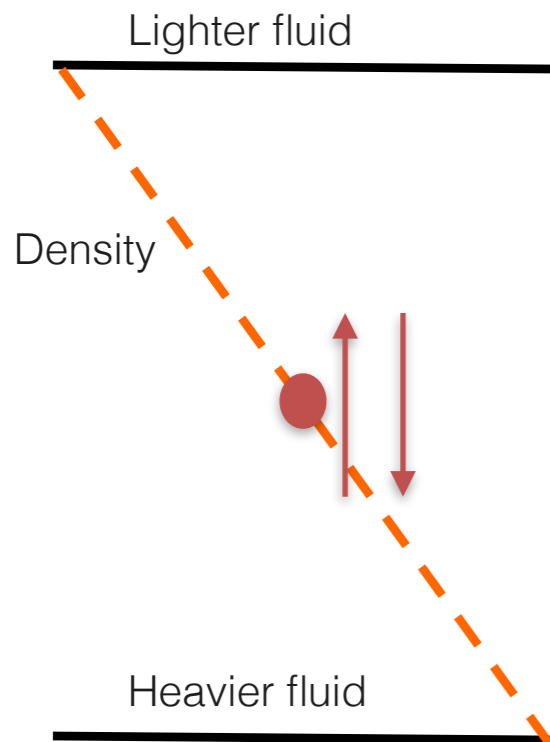


Stably stratified

- Lower mixing
- Internal waves act like barriers

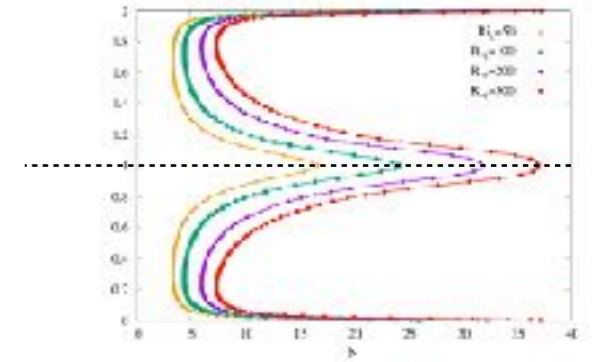
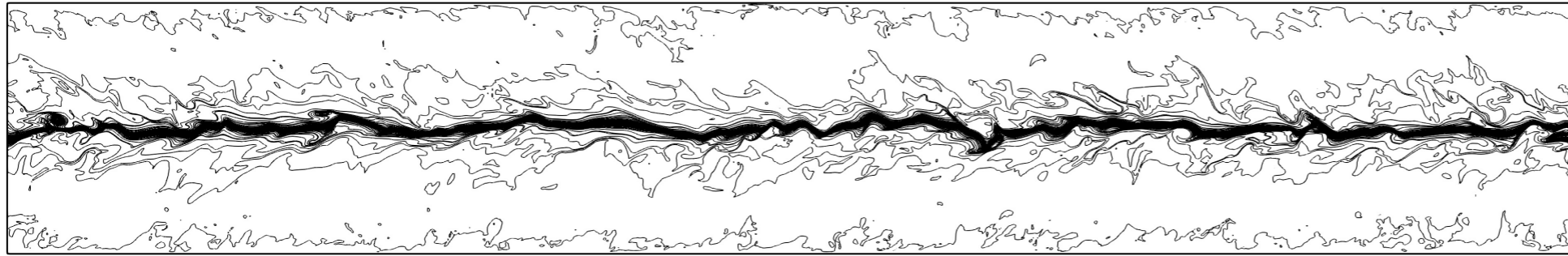
The origin of Internal Waves:

Force balance on a fluid particle



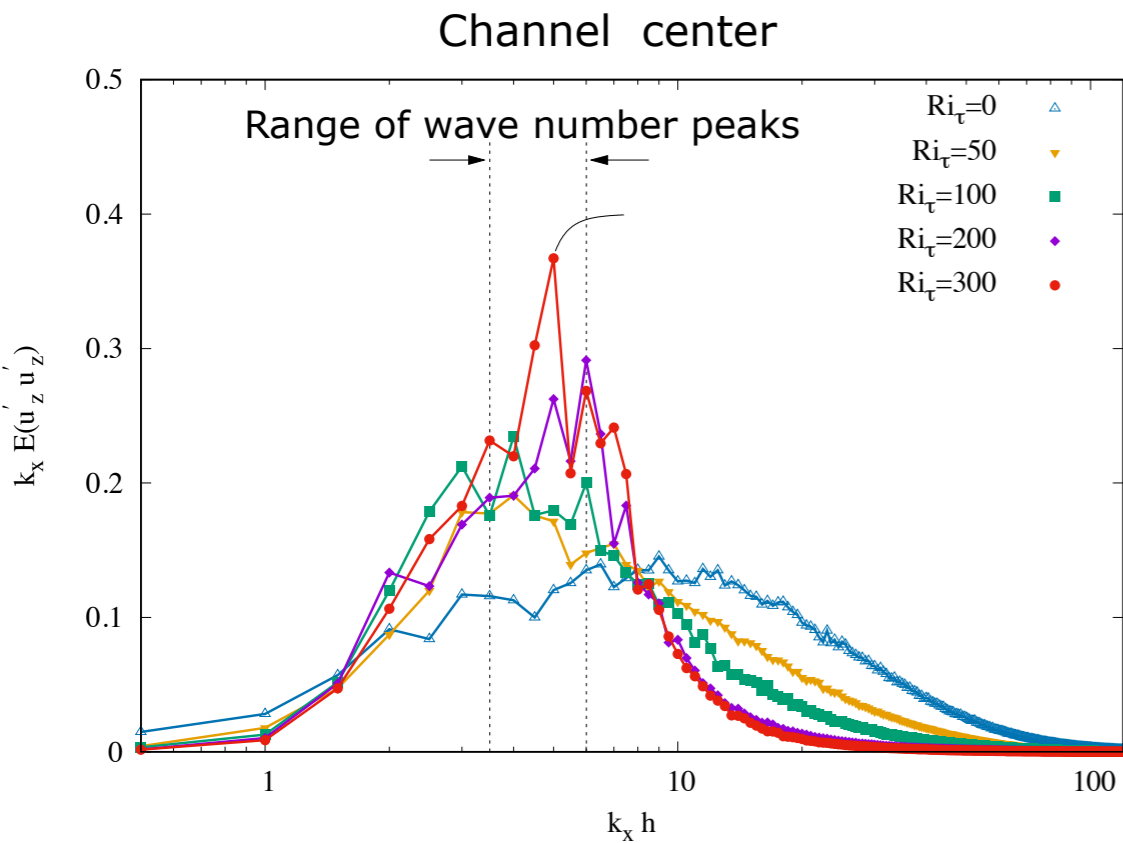
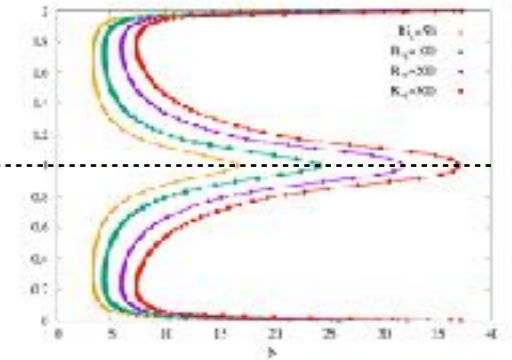
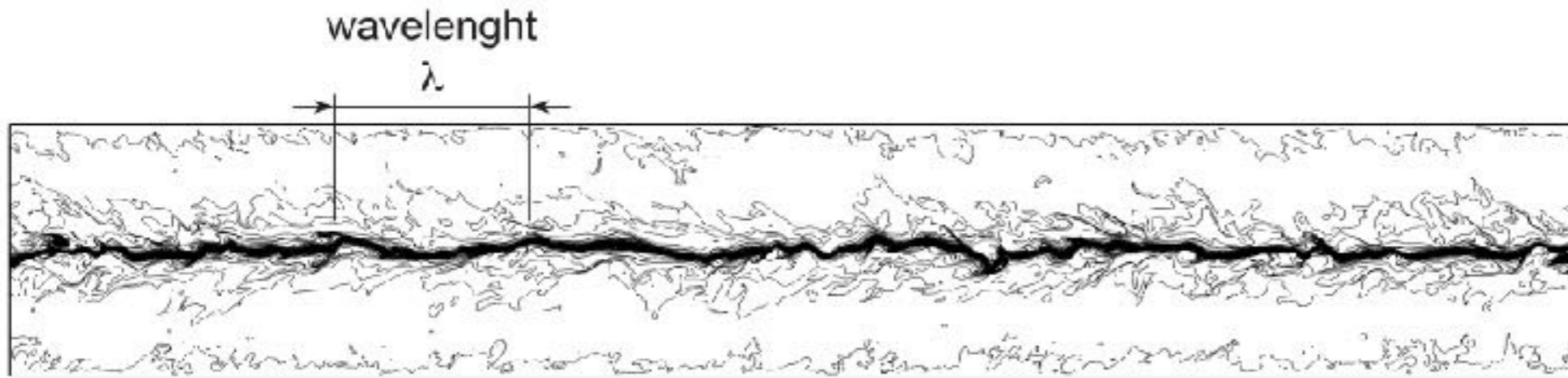
$$\left\{ \begin{aligned} \frac{\partial^2 Z}{\partial t^2} &= \frac{g}{\rho} \frac{\partial \rho}{\partial z} Z = N^2 Z \\ N &= \left(\frac{g}{\rho} \frac{\partial \rho}{\partial z} \right)^{1/2} \end{aligned} \right.$$

$$\left[N = \left(-g\beta \frac{\partial \langle \theta \rangle}{\partial z} \right)^{1/2} \right]$$



$$N = \sqrt{-g\beta \frac{\partial \langle \theta \rangle}{\partial z}}$$

$$N^* = Nu_\tau^* / h^*$$



Energy spectral density of the wall-normal velocity fluctuations

$$\langle u_{max} \rangle \longrightarrow \langle u_{wave} \rangle$$

$$u_{wave}^* = u_{wave} \times u_{\tau}^*$$

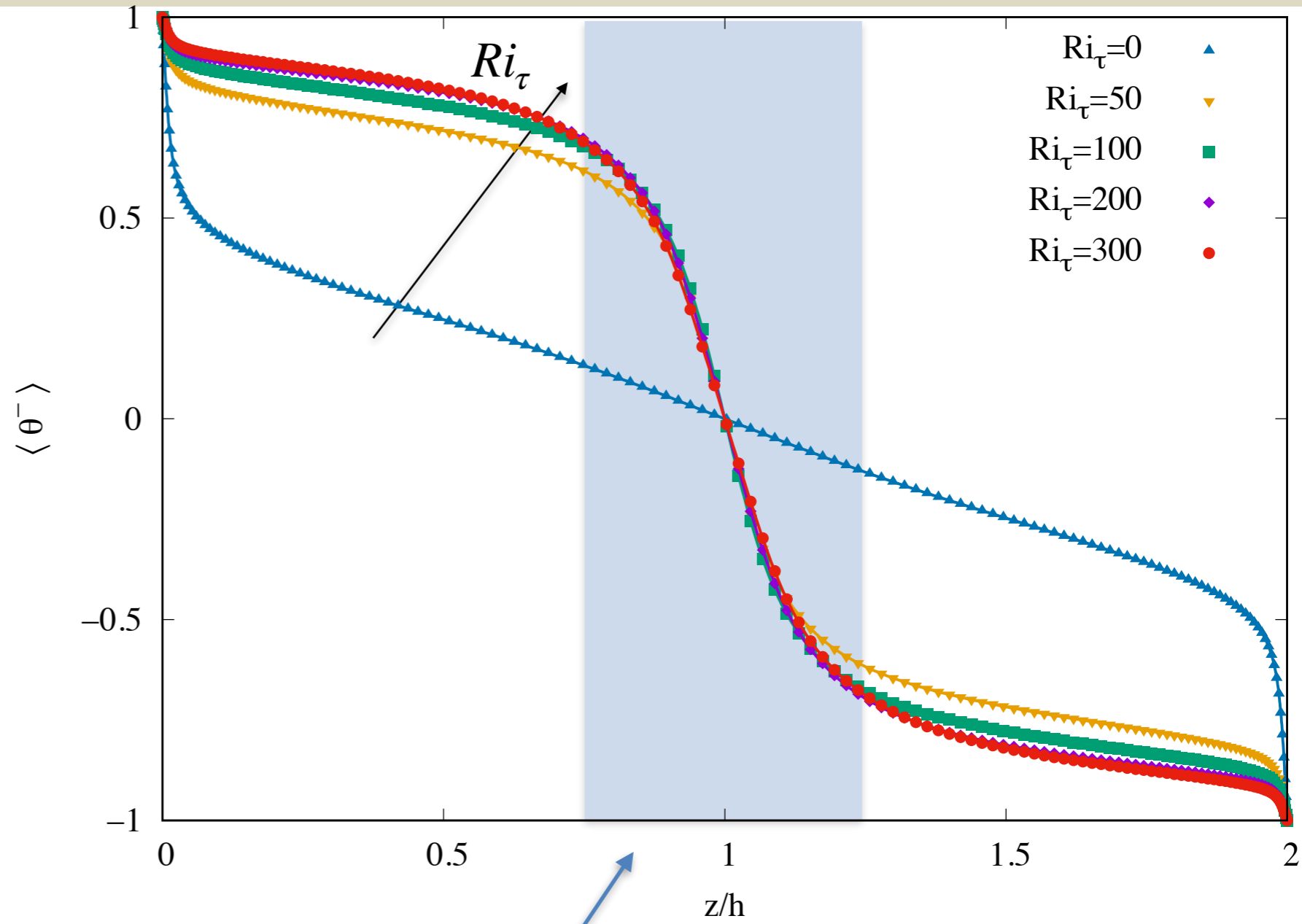
$$f_{wave}^* = u_{wave}^* \times k_x^* / 2\pi$$

$$N^* \sim f_{wave}^*$$

$$N = \sqrt{-g\beta \frac{\partial \langle \theta \rangle}{\partial z}}$$

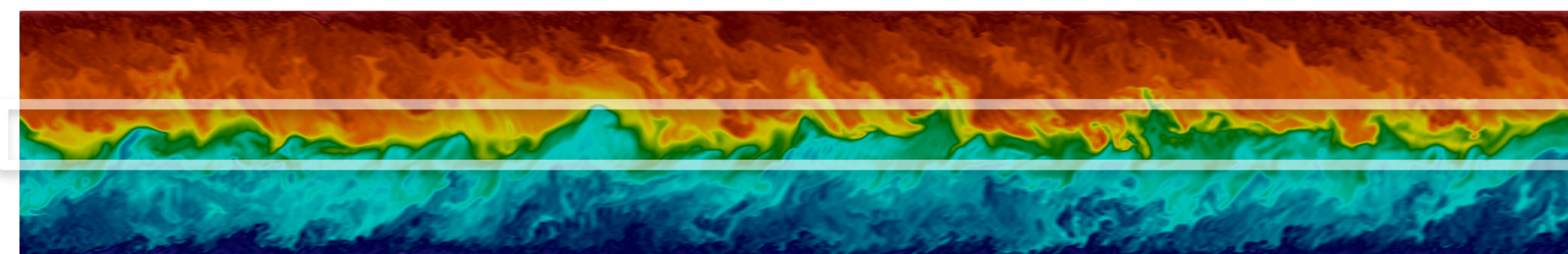
$$N^* = Nu_{\tau}^* / h^*$$

This links the Brunt-Väisälä frequency with the peak of the spectral energy

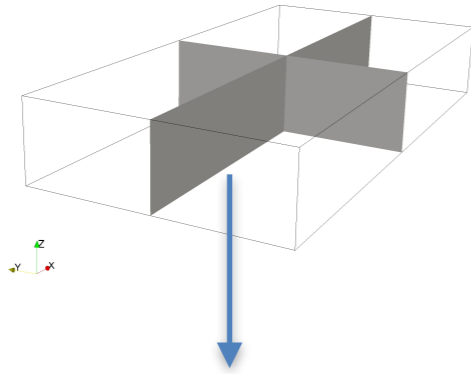


Main observations:

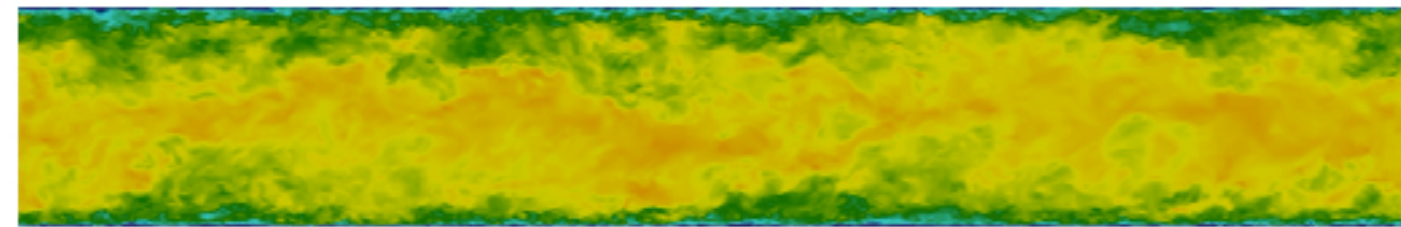
- IGWs (thermocline) behave like a thick interface
- Fluid parcels must climb a potential barrier
- Wall normal transport of heat is reduced



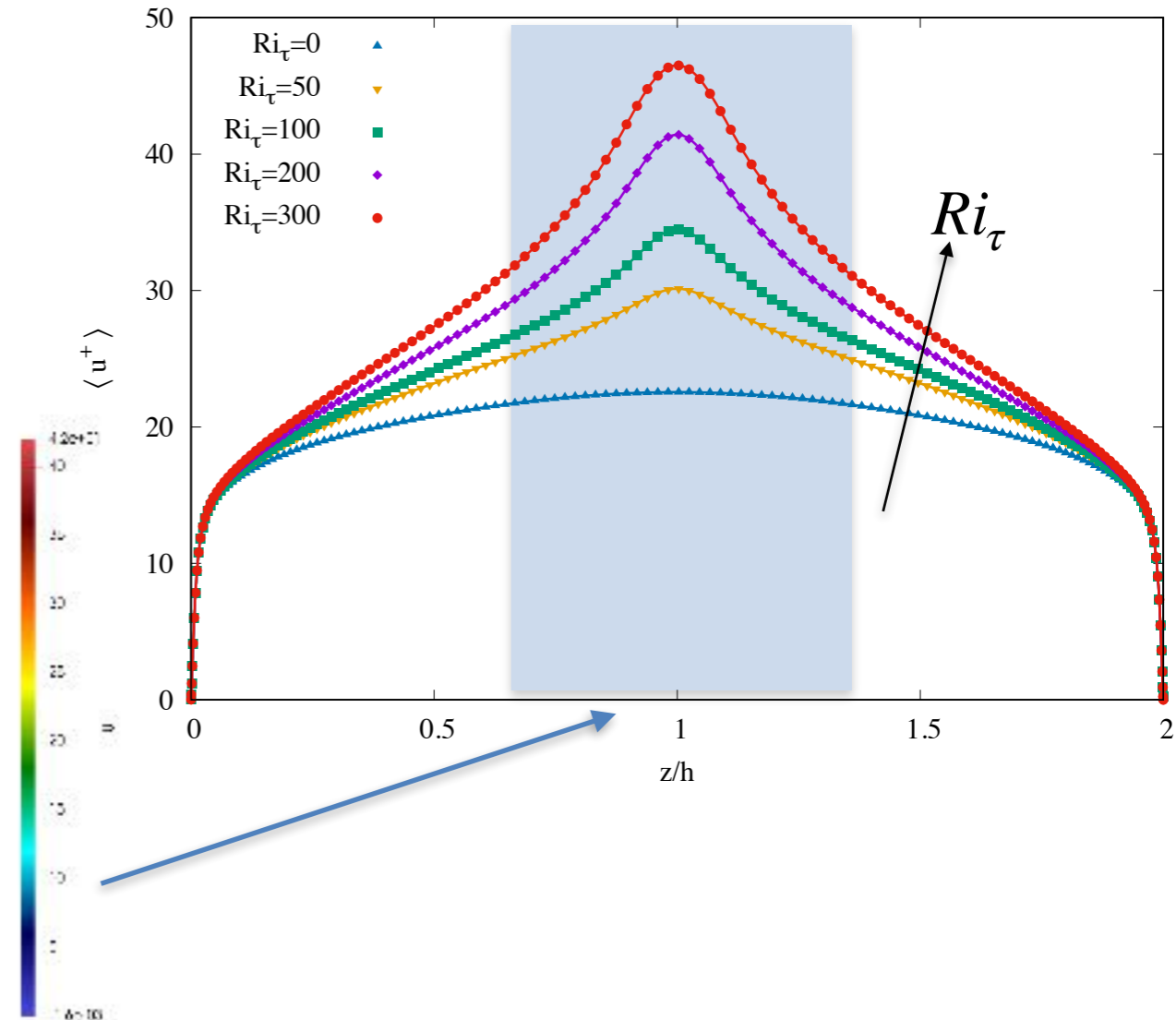
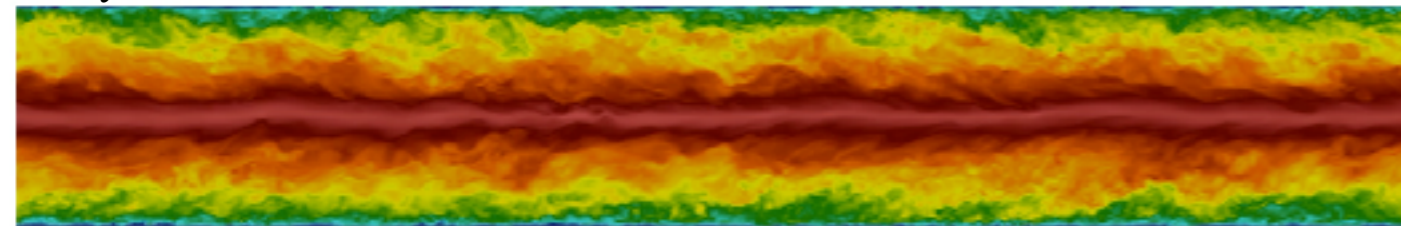
The temperature distribution has a remarkable effect on the velocity field



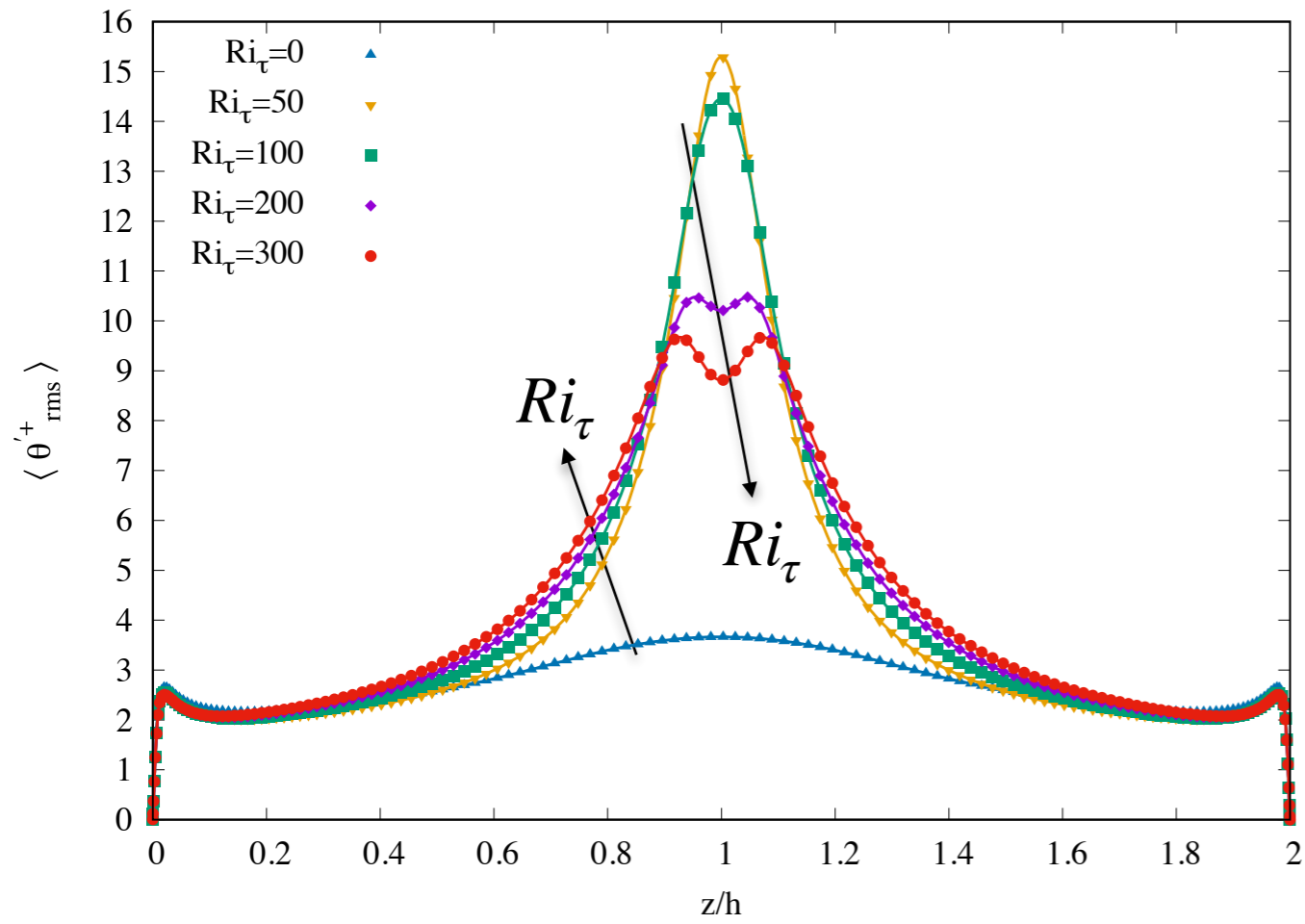
$Ri_\tau = 0$



$Ri_\tau = 200$

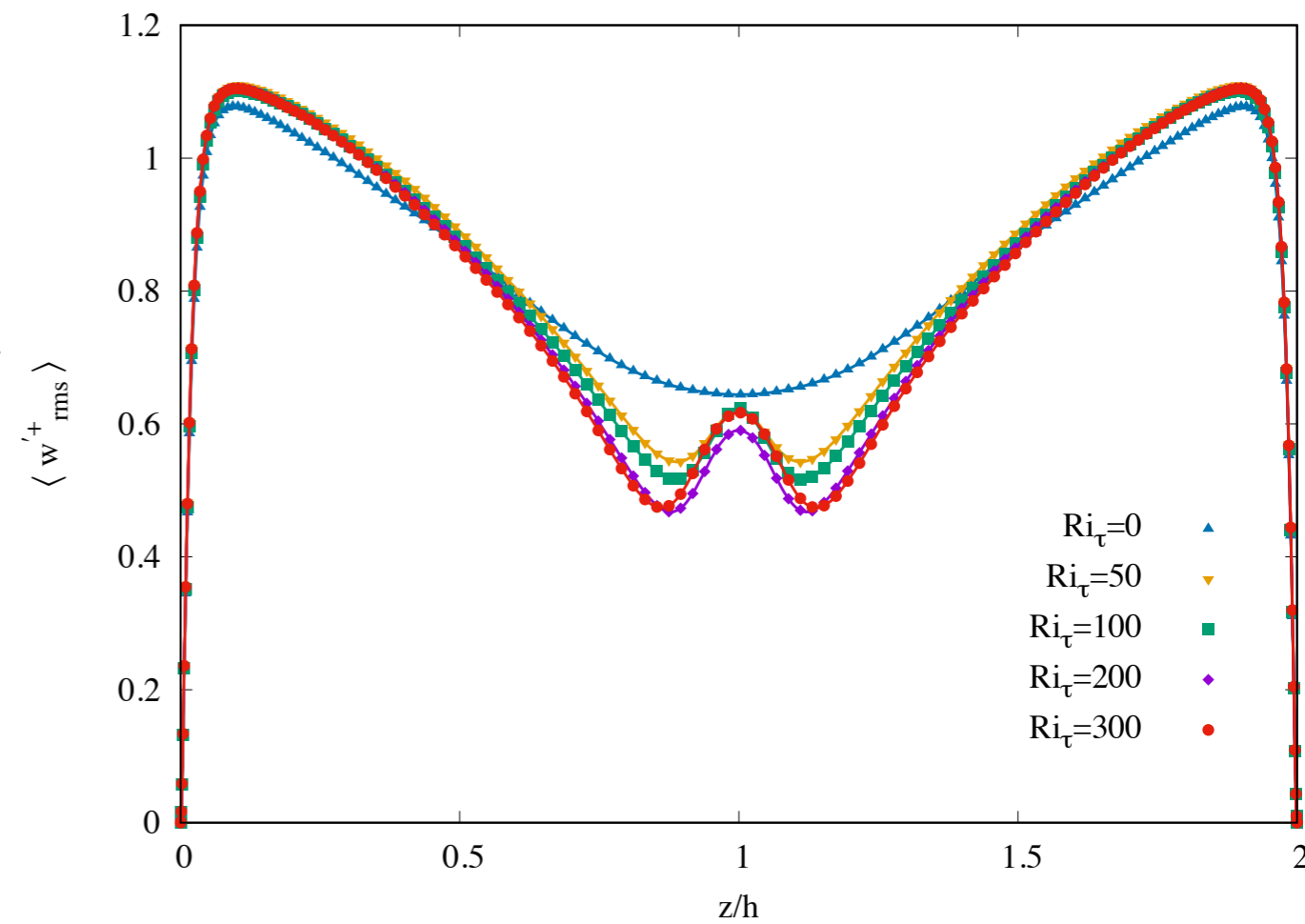


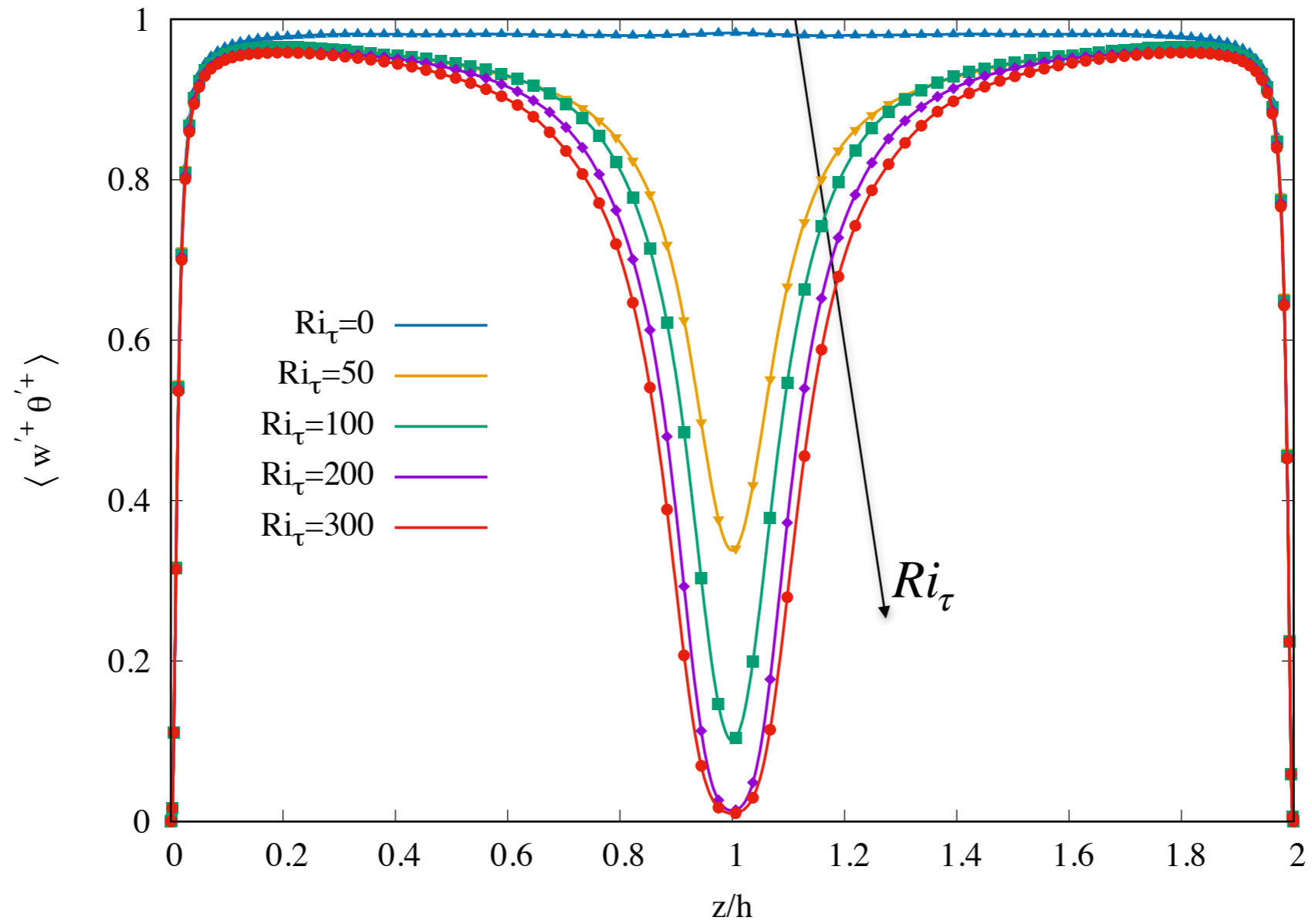
- Wall-normal transport of momentum is reduced (reduced friction factor)
- Stratified flow is “faster” and less homogeneous (layering)



Non-monotonic behavior

Large fluctuations at the channel center





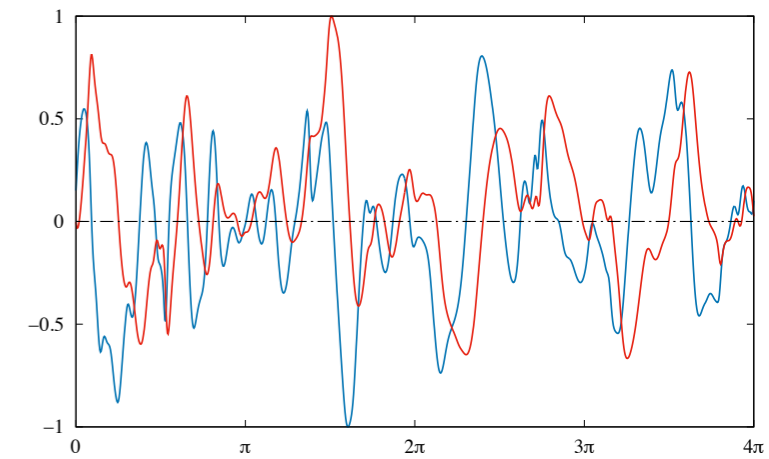
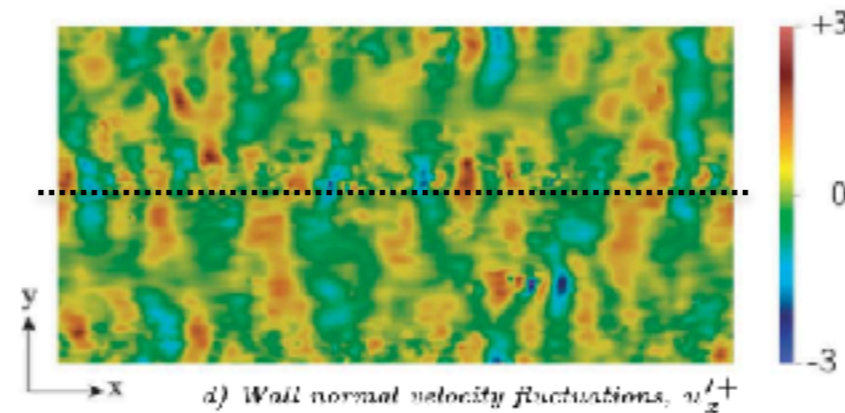
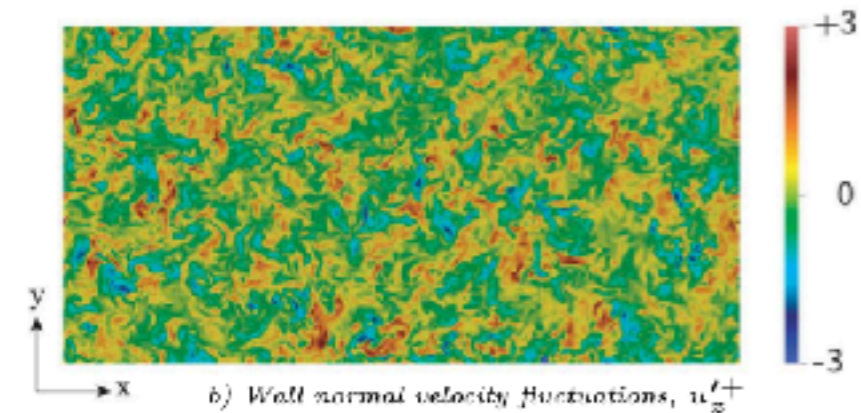
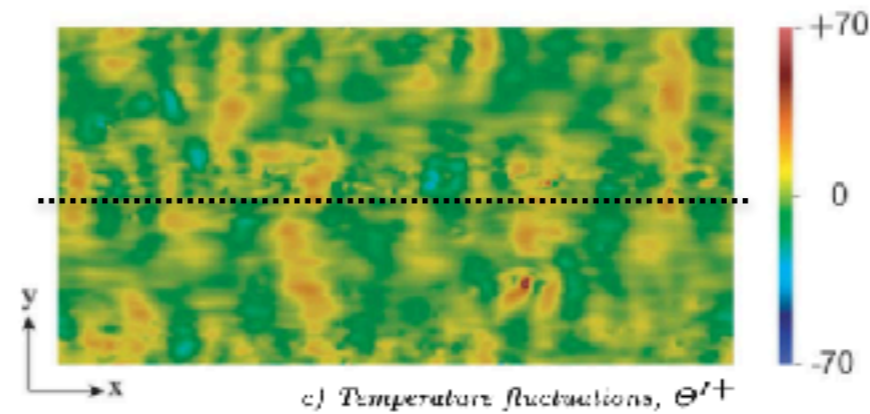
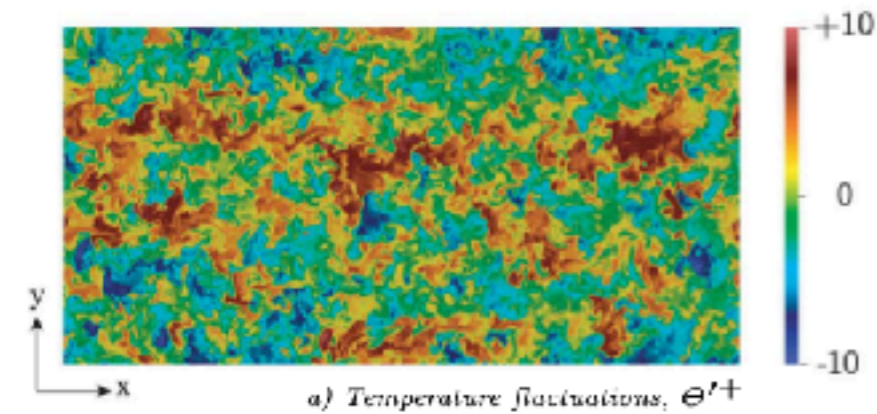
Main observation:

- Although fluctuations are large at the channel center, buoyancy flux is small...why?

Temperature and vertical velocity fluctuations in a horizontal plane parallel to the wall and located at the channel center, and corresponding cross spectrum

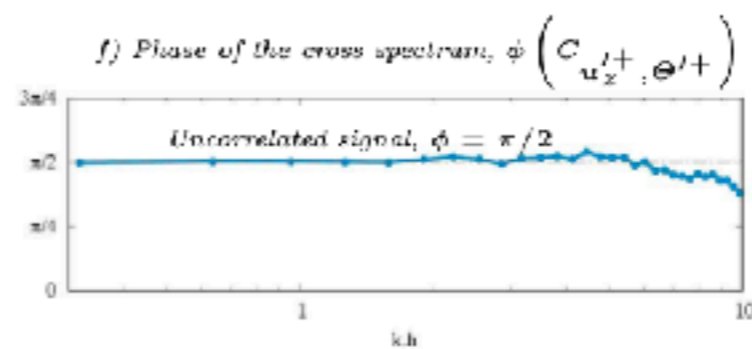
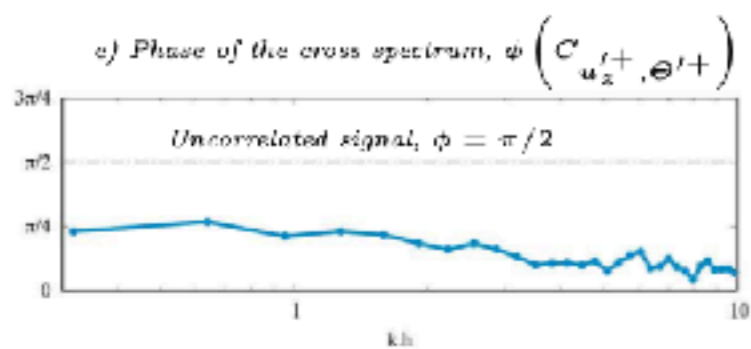
Neutrally buoyant, $Ri_\tau = 0$

Stratified, $Ri_\tau = 300$

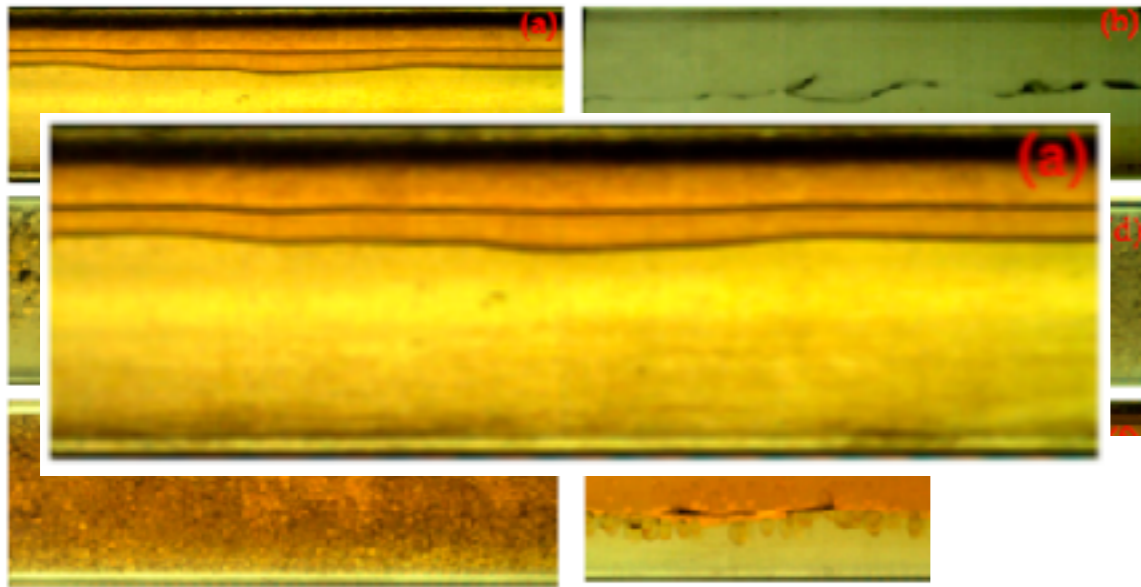


w'/w'_{max} —————

θ'/θ'_{max} —————



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2. Viscosity-stratified turbulence: interfacial waves

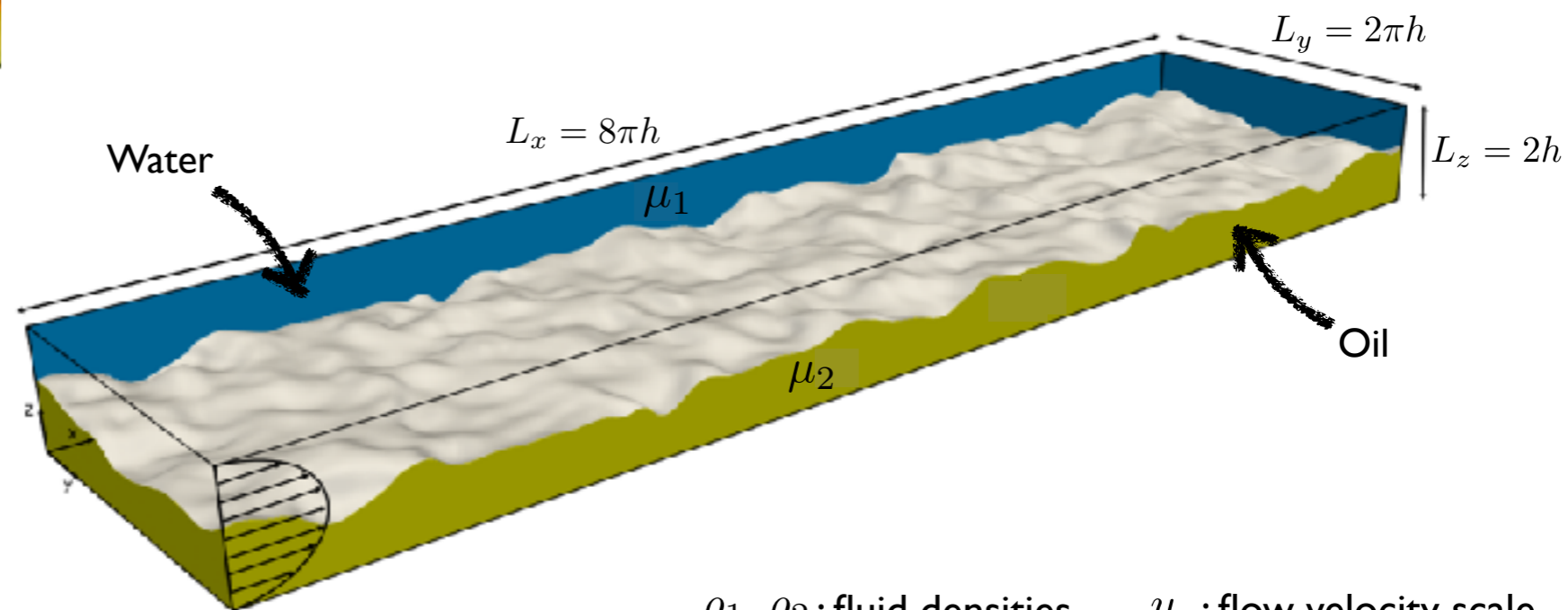


Phenomenology of oil-water flow in pipes and channels [1]

Reynolds number: $Re_\tau = \rho_2 u_\tau h / \mu_2$

Weber number: $We = \rho_2 u_\tau^2 h / \sigma$

Viscosity ratio: $\mu_r = \mu_1 / \mu_2$



ρ_1, ρ_2 : fluid densities u_τ : flow velocity scale
 μ_1, μ_2 : fluid viscosities σ : surface tension

Assumption: Matched densities $\rightarrow \rho_1 = \rho_2$

[1] Li, L.; Kong, L.; Xie, B.; Fang, X.; Kong, W.; Liu, X.; Wang, Y.; Zhao, F. The Influence on Response of a Combined Capacitance Sensor in Horizontal Oil–Water Two-Phase Flow. *Appl. Sci.* 2019, 9, 346.

Flow

Continuity: $\nabla \cdot \mathbf{u} = 0$

Navier-Stokes equations:

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \frac{1}{Re_\tau} \nabla \cdot [\mu(\phi, \mu_r)(\nabla \mathbf{u} + \nabla \mathbf{u}^T)] + \frac{3}{\sqrt{8}} \frac{Ch}{We} \nabla \cdot (|\nabla \phi|^2 \mathbf{I} - \nabla \phi \otimes \nabla \phi)$$

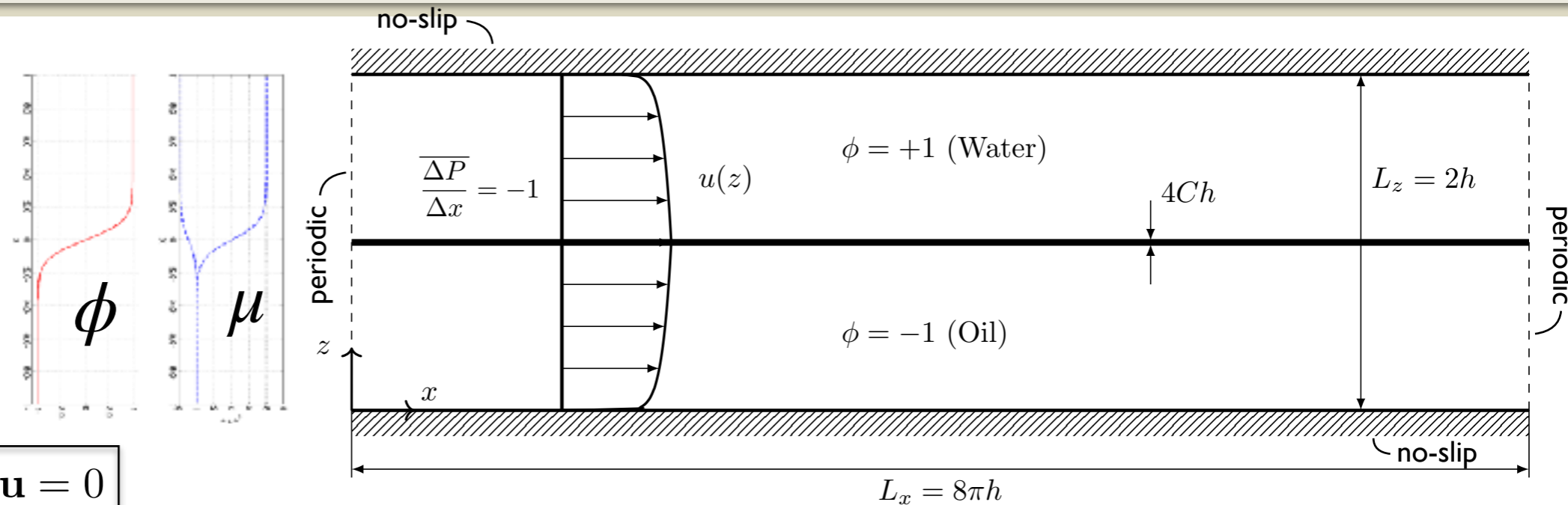
Interface

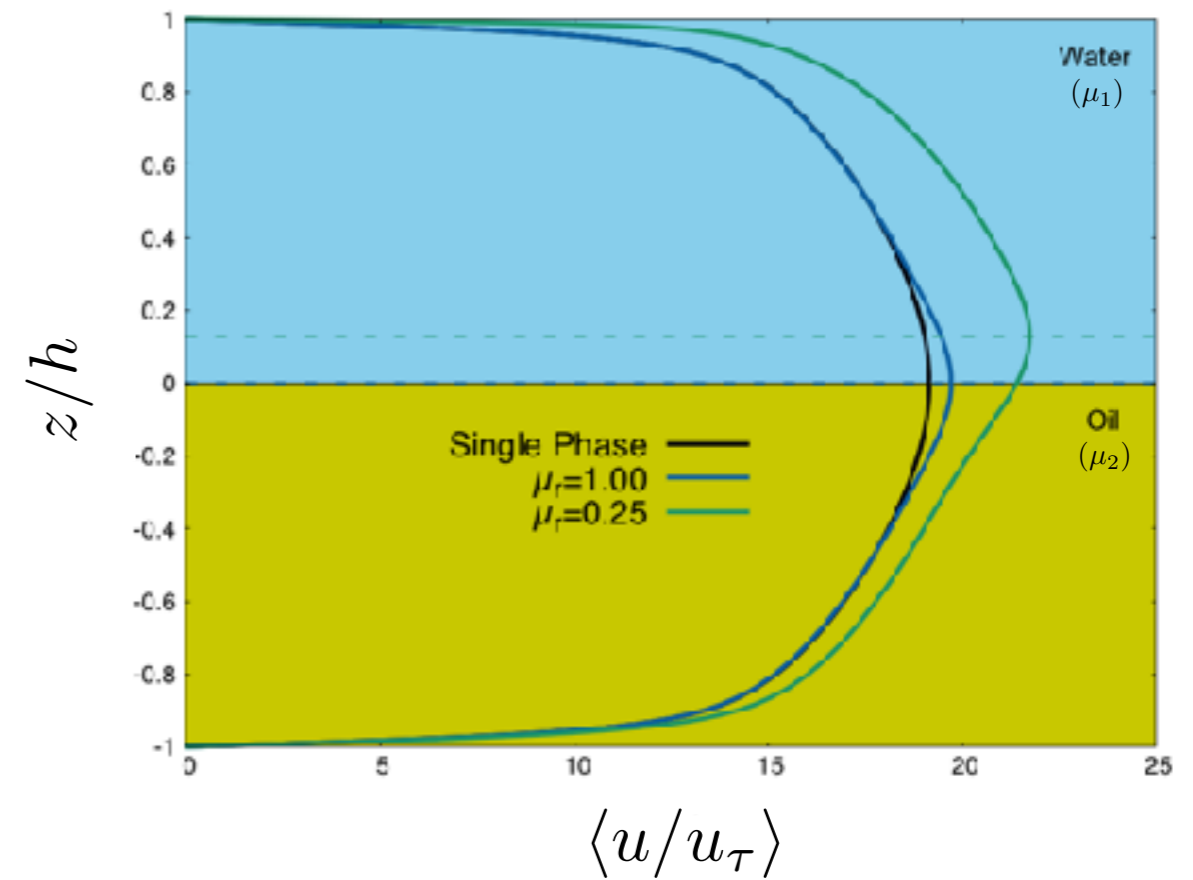
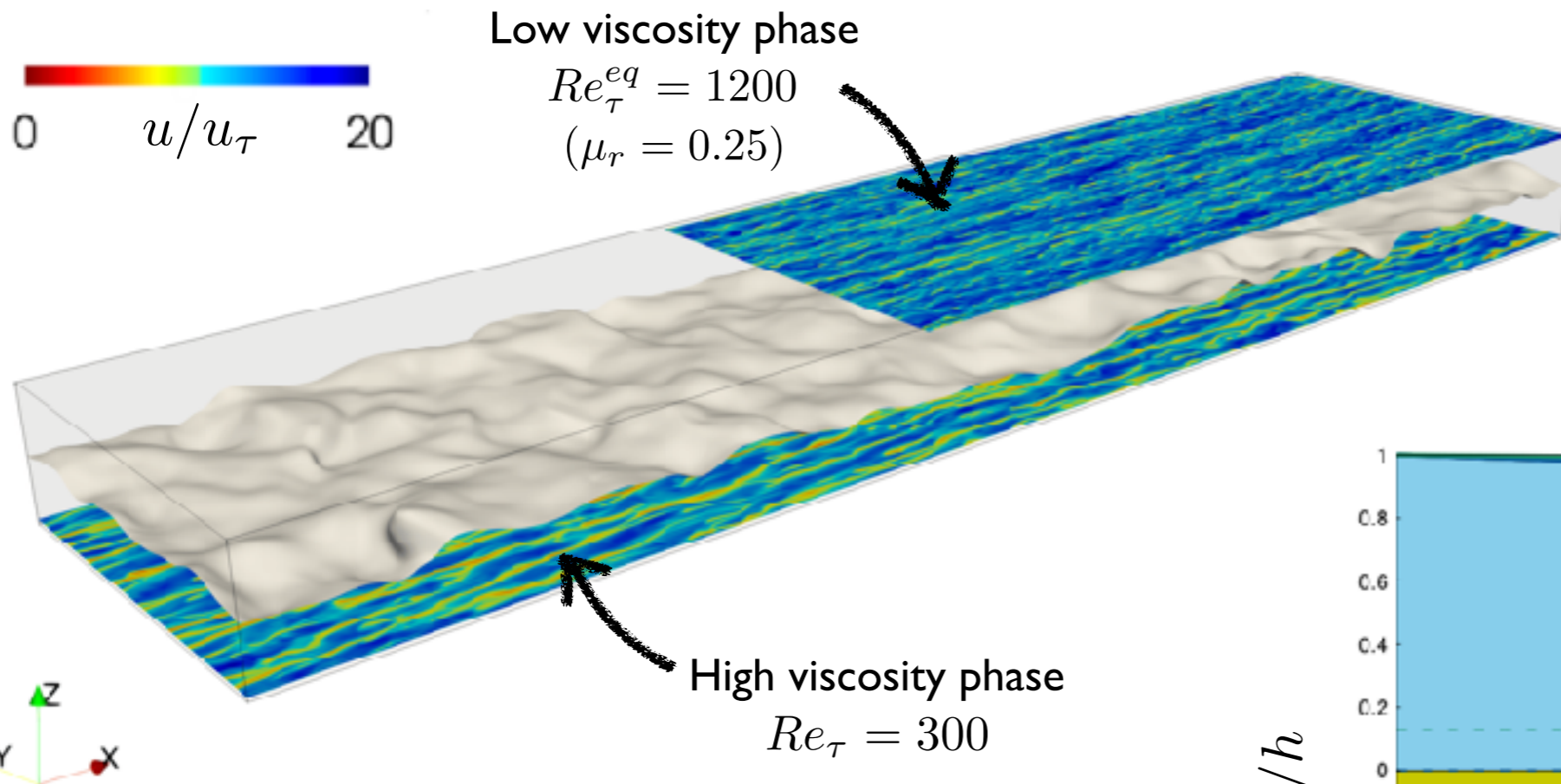
Advection-diffusion equation:

$$\frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla \phi = \frac{1}{Pe_\phi} \nabla^2 (\phi^3 - \phi - Ch^2 \nabla^2 \phi)$$

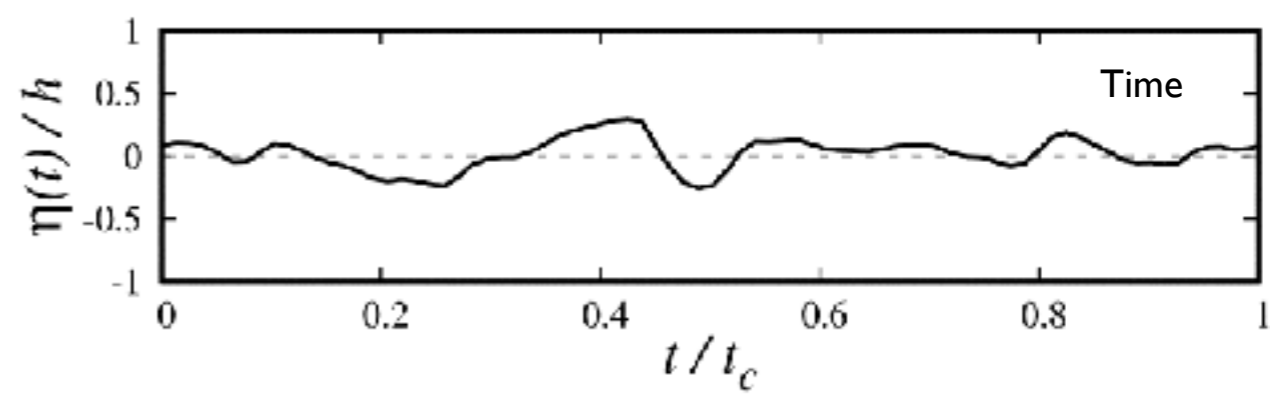
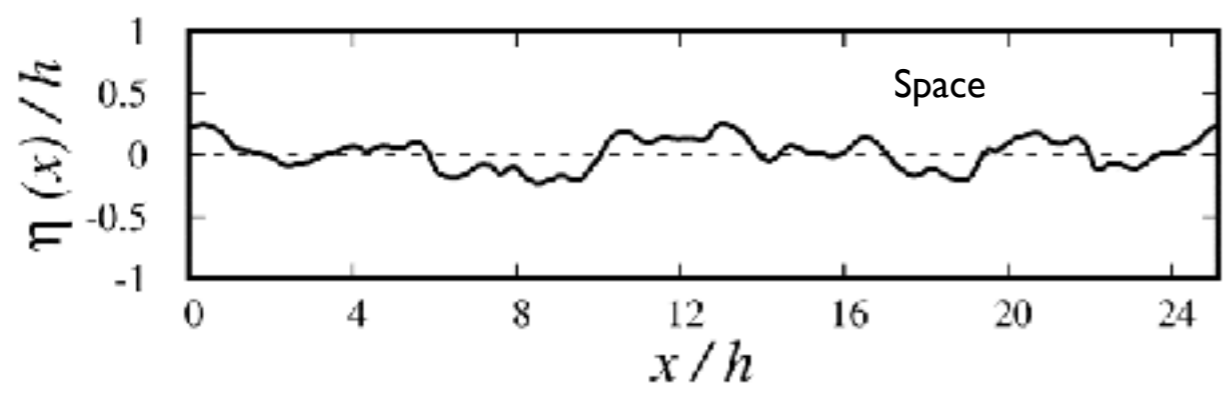
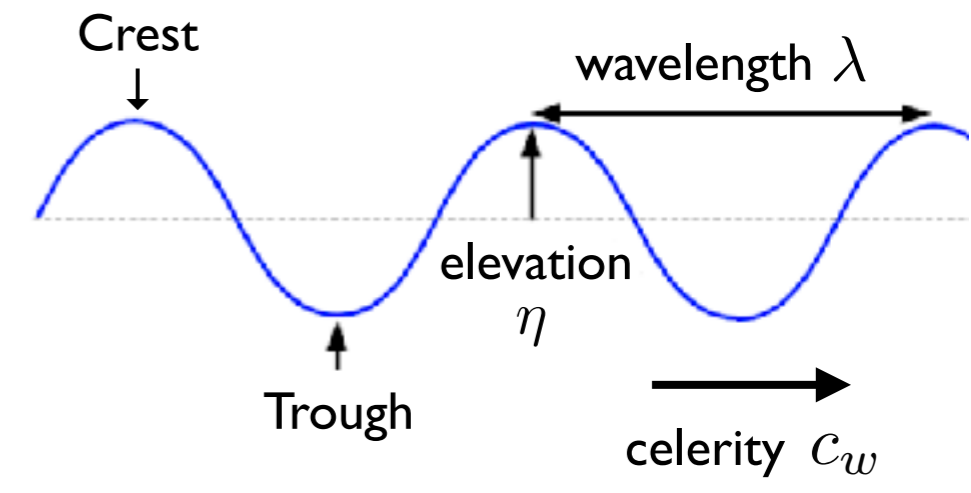
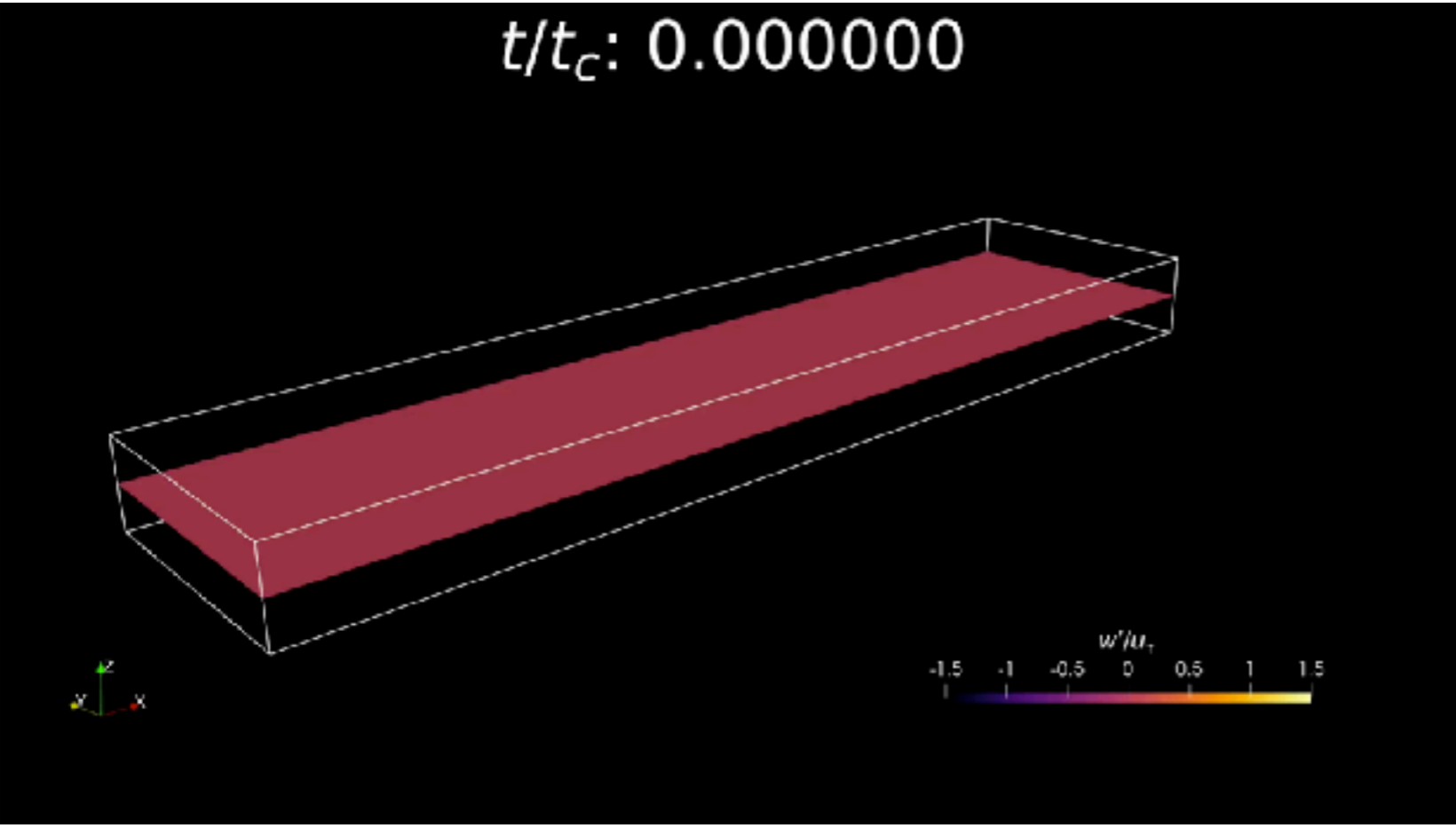
Cahn number: $Ch = \frac{\xi}{h}$

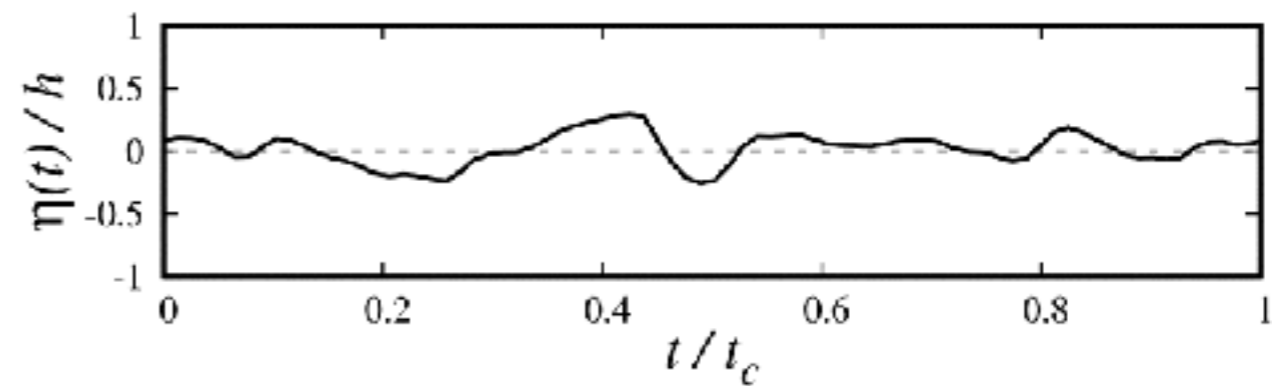
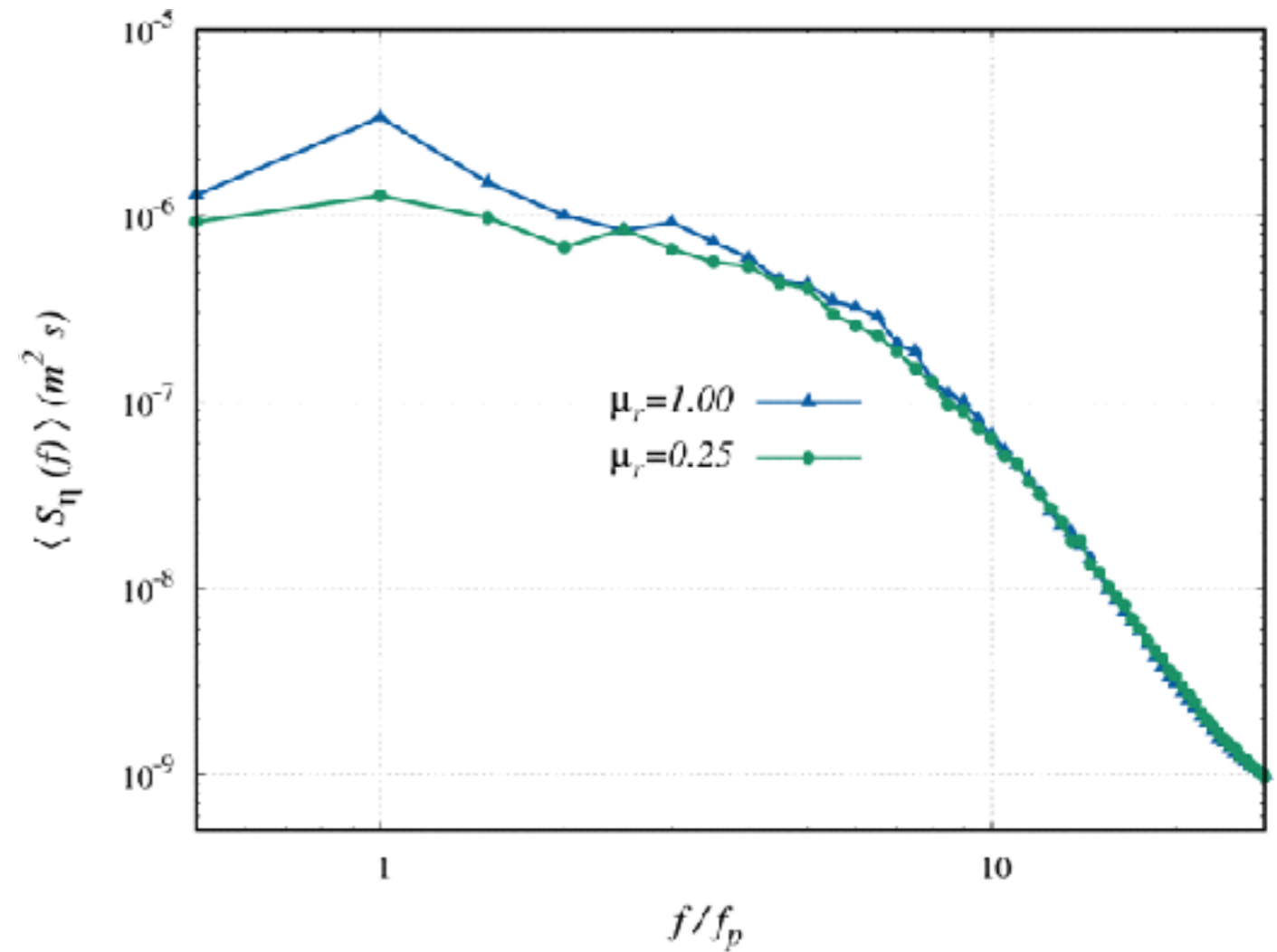
Péclet number: $Pe_\phi = \frac{u_\tau h}{\beta m_\phi}$





Simulation	Re_τ	We	N_x	N_z	N_y
Single Phase	300	-	1024	513	256
$\mu_r = 1.00$	300	1.0	1024	513	256
$\mu_r = 0.25$	300	1.0	2048	513	1024





* $f = c_w / \lambda$

[1] Zakharov, V. E. & Filonenko, N. N. 1967 Weak turbulence of capillary waves. Journal of Applied Mechanics and Technical Physics 8 (5), 37–40.

[2] Balkovsky, E., Falkovich, G., Lebedev, V. & Shapiro, I. Ya. 1995 Large-scale properties of wave turbulence. Phys. Rev. E 52, 4537–4540.

Wave Turbulence Theory predictions:

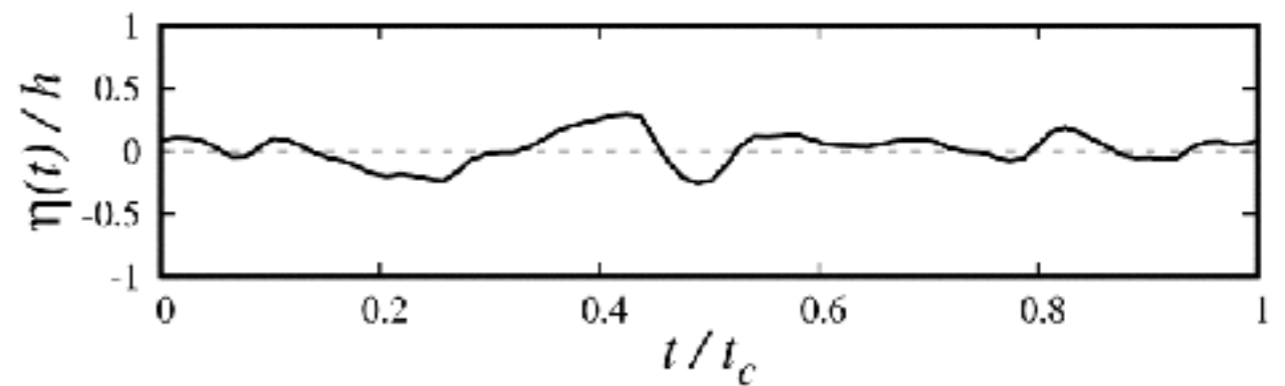
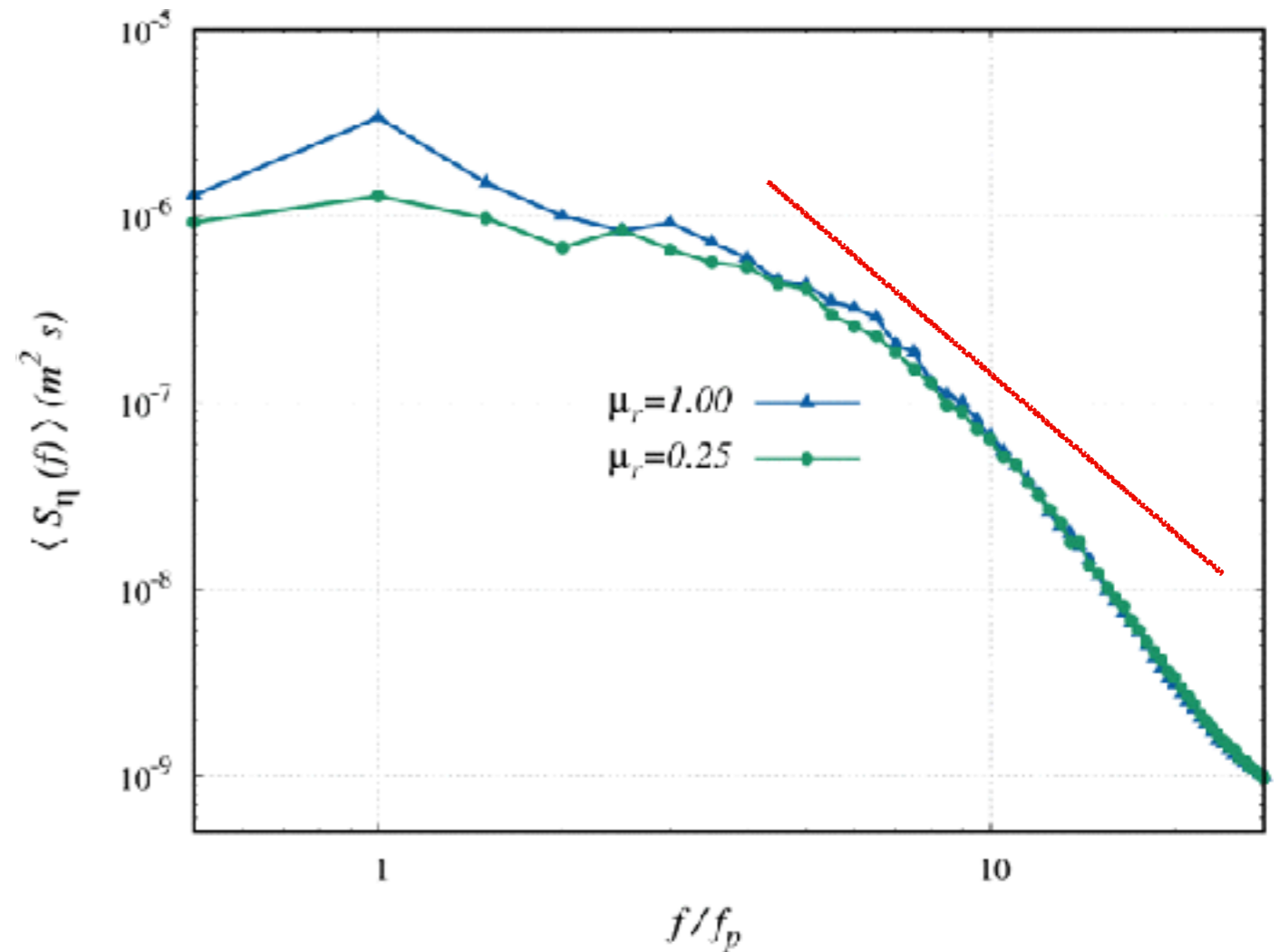
High-frequency waves [1]:

$$S_{\eta}(f) = f^{-8/3}$$

“Waves not at equilibrium -
energy transfer among wave scales”

Hp of Wave Turbulence Theory:
 Weak nonlinearities;
 Negligible dissipation;
 Inertial regime, far from injection and dissipation

* $f = c_w / \lambda$



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Wave Turbulence Theory predictions:

Low-frequency waves [2]:

$$S_{\eta}(f) = f^{-1}$$

“Waves at statistical equilibrium”

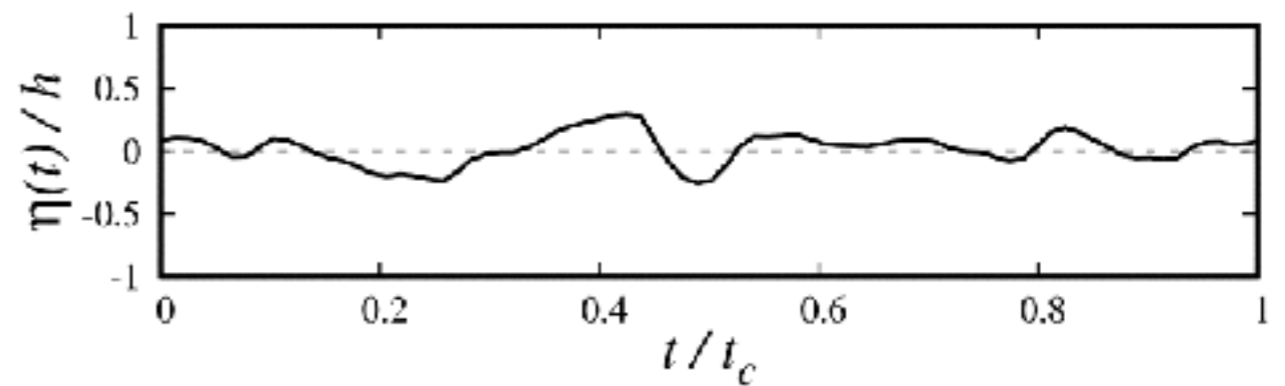
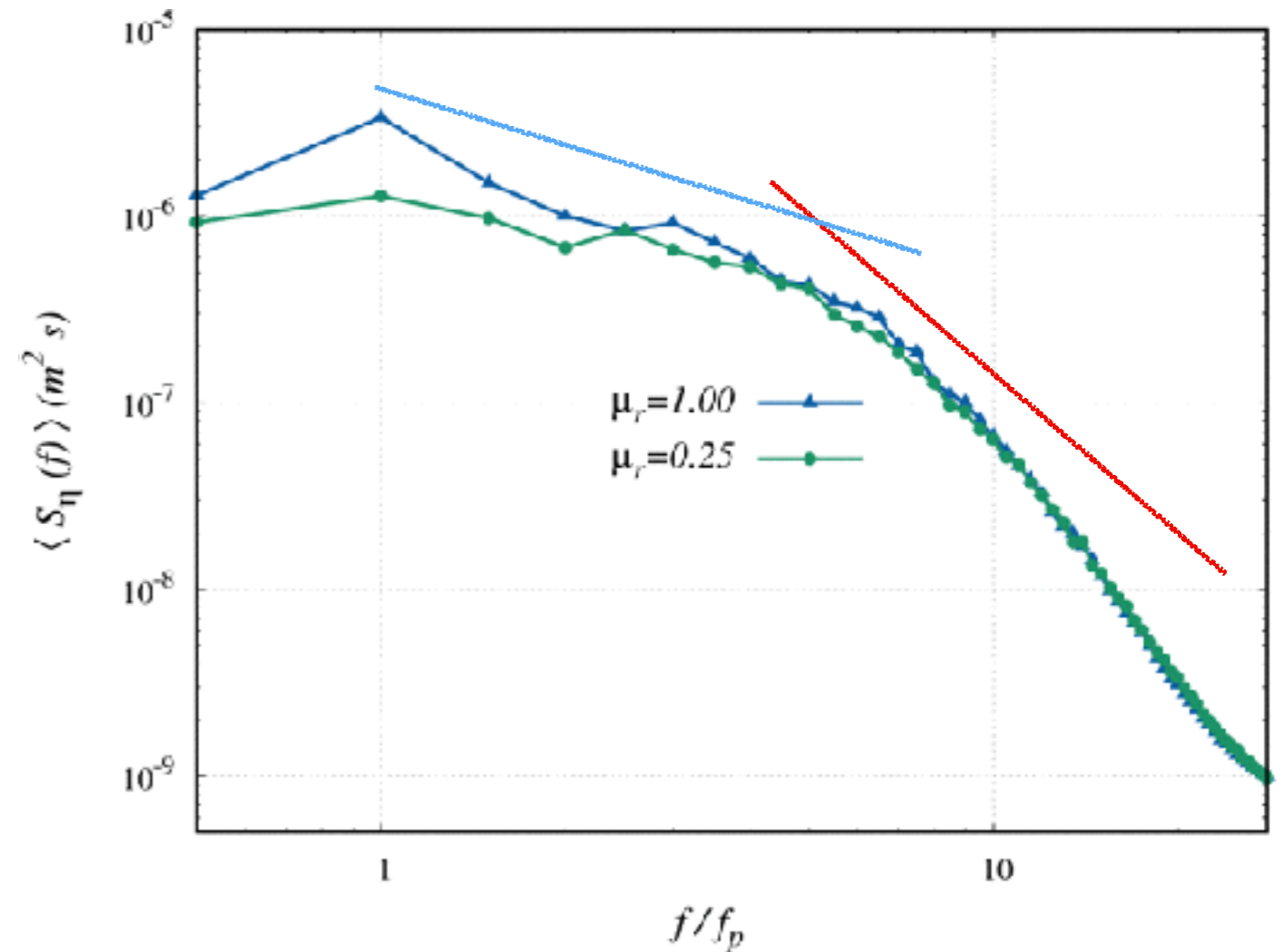
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[1] Zakharov, V. E. & Filonenko, N. N. 1967 Weak turbulence of capillary waves. Journal of Applied Mechanics and Technical Physics 8 (5), 37–40.

[2] Balkovsky, E., Falkovich, G., Lebedev, V. & Shapiro, I. Ya. 1995 Large-scale properties of wave turbulence. Phys. Rev. E 52, 4537–4540.

Wave Turbulence Theory predictions:

Large-scale waves [2]:

$$S_{\eta}(k) = k^{-1}$$

“Waves at statistical equilibrium”

First part (Super- Hinze scale waves [3]):

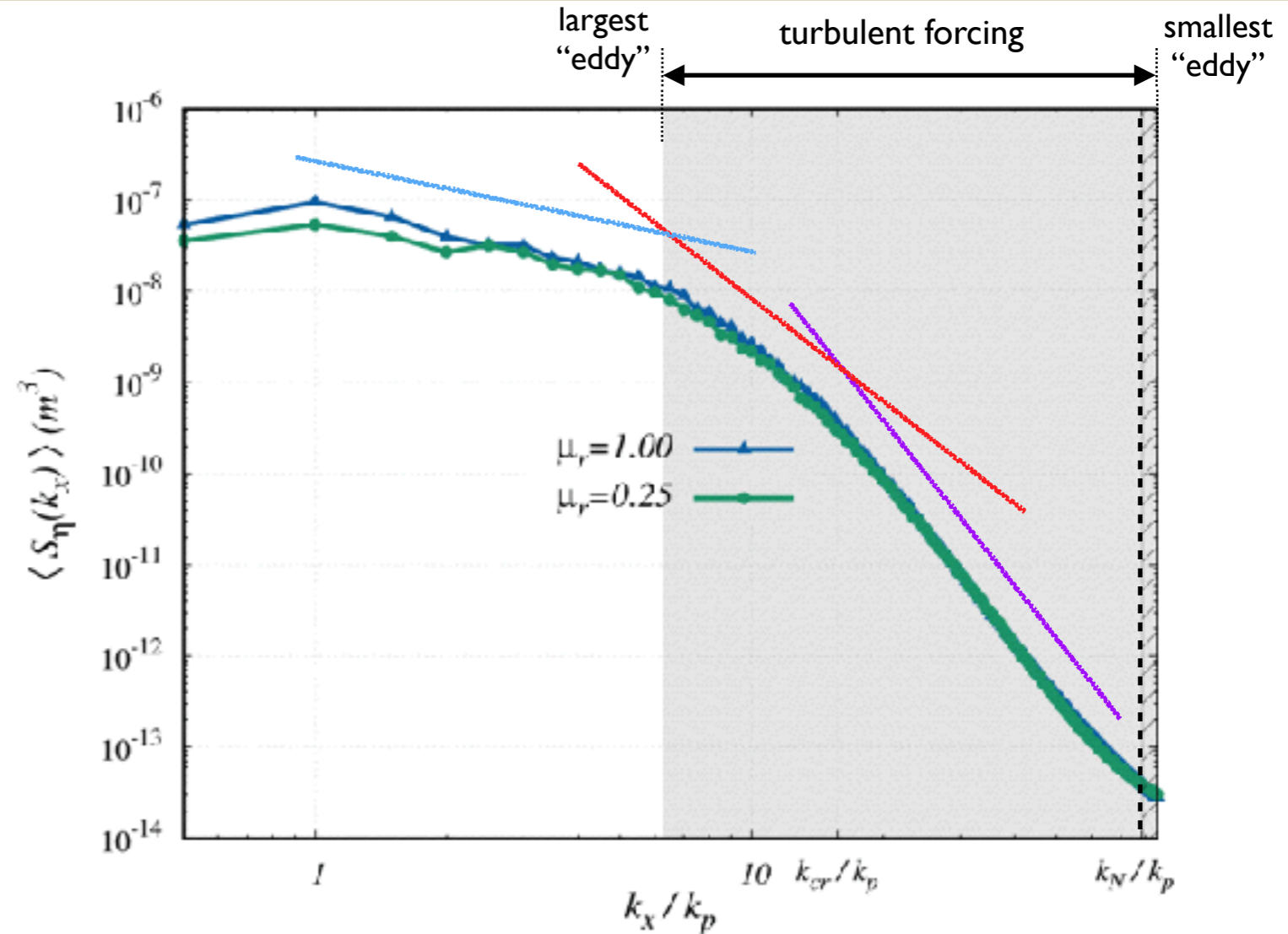
$$S_{\eta}(k) = k^{-4}$$

“Waves not at equilibrium - energy transfer among wave scales”

Second part (Sub- Hinze scale waves [4]):

$$S_{\eta}(k) = k^{-6}$$

“Strong damping due to surface tension”



--- Numerical resolution cut-off (shortest wave)

$$* k = 2\pi/\lambda$$

[1] Hinze, J. O. 1955 Fundamentals of the hydrodynamic mechanism of splitting in dispersion processes. *AIChE Journal* 1 (3), 289–295.

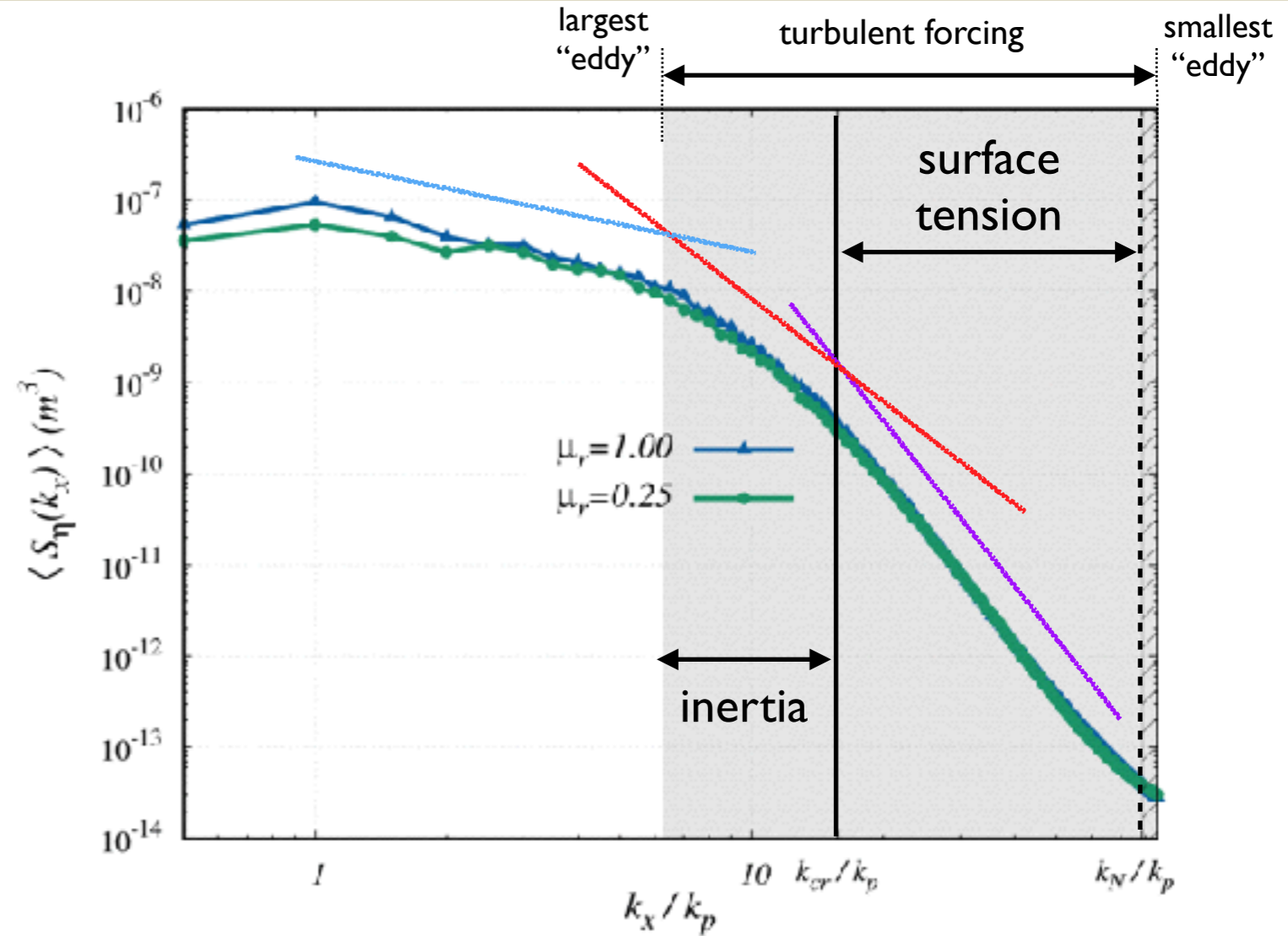
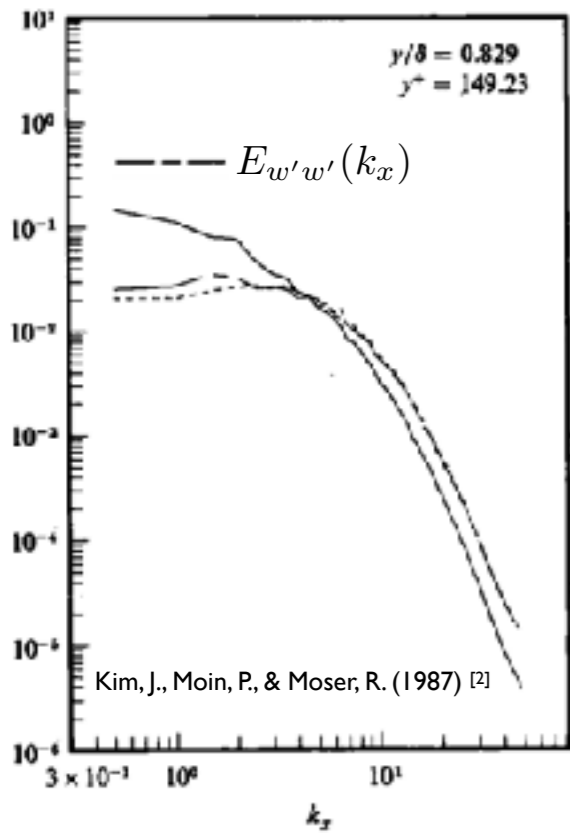
[2] Balkovsky, E., Falkovich, G., Lebedev, V. & Shapiro, I. Ya. 1995 Large-scale properties of wave turbulence. *Phys. Rev. E* 52, 4537–4540.

[3] Zakharov, V. E. & Filonenko, N. N. 1967 Weak turbulence of capillary waves. *Journal of Applied Mechanics and Technical Physics* 8 (5), 37–40.

[4] Savelsberg, R. & Van De Water, W. 2008 Turbulence of a free surface. *Phys. Rev. Lett.* 100, 034501

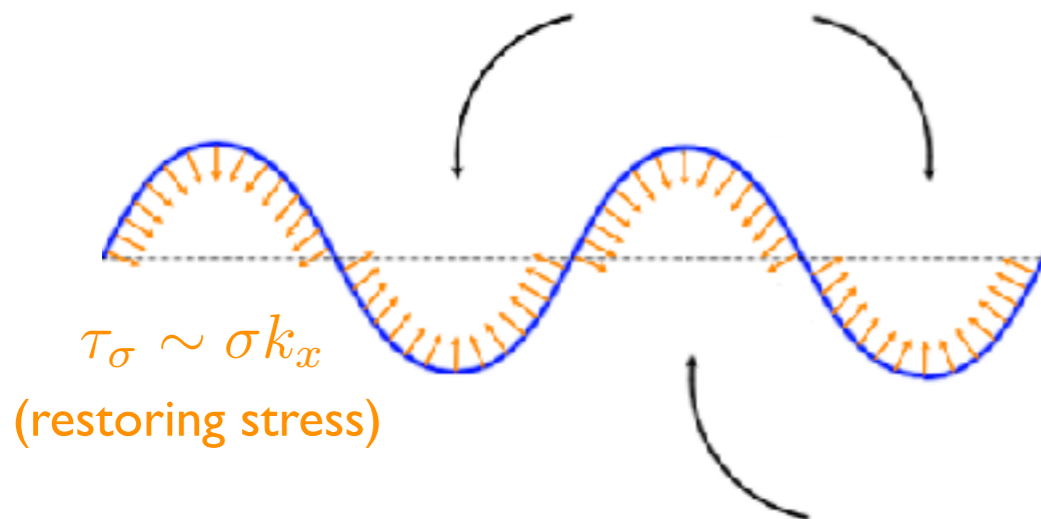
Wavenumber power spectra of wave elevation

Wave turbulence theory (WTT)



(perturbing pressure)

$$p' \sim \rho w'^2$$



— Kolmogorov- Hinze critical length scale [1]

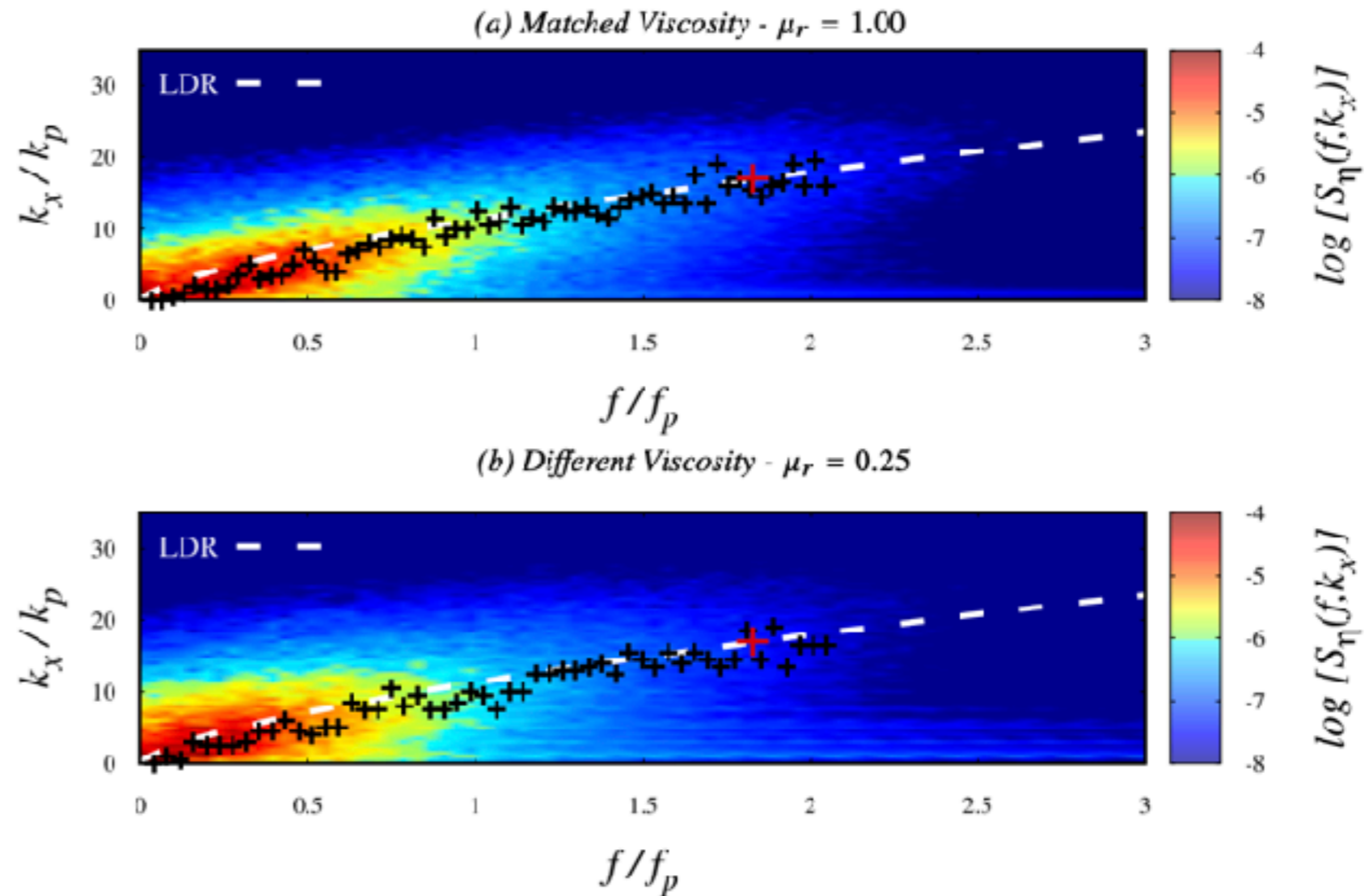
$$k_{cr} = 2\pi / (We^{-3/5} Re_{\tau}^{-2/5} |\epsilon_c|^{-2/5})$$

[2] Kim, J., Moin, P., & Moser, R. (1987). Turbulence statistics in fully developed channel flow at low Reynolds number. Journal of Fluid Mechanics, 177, 133-166.

Wave celerity (or frequency) is predicted by the Dispersion Relation (DR) [1]

$$\omega^2 = (\sigma/\rho)k^3 \tanh(kh) \quad (\text{wave celerity: } c_w = \omega/k)$$

Frequency -
wavenumber space



$$*\omega = 2\pi f$$

[1] Lamb H. 1932. Hydrodynamics. Springer-Verlag, Berlin

[2] Giamagas G, et al 2023. Journal of Fluid Mechanics, *In press*.

DNS of wall bounded stratified turbulence at $Re_\tau = 1000$ have been performed for different stratification levels, $Ri_\tau = 0 - 300$

Compared to forced convection, **stratification reduces mixing and wall normal transport of momentum/heat**

The stratified flow becomes "less homogenous"; **Internal waves are generated** at the channel center

Although velocity and temperature fluctuations may be large at the channel center, **the buoyancy flux is monotonically decreasing**

DNS of viscosity stratified oil-water channel turbulence has been performed

Analysis of interfacial waves shows: **energy equipartition** at scales larger than forcing

Kolmogorov-Hinze scale marks the threshold between growing waves and waves damped by surface tension

Propagation of waves in agreement with **linear dispersion relation**



Perspectives:

Thermally-stratified turbulence: Tracking particles/species in thermally stratified turbulence

Viscosity stratified turbulence: Adding surfactants; Wave breaking

CINECA



Thank you !