

Two-dimensional Vorticity Dynamics

2D Decaying Turbulence

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512 x 512 resolution
narrow band random initial condition

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19.00

Two-dimensional Vorticity Dynamics

$$\vec{u} = \begin{pmatrix} u(x, y, t) \\ v(x, y, t) \\ 0 \end{pmatrix} \quad \omega = \begin{pmatrix} 0 \\ 0 \\ \omega(x, y, t) \end{pmatrix} \quad \omega = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$$

Consider the stream function $\psi(x, y)$

$$\vec{A} = \psi(x, y) \vec{e}_z \quad \vec{u} = \nabla \times \vec{A}$$

$$\vec{\omega} = -\Delta \psi \vec{e}_z$$

Biot and Savart Law: 2D Flows

$$\psi(x, y, t) = -\frac{1}{2\pi} \iint \omega(x', y', t) \ln(R) dx' dy' + \psi_p$$

$$R^2 = (\vec{r} - \vec{r}')^2 = (x - x')^2 + (y - y')^2$$

$$\vec{u} = \frac{1}{2\pi} \iint \frac{\omega(x', y', t) \vec{e}_z \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^2} dx' dy' + \nabla \phi$$

For large r it decreases algebraically $\Gamma/(2\pi r)$

Viscous Flow $\nu \neq 0$

$$\frac{D\omega}{Dt} = \nu \Delta \omega$$

Diffusion of Vorticity

Cancellation of Vorticity

$$\frac{\partial \omega}{\partial t} = \nu \frac{\partial^2 \omega}{\partial y^2}$$

A unique vortex layer $\omega = \frac{\Gamma}{\sqrt{4\nu t}} \exp\left(-\frac{y^2}{4\nu t}\right)$

$$t = 0 \quad \omega \sim \Gamma \delta(t)$$

Vorticity diffuses by Viscous Effect

Circulation is conserved

$$\frac{\partial \omega}{\partial t} = \nu \frac{\partial^2 \omega}{\partial y^2}$$

An array of vortex layers $\omega = \omega_0 \sin(ky)$

$$\omega(y, t) = \omega_0 \exp(-k^2 t) \sin(ky)$$

Vorticity Cancellation by Viscous Effect

The total vorticity field $\nu \neq 0$

$$\Gamma = \int \int \omega dS$$

$$\vec{OC} = \frac{\int \int \omega \vec{OM} dS}{\int \int \omega dS}$$

C : Position of the vorticity centroid

$$a^2 \equiv \frac{\int \int \omega |\vec{CM}|^2 dS}{\int \int \omega dS}$$

Size of total vorticity repartition

$$a^2 = a_0^2 + 4\nu t$$

Unique Vortex $\nu \neq 0$

$$\Gamma = \int \int \omega dS$$

$$\vec{OC} = \frac{\int \int \omega \vec{OM} dS}{\int \int \omega dS} \quad \text{C : Position of the vortex center the centroid}$$

$$a^2 \equiv \frac{\int \int \omega |\vec{CM}|^2 dS}{\int \int \omega dS} \quad \text{Size of the vortex}$$

$$a^2 = a_0^2 + 4\nu t$$

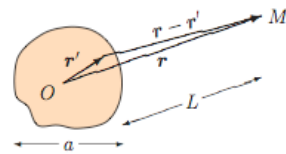
Inviscid vortex flow $\nu = 0$

Conservation of vorticity pointwise

In an assembly of vortices this means :

⇒ Conservation of circulation of any given vortex

⇒ Size of any given vortex
(only for a uniform vorticity)



$$\vec{O}C = \frac{\int \int \omega O\vec{M} dS}{\int \int \omega dS}$$

$$|\vec{r}'| \ll |\vec{r}|$$

$$\begin{aligned} \ln(R) &= \frac{1}{2} \ln |\vec{r} - \vec{r}'|^2 = \frac{1}{2} \ln \left[r^2 \left(1 - \frac{2}{r^2} \vec{r} \cdot \vec{r}' \right) \right] \\ &\sim \ln(r) - \frac{\vec{r} \cdot \vec{r}'}{r^2} + \mathcal{O}\left(\frac{a^2}{L^2}\right) \end{aligned}$$

$$\psi = -\frac{\ln(r)}{2\pi} \underbrace{\int \int \omega dx' dy'}_{\Gamma} + \frac{\vec{r}}{2\pi r^2} \underbrace{\int \vec{r}' \omega dx' dy'}_{\vec{r}_C = \vec{0}} + \mathcal{O}\left(\frac{a^2}{r^2}\right)$$

Potential vortex flow

$$\psi \sim \underbrace{-\frac{\Gamma}{2\pi} \ln(r)}_{\text{point vortex}} + \mathcal{O}\left(\frac{a^2}{r^2}\right)$$

At a point far away from a vortex location,
the velocity field produced by a vortex is similar to a

point-vortex velocity $\omega = \Gamma \delta(x - x_c) \delta(y - y_c)$

Potential vortex flow

$$u_r(r) = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = 0, \quad u_\theta(r) = -\frac{\partial \psi}{\partial r} = \frac{\Gamma}{2\pi r}$$

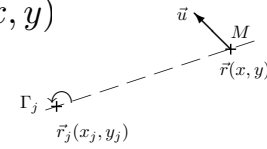
For the velocity field at M, only two quantities matter

The distance OM to the vortex centroid O

The vortex circulation Γ

Velocity due to vortex j at point $\vec{r}(x, y)$

$$\vec{u}_j = \frac{\Gamma}{2\pi |\vec{r} - \vec{r}_j|^2} \begin{pmatrix} -(y - y_j) \\ x - x_j \\ 0 \end{pmatrix}$$



A system of point vortices : Dynamical system

$$\begin{pmatrix} \dot{x}_i \\ \dot{y}_i \end{pmatrix} = \sum_{j \neq i} \frac{\Gamma_j}{2\pi R_{ij}^2} \begin{pmatrix} -(y_i - y_j) \\ x_i - x_j \end{pmatrix}, \quad R_{ij}^2 = |\vec{r}_i - \vec{r}_j|^2$$

This system is an Hamiltonian system

$$H = -\frac{1}{4\pi} \sum_{k=1}^N \sum_{j \neq k} \Gamma_k \Gamma_j \ln(R_{kj})$$

Interaction energy
Kinetic energy is infinite

$$q_i \equiv x_i \quad p_i \equiv \Gamma_i y_i$$

$$\dot{q}_i = \frac{\partial H}{\partial p_i}, \quad \dot{p}_i = -\frac{\partial H}{\partial q_i}$$

H is a conserved energy

Other Conserved Quantities

A centroid

$$\begin{pmatrix} x_C \\ y_C \end{pmatrix} = \frac{1}{\Gamma} \sum_{i=1}^N \Gamma_i \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

A dispersion radius

$$d^2 = \frac{1}{\Gamma} \sum_{i=1}^N \Gamma_i [(x_i - x_C)^2 + (y_i - y_C)^2]$$



System of Two point vortices

4 dynamical variables x_1, y_1, x_2, y_2

3 Conserved quantities

$$H = -\frac{1}{2\pi} \Gamma_1 \Gamma_2 \ln(R_{12}) = Cst \quad \Rightarrow \quad R_{12} = |\vec{r}_1 - \vec{r}_2| = Cst$$

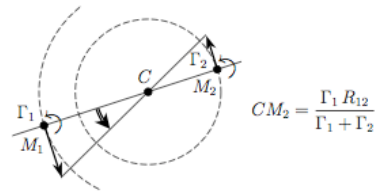
The dispersion radius gives again $R_{12} = |\vec{r}_1 - \vec{r}_2| = Cst$

$$\Gamma_1 y_1 + \Gamma_2 y_2 = cte \quad \Gamma_1 x_1 + \Gamma_2 x_2 = cte$$

Two point vortices

$$\Gamma_1 \Gamma_2 > 0$$

C is located between 1 and 2.

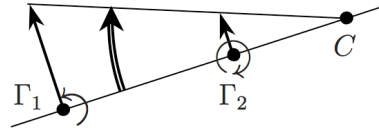


The system rotates around C at angular velocity

$$\Omega = \frac{\Gamma_1}{2\pi R_{12}} \frac{\Gamma_1 + \Gamma_2}{\Gamma_1 R_{12}} = \frac{\Gamma_1 + \Gamma_2}{2\pi R_{12}^2}$$

Two point vortices

$$\Gamma_1 \Gamma_2 < 0$$



C is located on the side of the more intense vortex

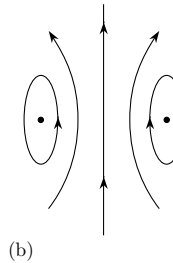
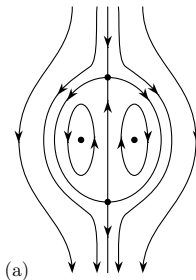
Two Point Vortices: a Dipole $\Gamma_1 + \Gamma_2 = 0$

The dipole does not rotate, and translates at velocity

$$\Gamma / (2\pi R_{12})$$

streamlines in the comoving Frame

streamlines in the Laboratory



Trapping and Transporting Fluid

Three Point Vortices

Simple dynamics : an integrable motion

Triangle shape oscillates with a period T_1

Triangle rotates with a period T_2

$\frac{T_1}{T_2}$ is rational number (initial condition) \Rightarrow periodic

$\frac{T_1}{T_2}$ is in general not in an irrational number
 \Rightarrow Quasi-periodic Motion

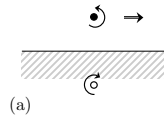
Four or more point Vortices

8 dynamical variables

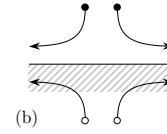
4 Conserved quantities

The dynamical system is non integrable: chaotic

Point Vortices with a Wall



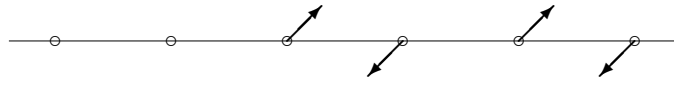
(a) A unique vortex along a wall



(b) A Dipole along a wall

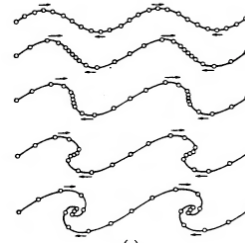
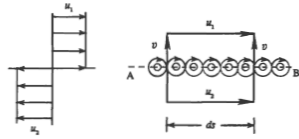
Image theory

Infinite Row of Point Vortices

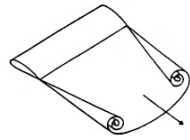
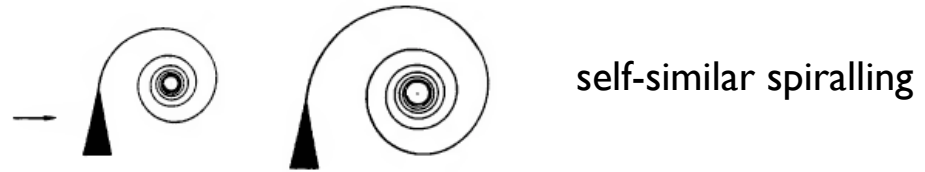


Pairing instability

From an Infinite Row of Point Vortices towards Vortex Sheet

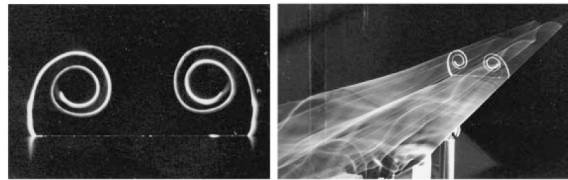


Vortex Sheet

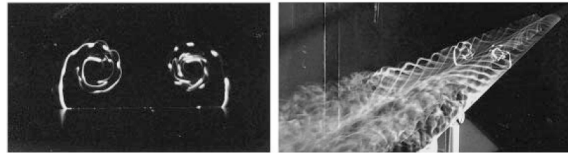


Roll-up behind a wing

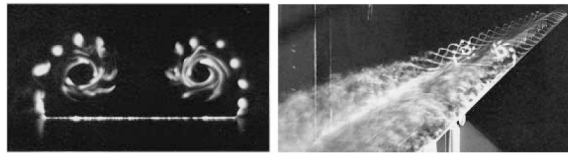
Vortex sheet near trailing edges of a Delta wing



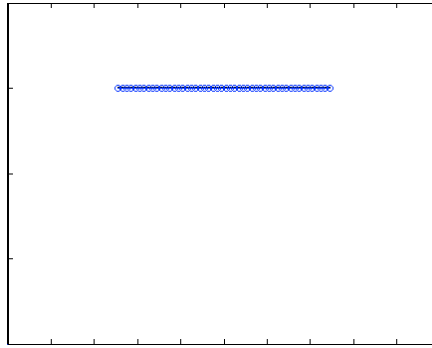
(a) $Re_x = 25\,200, Re_c = 41\,900$



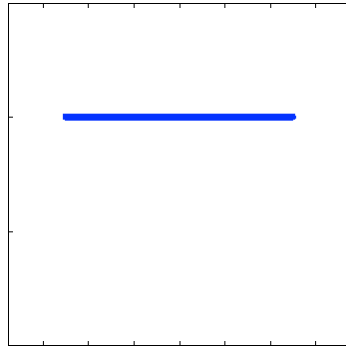
(b) $Re_x = 61\,100, Re_c = 101\,900$



(c) $Re_x = 118\,000, Re_c = 196\,700$



A First Vortex sheet with zero total circulation



Another Vortex sheet with non nul vortex
circulation

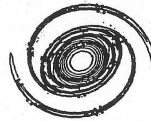
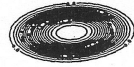
EXTENDED 2D VORTICES

Vortex axisymmetrization

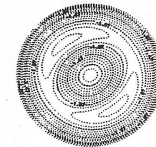
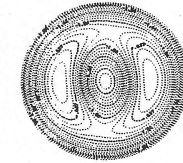
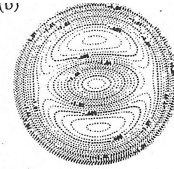
Vorticity

Streamfunction

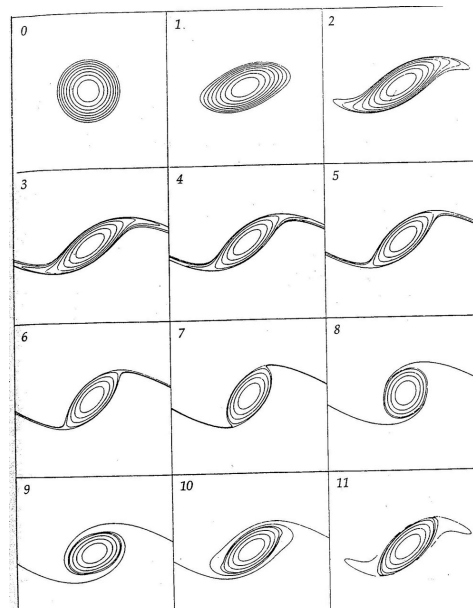
(a)



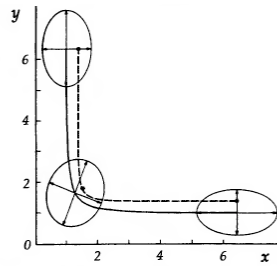
(ψ)



Vortex in a Potential Strain field

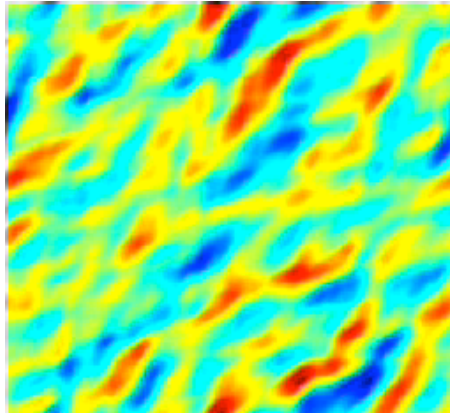


Stripping



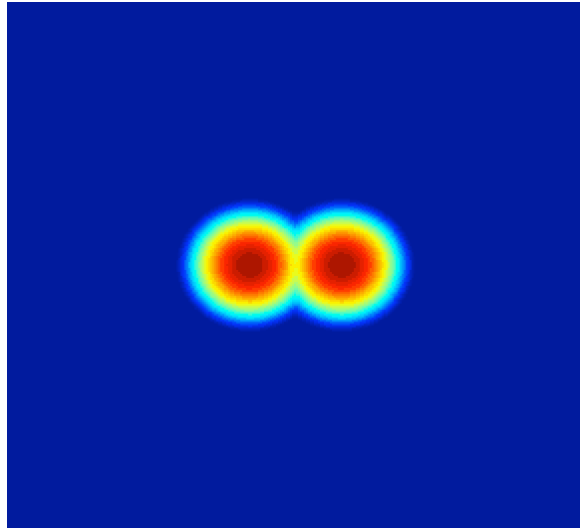
Vortex along a wall or a Dipole

Merging of Vortices



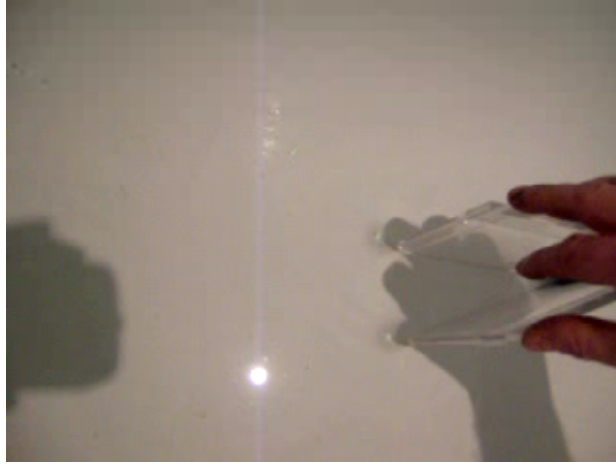
Two-dimensional Turbulence : Inverse Cascade
Gerris (S.Popinet)

Vortex Merging



Josserand and Rossi 2005

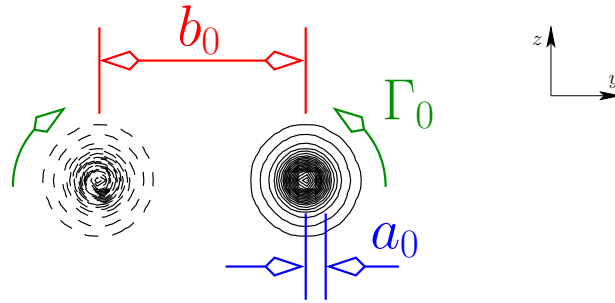
Vortex Dipole

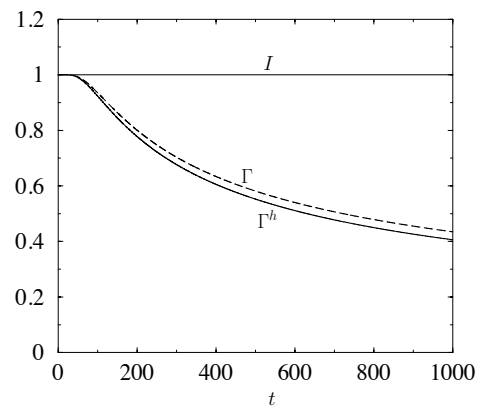


Vortex In a bath

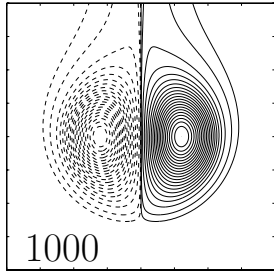
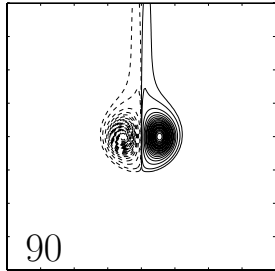
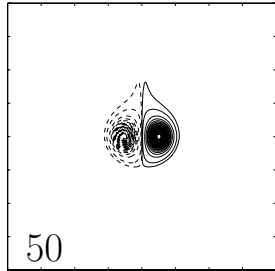
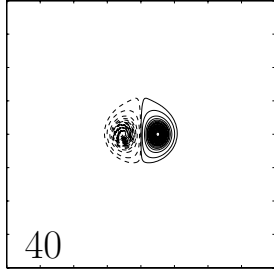
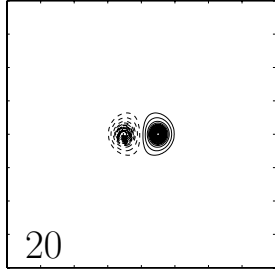
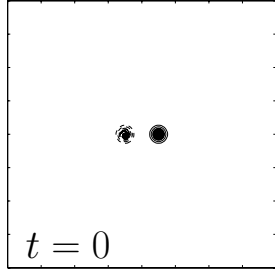
$$\omega(y, z, t = 0) = \omega^{(+)}(y, z) + \omega^{(-)}(y, z)$$

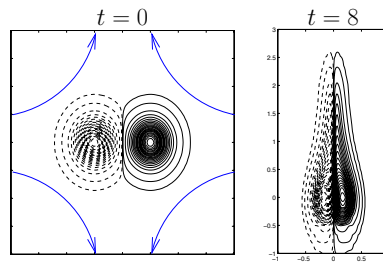
$$\omega^{(\pm)}(y, z) = \pm \frac{\Gamma_0}{\pi a_0^2} \exp\left(-\frac{(y \mp b_0/2)^2 + z^2}{a_0^2}\right)$$





Cancellation of Vorticity





Strain dipole in the Plane



Numerical Simulation Shallow Water equations
(Code Gerris S.Popinet)